

Cosmic Large-scale Structure Formations

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18 weeks

outline

Background (1 w)

- universe geometry and matter components (1 hr)
- Standard candle (SNIa) (0.5 hr)
- Standard ruler (BAO) (0.5 hr)

Linear perturbation (9 w)

- relativistic treatment perturbation (2 hr)
- primordial power spectrum (2 hr)
- linear growth rate (2 hr)
- galaxy 2-pt correlation function (2 hr)
- Baryon Acoustic Oscillation (BAO) (2 hr)
- Redshift Space Distortion (RSD) (2 hr)
- Weak Lensing (2 hr)
- Einstein-Boltzmann codes (2 hr)

Non-linear perturbation (6 w)

- Non-linear power spectrum (2 hr)
- halo model (2 hr)
- N-body simulation algorithms (2 hr)
- Press-Schechter (PS) halo mass function (2 hr)
- Extended-PS (EPS) halo mass function (2 hr)
- halo bias & halo density profile (2 hr)

Statistical analysis (2 w)

- Monte-Carlo Markov Chain sampler (2 hr)
- CosmoMC use (2 hr)

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outline

Non-linear perturbation (6 w)

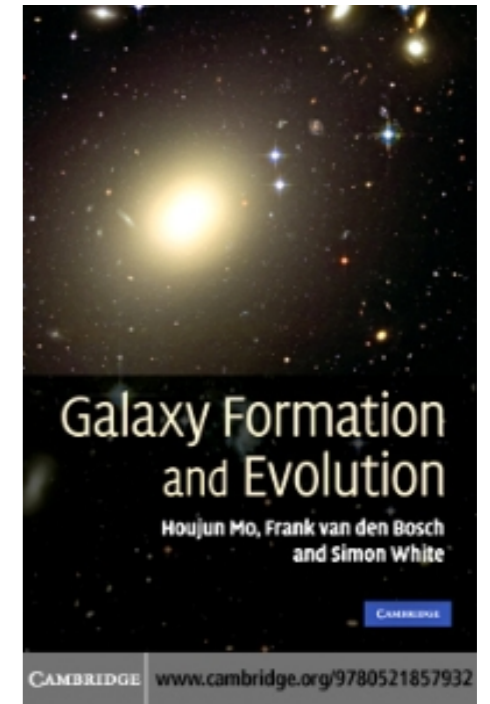
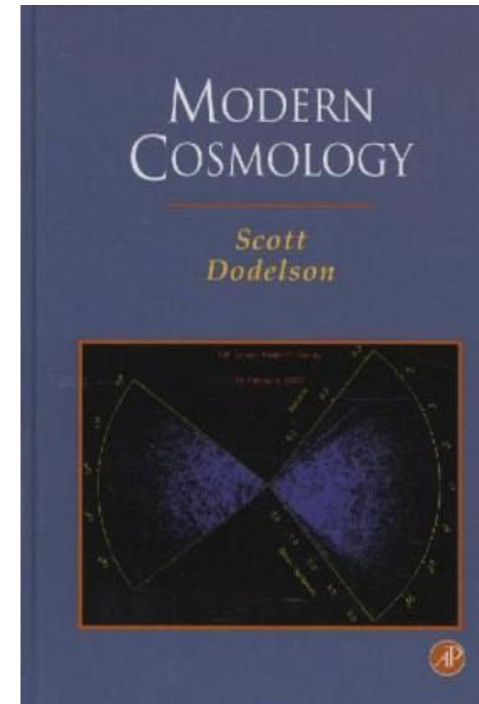
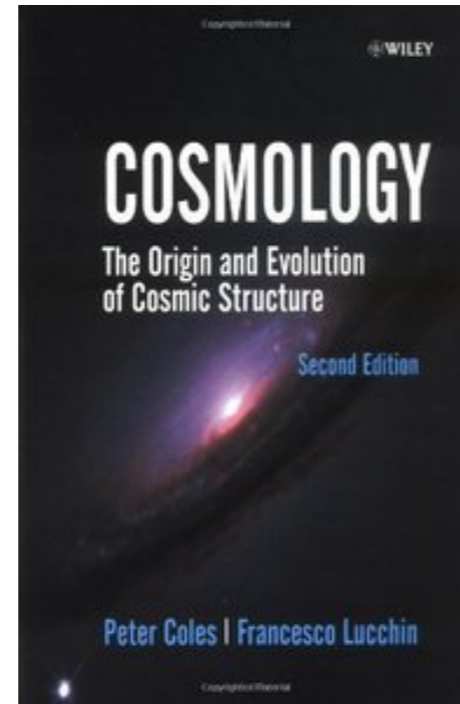
- Non-linear power spectrum (2 hr)
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Statistical analysis (2 w)

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1. Gaussian Random Field/ Power spectrum/Correlation function/ Phase
2. BAO
3. Galaxy Clustering
4. RSD
5. Lensing: WL/ Strong Lensing
6. Linear Growth
7. Nonlinear growth (spherical collapse)
8. Halo model: Press-Schechter formalism, merge tree

references



- 宇宙大尺度结构的形成 向守平、冯珑珑
- Cosmology Peter Coles & Francesco Lucchin
- Modern Cosmology Scott Dodelson
- Galaxy Formation and Evolution Houjun Mo & van den Bosch and Simon White
- Baumann lecture note

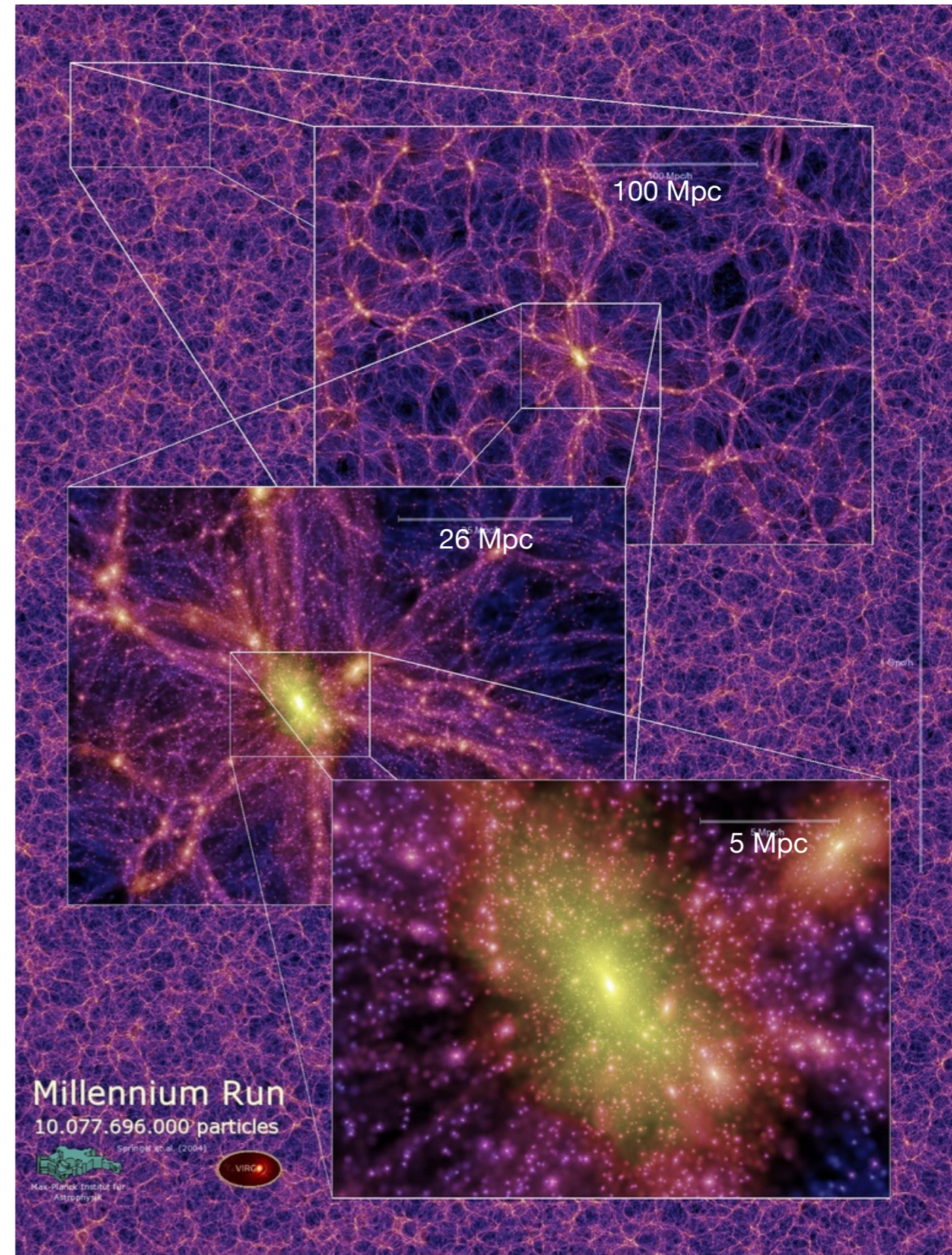
<http://101.96.8.165/www.damtp.cam.ac.uk/user/db275/Cosmology/Lectures.pdf>

成绩计算（百分制）

1. 平时作业：40%
2. 期末随堂考试：60%

Lecture 1

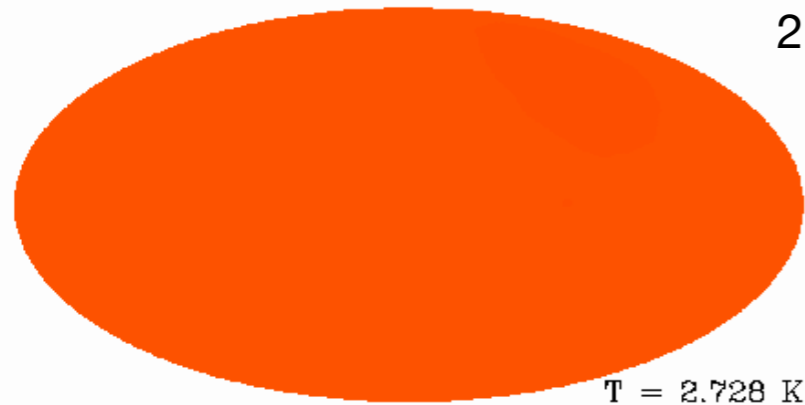
non-relativistic matter distribution



1. cosmological principle/CP

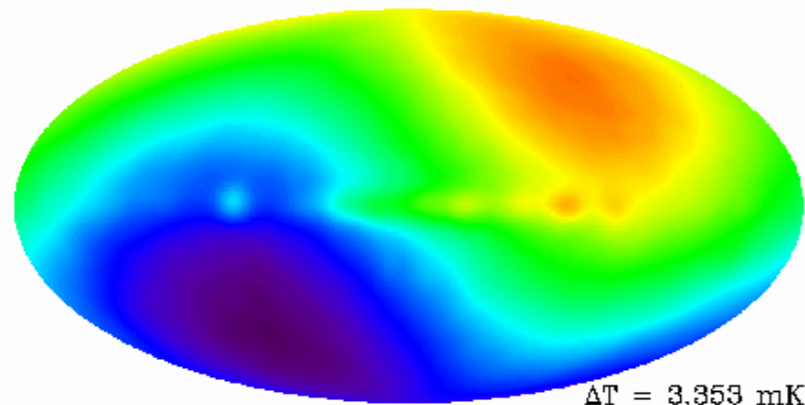
For a **co-moving observer**, on the **large** scale, the universe is **homogenous** and **isotropic**.

1. Observer: co-move with the background expansion

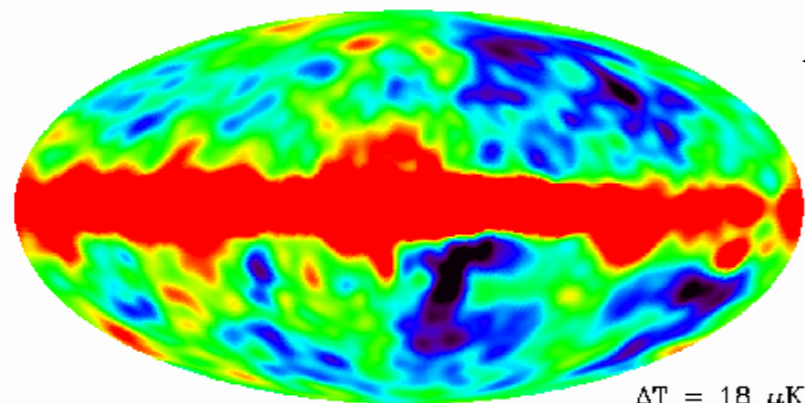


2. On this scale (> 1 Mpc):
each galaxy is like a test particle

[milky way ~ 15 kpc,
1 pc ~ 3 ly]



relativistic
photon
distribution



Galaxy clustering



2. FRWL metric [Friedmann–Robertson–Walker-Lemaître]

$$[ds^2 = g_{\mu\nu} dx^\mu dx^\nu] \quad g_{\mu\nu} \rightarrow \text{describe the d.o.f. gravity sector}$$

similar to: E/B field in Maxwell eq.

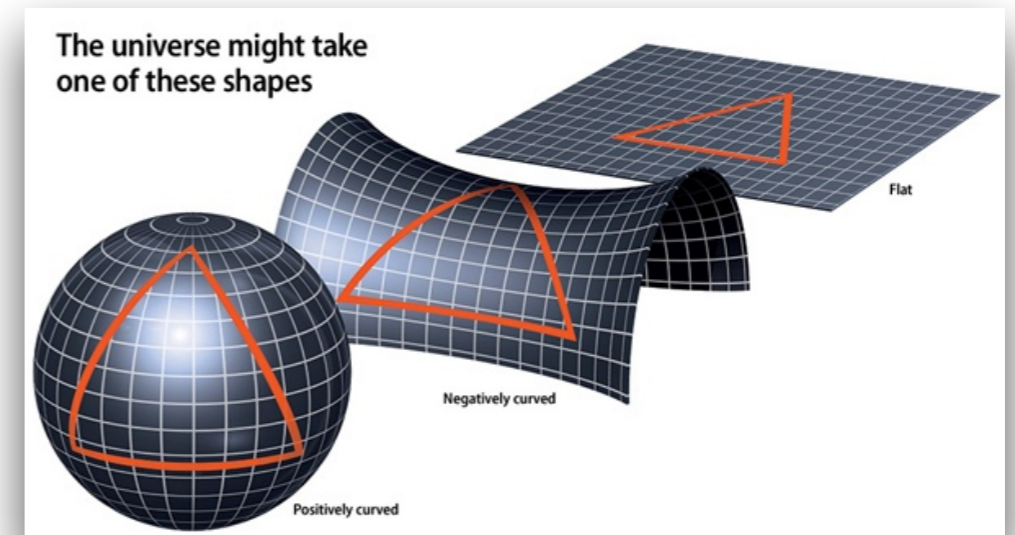
$$\begin{aligned} \nabla \cdot \vec{E} &= \frac{\rho}{\epsilon_0} \\ \nabla \cdot \vec{B} &= 0 \\ \nabla \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} \\ \nabla \times \vec{B} &= \mu_0 \vec{J} + \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t} \end{aligned}$$

$$\xrightarrow[\text{format}]{\text{covariant}} \quad F^{\mu\nu}{}_{;v} = J^\mu$$

$$[c=1] \quad ds^2 = -dt^2 + a^2(t) \left[\frac{dr^2}{1 - K(t)r^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\varphi^2 \right]$$

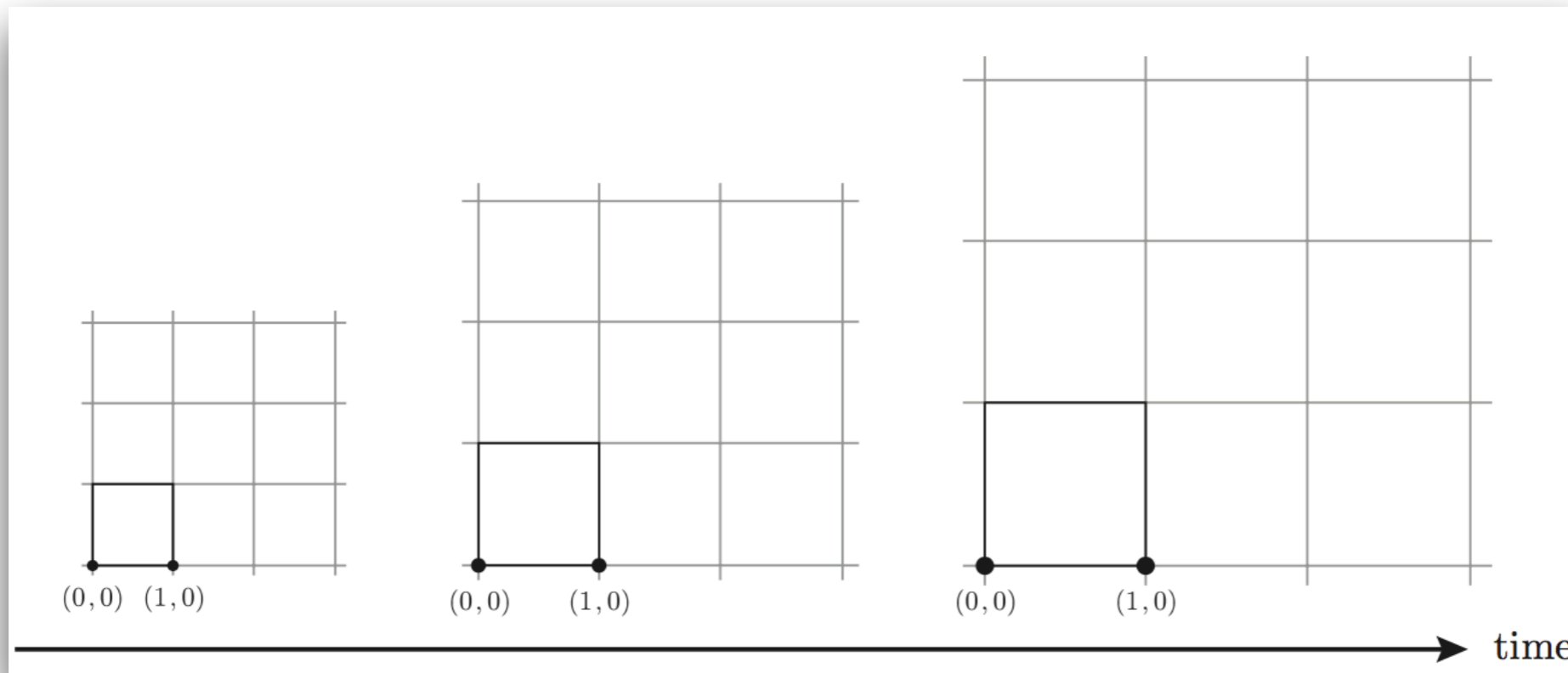
3 spatial curvature

$$K(t) = \begin{cases} > 0 \rightarrow (\textit{close}) \\ 0 \rightarrow (\textit{flat}) \\ < 0 \rightarrow (\textit{open}) \end{cases}$$



$a(t)$ [scale factor 标度因子]: tells the physical size of the universe.

cosmic redshift: $a = 1 / (1 + z)$ or $z = 1/a - 1$ with $a_0 = 1, z = 0$



$$ds^2 = -dt^2 + a^2(t) \left[\frac{dr^2}{1 - K(t)r^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right] \rightarrow dv^2 \text{ (physical)}$$

\downarrow
 (co-moving) du^2

$dv = a du$

here, $du = 1$, but dv increase w.r.t. time

physical meaning of FRWL metric:

For a **co-moving observer**, on the **large** scale, the universe is **homogenous** and **isotropic**.

the only metric compatible with the cosmological principle!

e.g. flat case $ds^2 = -dt^2 + a^2(t)[dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2] \rightarrow$ isotropic (rotation symm.)

coordinate transformation $\rightarrow ds^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j \rightarrow$ homo (spatial shift symm.)

$$[ds^2 = g_{\mu\nu}dx^\mu dx^\nu = \tilde{g}_{\mu\nu}d\tilde{x}^\mu d\tilde{x}^\nu]$$

iso/hom is purely geometry property of the space-time, so it does **NOT** depends on the coord. But, some properties are more easily demonstrated in some specific coordinates.

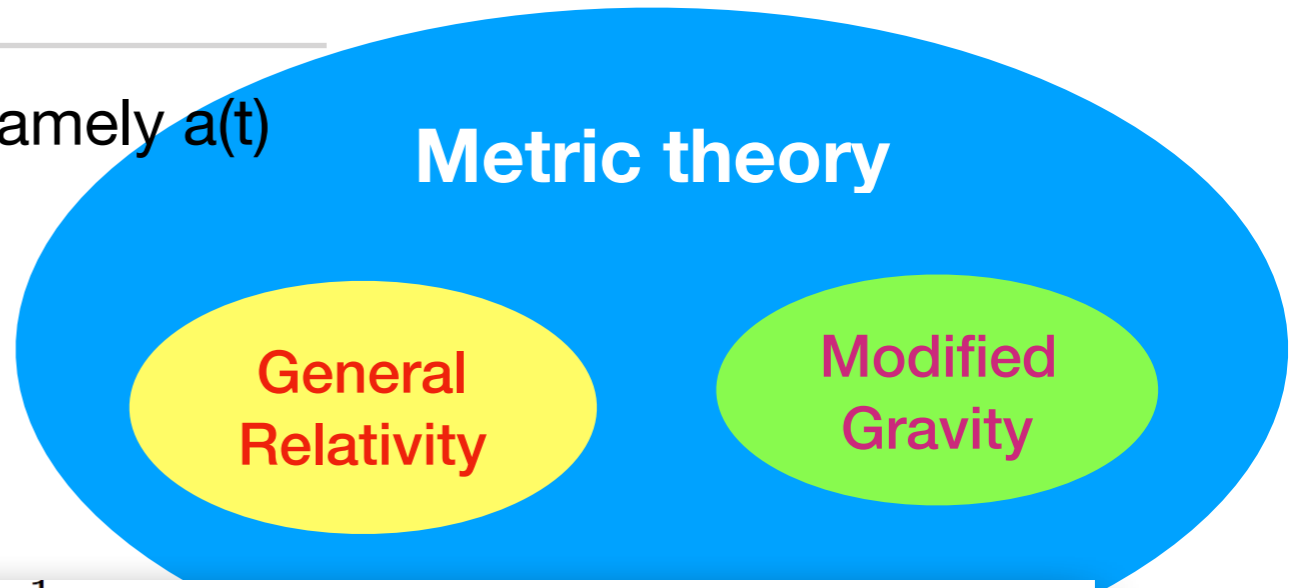
More importantly, FRWL metric defines **a unique clock** with time coordinate 't'

All the observers, **who satisfy the CP**, have to **co-move with this clock**

Now, the metric is fixed up to a function of 't', namely a(t)

This is because, up to now, we only use the geometric/symm. property of the space-time

In order to fix a(t), we need to solve the dynamical equation of gravity sector



e.g. GR: $G_{\mu\nu} = 8\pi G T_{\mu\nu}$ f(R) gravity: $F(R)R_{\mu\nu}(g) - \frac{1}{2}f(R)g_{\mu\nu} - \nabla_\mu \nabla_\nu F(R) + g_{\mu\nu} \square F(R) = \kappa^2 T_{\mu\nu}^{(M)}$

3. Friedmann eq.



John Wheeler

Spacetime tells matter how to move; matter tells spacetime how to curve

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

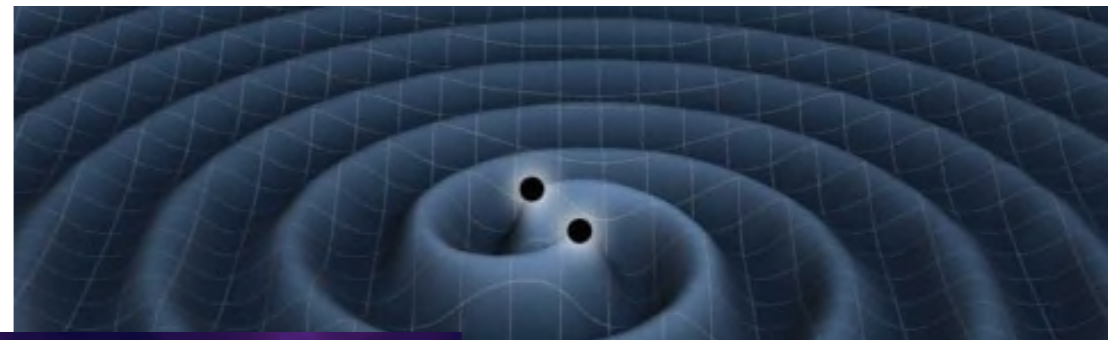
gravity sector
(Geometry of space-time)

gravitational
coupling constant

matter sector
(stress-energy tensor)

(EoM of gravitational d.o.f.)

Einstein eq.



4-velocity relative
to the obs

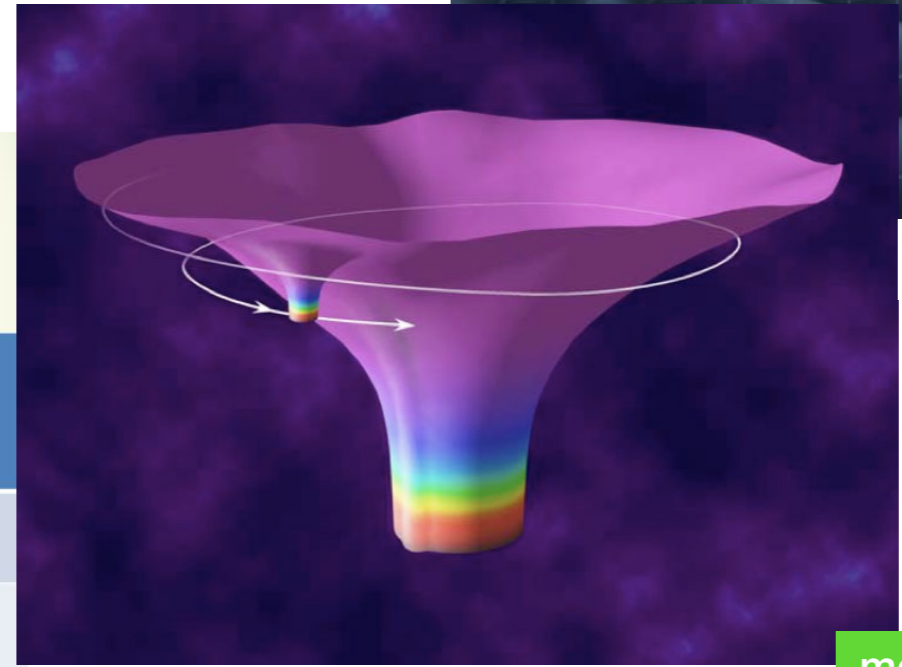
$$T_{\mu\nu} = (\rho + P)U_{\mu}U_{\nu} - P g_{\mu\nu}$$

$$\begin{pmatrix} -\rho & 0 & 0 & 0 \\ 0 & +P & 0 & 0 \\ 0 & 0 & +P & 0 \\ 0 & 0 & 0 & +P \end{pmatrix}$$

model cosmic matter distribution
by the fluid approach!

Gravitation and the other fundamental interactions

Fundamental Interaction	Crucial years	Fundamental constant	Normalized Intensity
Gravity	1687	$Gm_p^2/\hbar c$	5.1×10^{-39}
Weak nuclear force	1934	$G_{Fermi} (m_p c^2)^2$	1.03×10^{-5}
Electromagnetism	1864	$e^2/(4\pi\epsilon_0\hbar c)$	$7.3 \times 10^{-3} \simeq 1/137$
Strong nuclear force	1935/1947	α_s	0.119



need a lots of energy to bend the space-time!

$$\nabla_{\mu} T^{\mu\nu} = 0$$

energy-momentum conservation eq.
(EoM of matter)

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}R g_{\mu\nu} \quad \text{Einstein tensor}$$

$$R_{\mu\nu} \equiv \partial_\lambda \Gamma_{\mu\nu}^\lambda - \partial_\nu \Gamma_{\mu\lambda}^\lambda + \Gamma_{\lambda\rho}^\lambda \Gamma_{\mu\nu}^\rho - \Gamma_{\mu\lambda}^\rho \Gamma_{\nu\rho}^\lambda \quad \text{Ricci tensor}$$

$$R = R^\mu{}_\mu = g^{\mu\nu} R_{\mu\nu} \quad \text{Ricci scalar}$$

$$\Gamma_{\alpha\beta}^\mu \equiv \frac{1}{2}g^{\mu\lambda}(\partial_\alpha g_{\beta\lambda} + \partial_\beta g_{\alpha\lambda} - \partial_\lambda g_{\alpha\beta}) \quad \text{connection/Christoffel symbol}$$

$G_{\mu\nu}$ contains $(g_{\mu\nu}, \dot{g}_{\mu\nu}, \ddot{g}_{\mu\nu}) \longrightarrow$ good! does not need acceleration of acceleration

Classical dynamics tell us: a canonical dynamical system, shall at most contain the 2nd order time derivative of its dynamical variables.

However, $G_{\mu\nu}(g_{\mu\nu})$ is a **non-linear** functional. \longrightarrow bad! very hard to solve

e.g. for merger stage of binary black hole system, EE is very very hard to solve!

EE is written @1915, but the first bbh solution is got @2005

For FRWL metric:

$$ds^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j$$

[Pb1.] $R_{00} = -3\frac{\ddot{a}}{a}$

$$R_{ij} = [a\ddot{a} + 2\dot{a}^2]\delta_{ij}$$

$$G^0_0 = 3\left[\left(\frac{\dot{a}}{a}\right)^2 + \frac{k}{a^2}\right],$$

$$G^i_j = \left[2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2 + \frac{k}{a^2}\right]\delta^i_j$$

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2},$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P)$$

1st Friedmann eq.

2nd Friedmann eq.

We need:

For a co-moving obs: $U^\mu = (-1, 0, 0, 0)$

For a perfect fluid: $T^\mu_\nu = \begin{pmatrix} -\rho & 0 & 0 & 0 \\ 0 & +P & 0 & 0 \\ 0 & 0 & +P & 0 \\ 0 & 0 & 0 & +P \end{pmatrix}$

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2},$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P)$$

Not independent with each other

[Pb2.] Check the relationship between 1st & 2nd Friedmann eq.

$$\nabla_\mu T^{\mu\nu} = 0 \longrightarrow \dot{\rho} + 3\frac{\dot{a}}{a}(\rho + P) = 0$$

Besides the conservation eq. we also need the thermal dynamical info of the fluid

$$\frac{H^2}{H_0^2} = \Omega_r a^{-4} + \Omega_m a^{-3} + \Omega_k a^{-2} + \Omega_\Lambda$$

(CAN SNela measure H_0 ?)

e.g. Equation of State $w = P / \rho$

$$H^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} \quad V \sim a^{-3}$$

- for non-relativistic particle, E is conserved
- for relativistic photon, E is **NOT** conserved!

$$\rho \propto a^{-3(1+w)},$$

$$\rho \propto \begin{cases} a^{-3} & \text{matter} \\ a^{-4} & \text{radiation} \\ a^0 & \text{vacuum} \end{cases}$$

$$\rho_{\text{crit},0} = \frac{3H_0^2}{8\pi G} = 1.9 \times 10^{-29} h^2 \text{ grams cm}^{-3}$$

$$= 2.8 \times 10^{11} h^2 M_\odot \text{ Mpc}^{-3}$$

$$= 1.1 \times 10^{-5} h^2 \text{ protons cm}^{-3}$$

$$\Omega_{I,0} \equiv \frac{\rho_{I,0}}{\rho_{\text{crit},0}}$$

$$E(v_a) > E(v_b)$$

- for vacuum energy, E is **NOT** conserved!

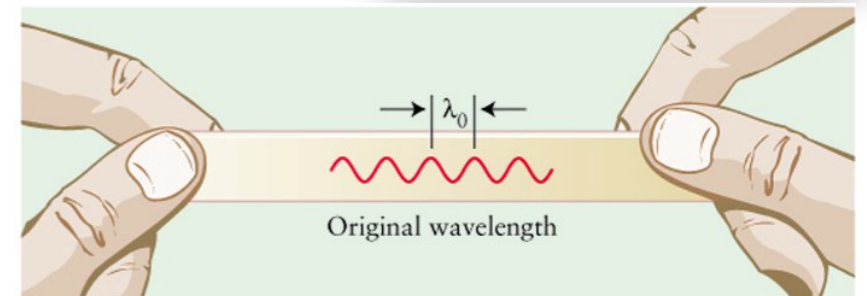
$$dU = -PdV$$

vacuum energy:

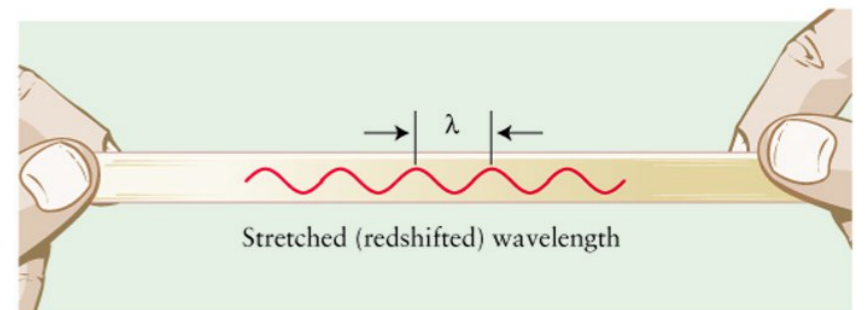
where there is space, there it is

$$dU = \rho dV$$

A negative EoS means, after a system work to the environment, its internal energy is **increased** instead of decreased.



(a) A wave drawn on a rubber band ...



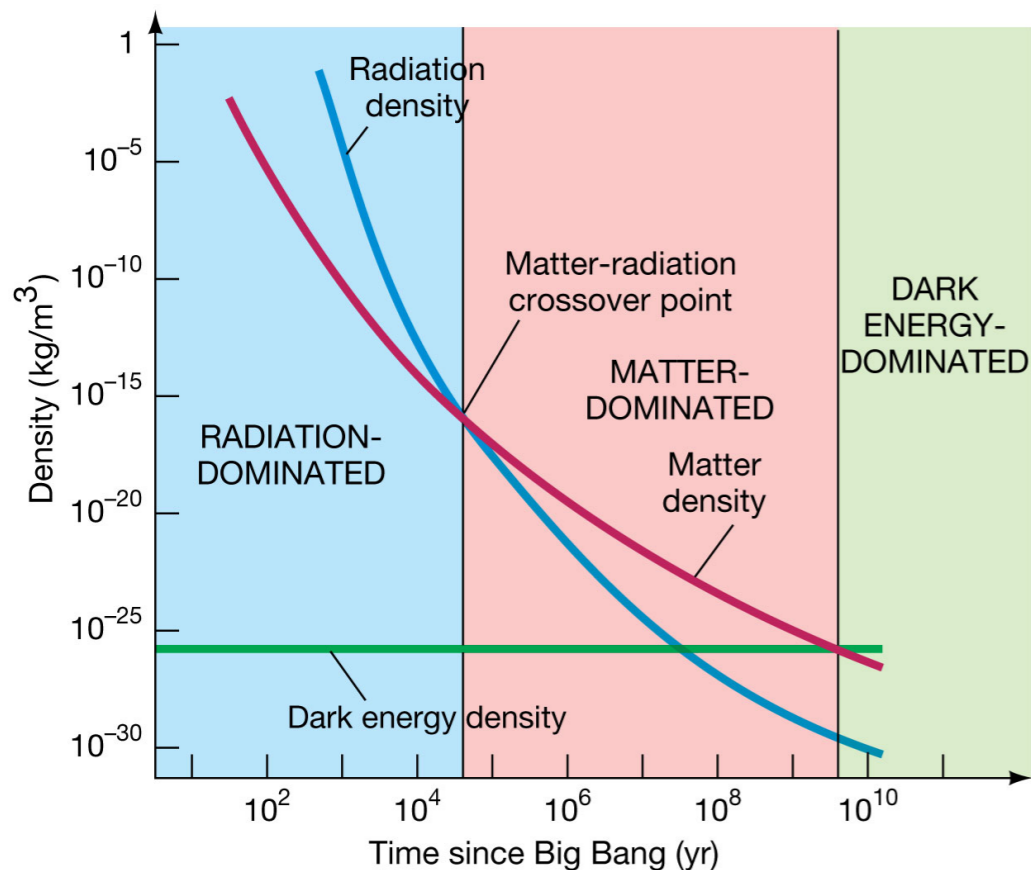
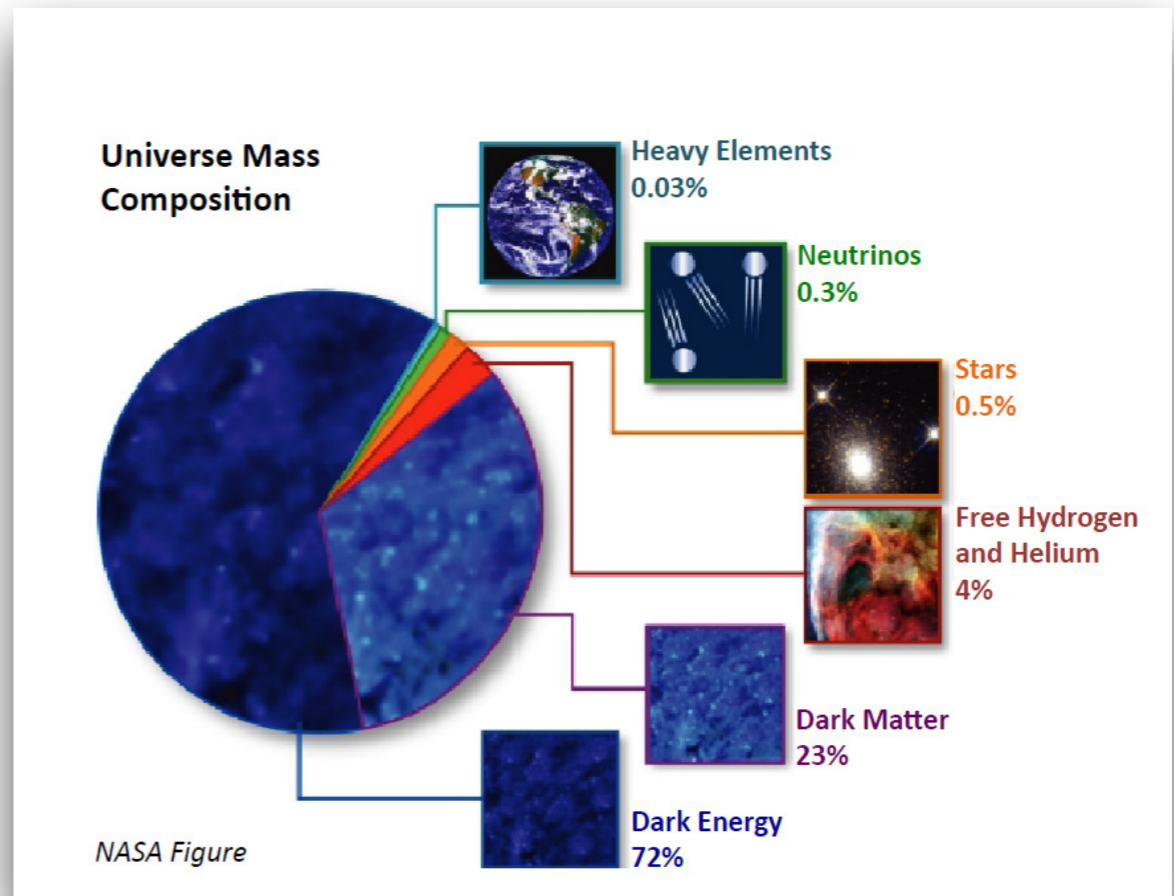
(b) ... increases in wavelength as the rubber band is stretched.

single component universe solution

$$a(t) \propto \begin{cases} t^{2/3(1+w_I)} & w_I \neq -1 \\ e^{Ht} & w_I = -1 \end{cases} \begin{matrix} t^{2/3} & \text{MD} \\ t^{1/2} & \text{RD} \end{matrix}$$

$$a(\tau) \propto \begin{cases} \tau^{2/(1+3w_I)} & w_I \neq -1 \\ (-\tau)^{-1} & w_I = -1 \end{cases} \begin{matrix} \tau^2 & \text{MD} \\ \tau & \text{RD} \\ \Lambda\text{D} & \end{matrix}$$

matter ingredient



baryon & DM is **indistinguishable** on the large scale

on the small scale, baryon stop collapsing once below its jeans radius
DM will keep collapsing until $r \sim 0$

photon & neutrino is **indistinguishable** in the early stage ($z > 200$)
once $z < 200$, neutrino will becomes non-relativistic, behaves more like DM

As of DE:



4. Distance

$$ds^2 = -dt^2 + a^2(t)[dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2] \quad \text{Null-like geodesic} \longrightarrow ds = 0$$

a light co-movingly propagate with background expansion along the radial direction from z_1 to z_0 $\longrightarrow dr = dt / a$

- **co-moving distance [along line-of-sight]:**

$$\chi = \int_{z_1}^{z_0} dr = \int_{z_1}^{z_0} \frac{dt}{a} = \int_{z_1}^{z_0} \frac{dz}{H(z)}$$

- **diameter distance [transverse]:**

$r^* \longrightarrow$ known by prior (physical scale)

measure the angular separation θ

assuming Euclidean geometry, we can define

But, this is WRONG! Physical geometry is NOT Euclidean, the co-moving one does! We need re-scale r^* to the co-moving one, namely r^*/a_1

$$D_A = \frac{\chi}{1+z}$$

- **luminosity distance:**

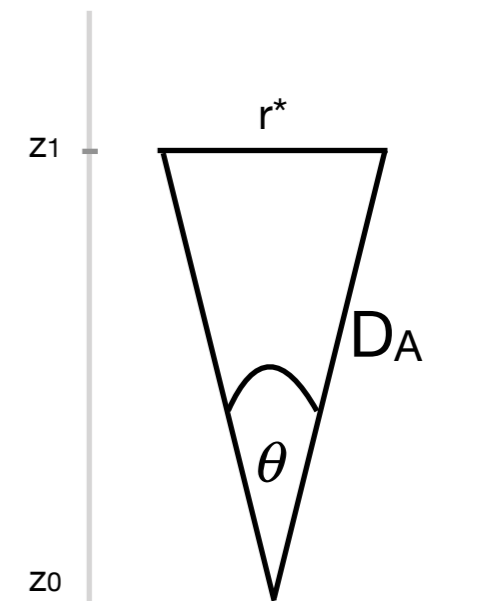
$$F_{obs}(z_0) = \frac{L_{ABS}(z_1)}{4\pi * D_L^2}$$

In Euclidean geometry, we shall have

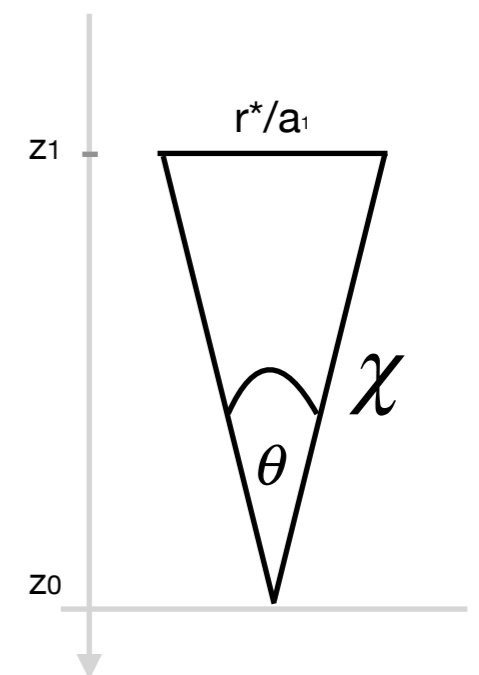
$$F_{obs}(z_0) = \frac{L_{ABS}(z_0)}{4\pi * \chi^2}$$

$$L_{ABS}(z_0) = \frac{E_0}{\delta t_0} = \frac{E_1 / (1+z)}{(1+z)\delta t_1} = L_{ABS}(z_1) / (1+z)^2$$

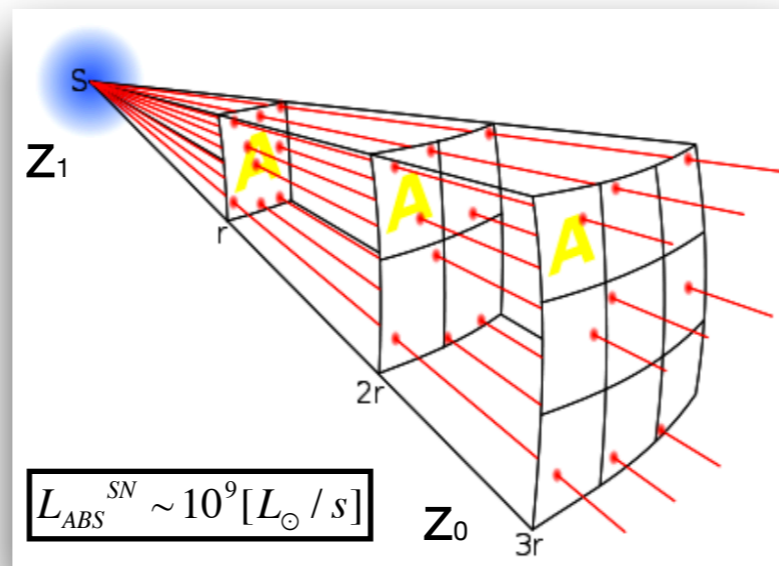
$$D_L = (1+z)\chi$$



$$D_A = \frac{r^*}{\theta}$$



$$\chi = \frac{(1+z_1)r^*}{\theta}$$



为什么宇宙的年龄是 130 亿年，而我们却能看到 470 亿光年远的东西？

2017-09-04 土豆泥 超级数学建模



麻烦来个通俗解释

宇宙的年龄约130亿年，可观宇宙半径约为470亿光年。资料在维基百科等很多地方都可以查到。我只是不理解如果光速不可超越的话，怎么会在130亿年时间里产生了470亿光年的距离，并且是人类可观测到的，也就

[Pb3.]

calculate this number now!

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$$\Omega_{m,0} = 0.3$$

$$\Omega_{\Lambda,0} = 0.7$$

$$\Omega_{r,0} = 10^{-5}$$

$$H_0 = 68[km / s / Mpc]$$

$$c = 30 * 10^4 km / s$$

$$z_0 = 0; z_1 = 1100$$

$$\chi = \frac{c}{H_0} \int_{z_1}^{z_0} \frac{dz}{E(z)}; E^2(z) = \Omega_{m,0}(1+z)^3 + \Omega_{r,0}(1+z)^4 + \Omega_{\Lambda,0}$$

~14 Gpc

$$t = \frac{c}{H_0} \int_{z_1}^{z_0} \frac{dz}{E(z)(1+z)}$$

~138 billion yr

5. Standard candle

$$F_{obs}(z_0) = \frac{L_{ABS}(z_1)}{4\pi * D_L^2}$$

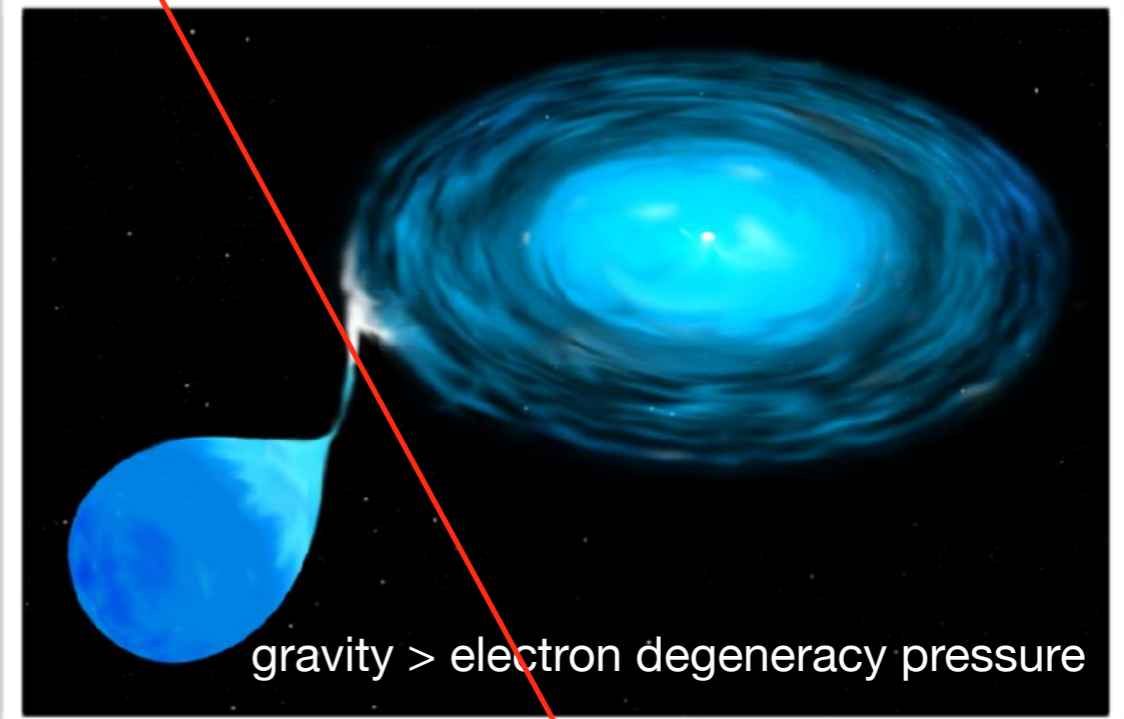
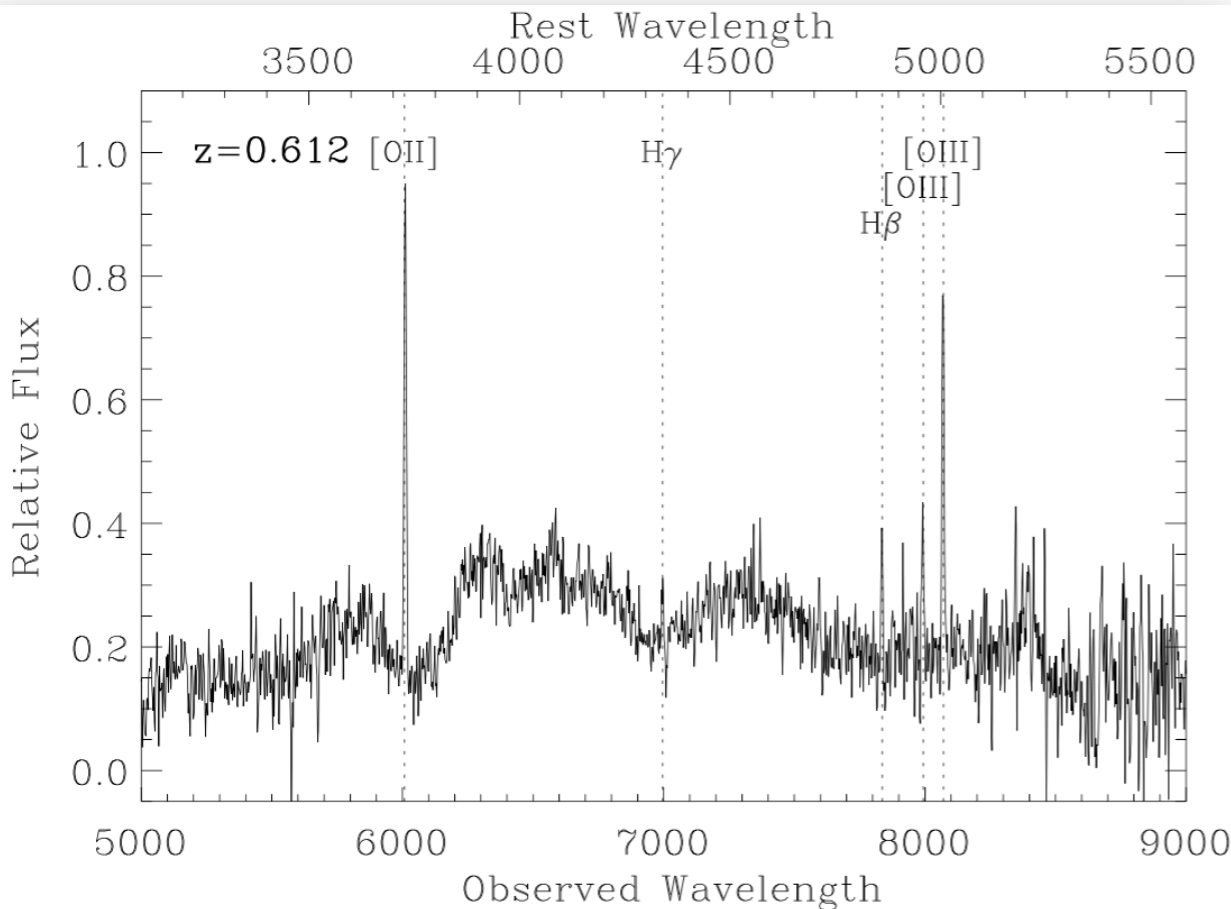
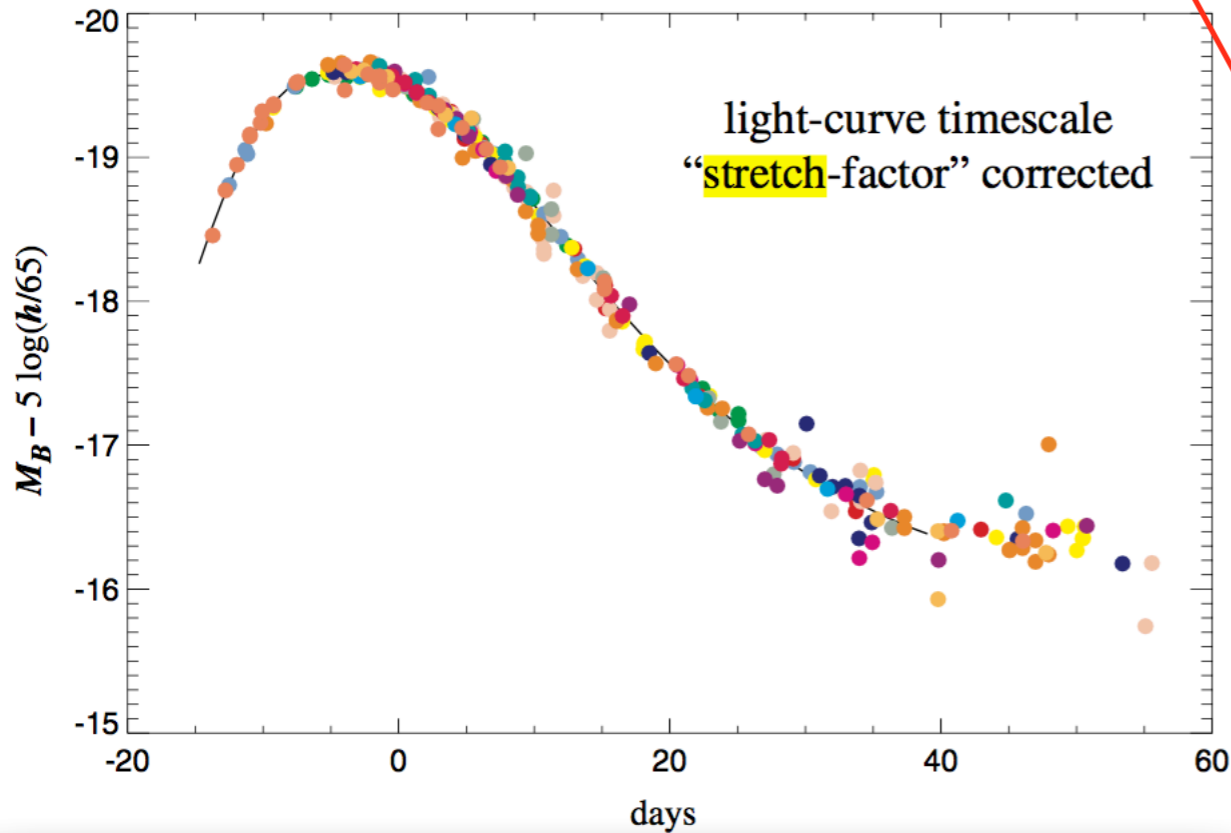
$$L_{ABS}^{SN} \sim 10^9 [L_{\odot} / s]$$

$$M_{WD} > 1.44 M_{\odot}$$

(Chandrasekhar limit)



SN Ia explode from a binary star system, typically one white dwarf, one giant star



$$D_L = (1+z)\chi$$

$$\lambda_0 \sim 3700 \text{ \AA}$$

$$\chi = \int_{z_1}^{z_0} dr = \int_{z_1}^{z_0} \frac{dt}{a} = \int_{z_1}^{z_0} \frac{dz}{H(z)}$$

$$z \sim 6000/3700 - 1 \sim 0.6$$

test cosmology

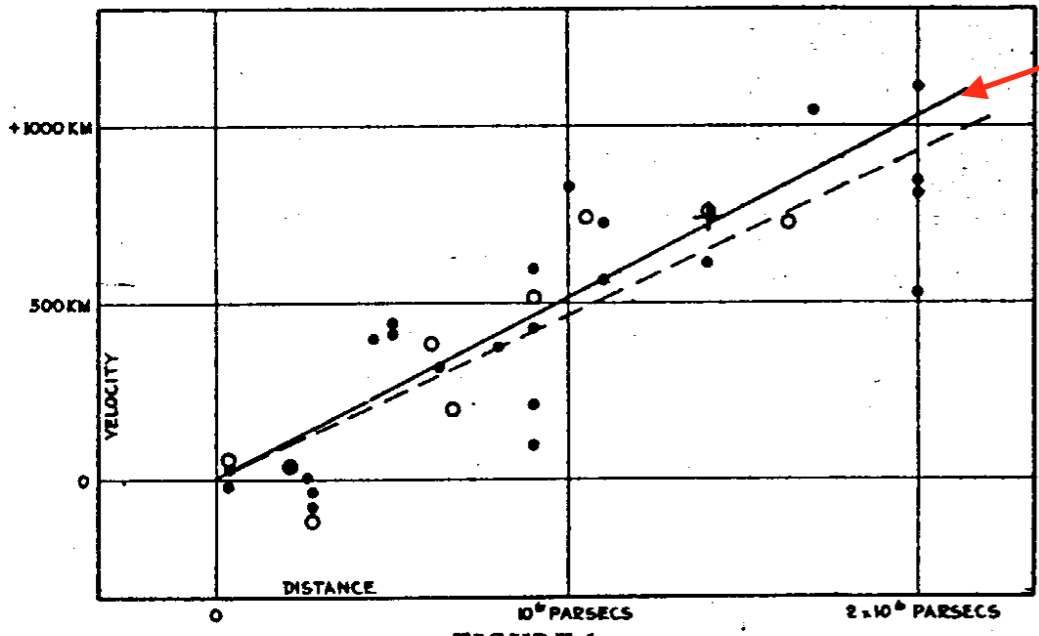
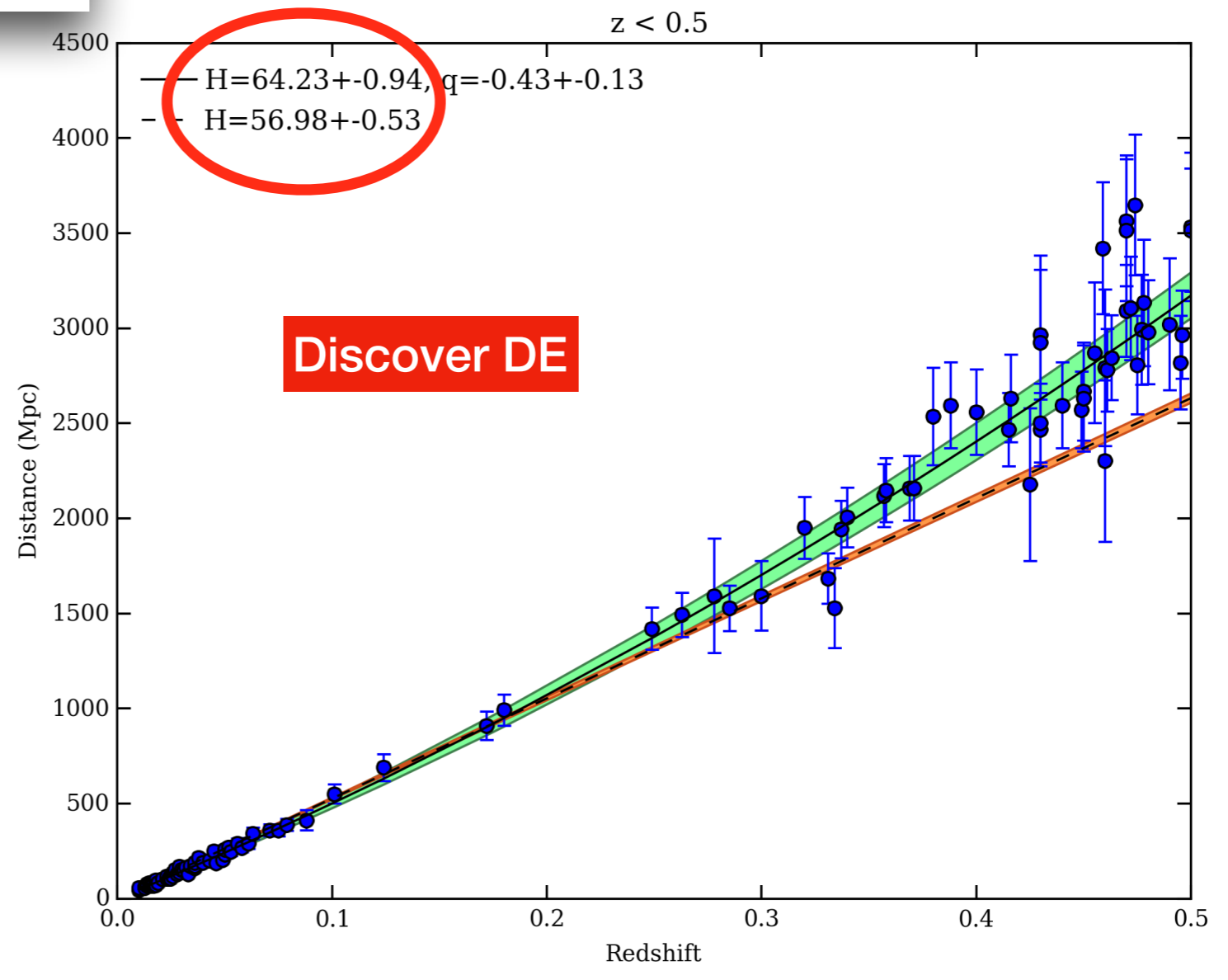


FIGURE 1

$H_0=500$

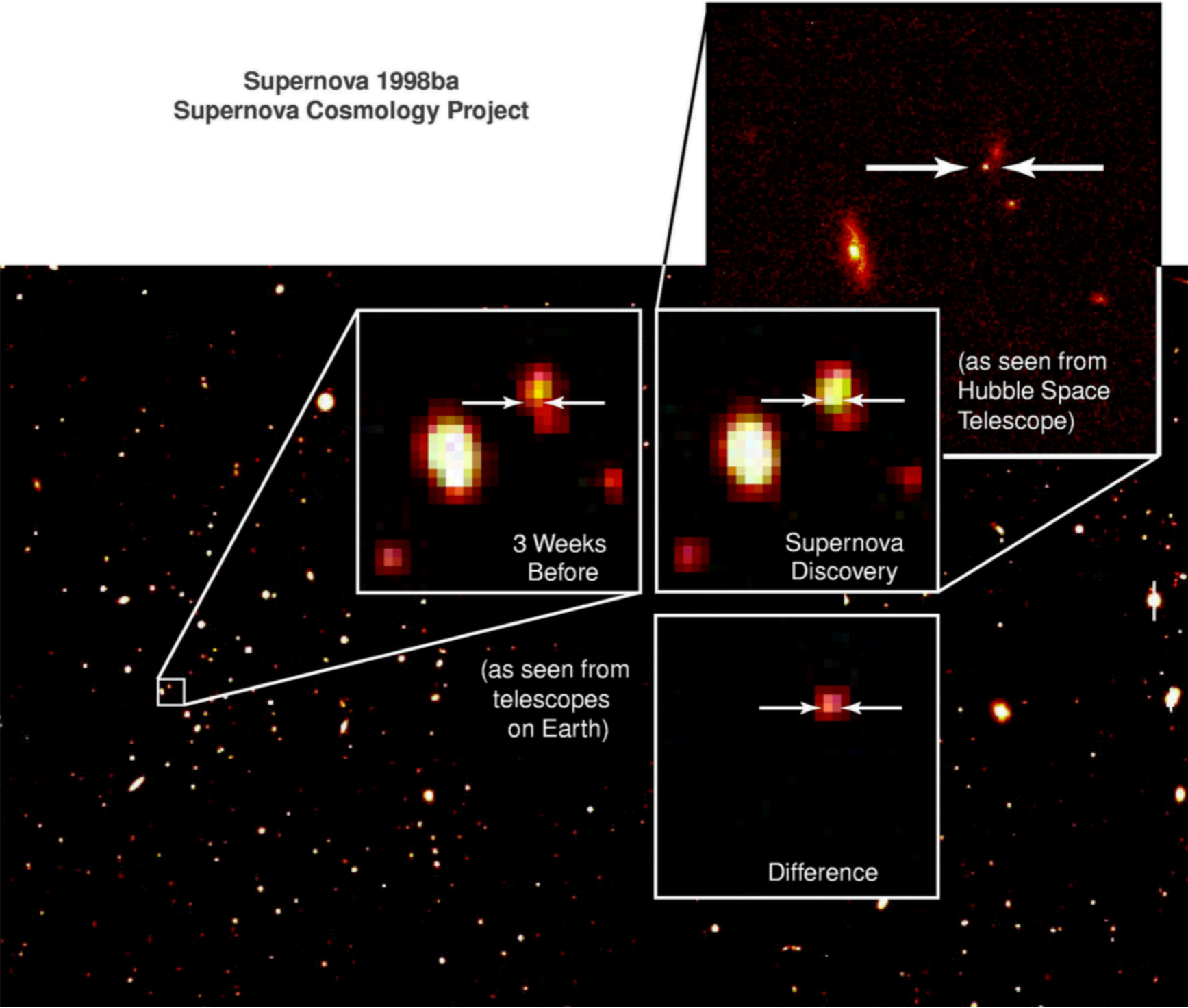
Hubble 1929

Now

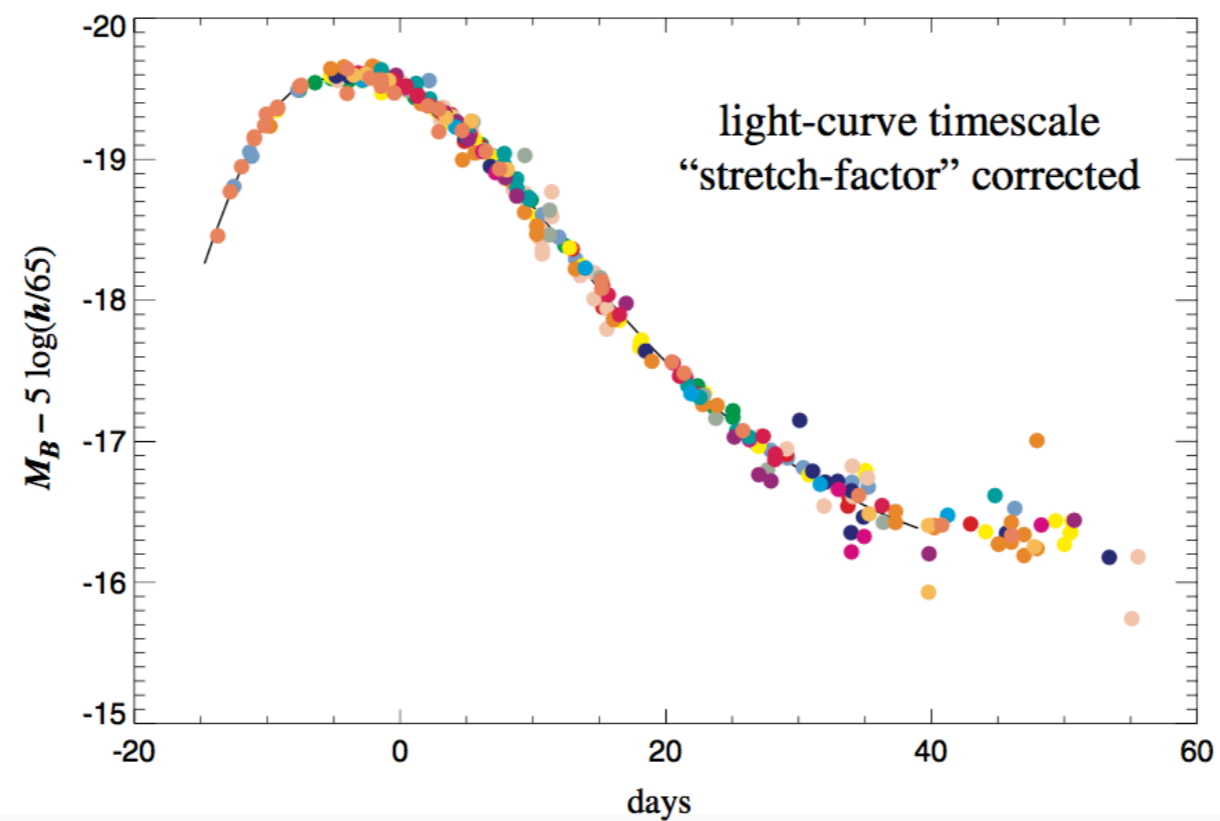
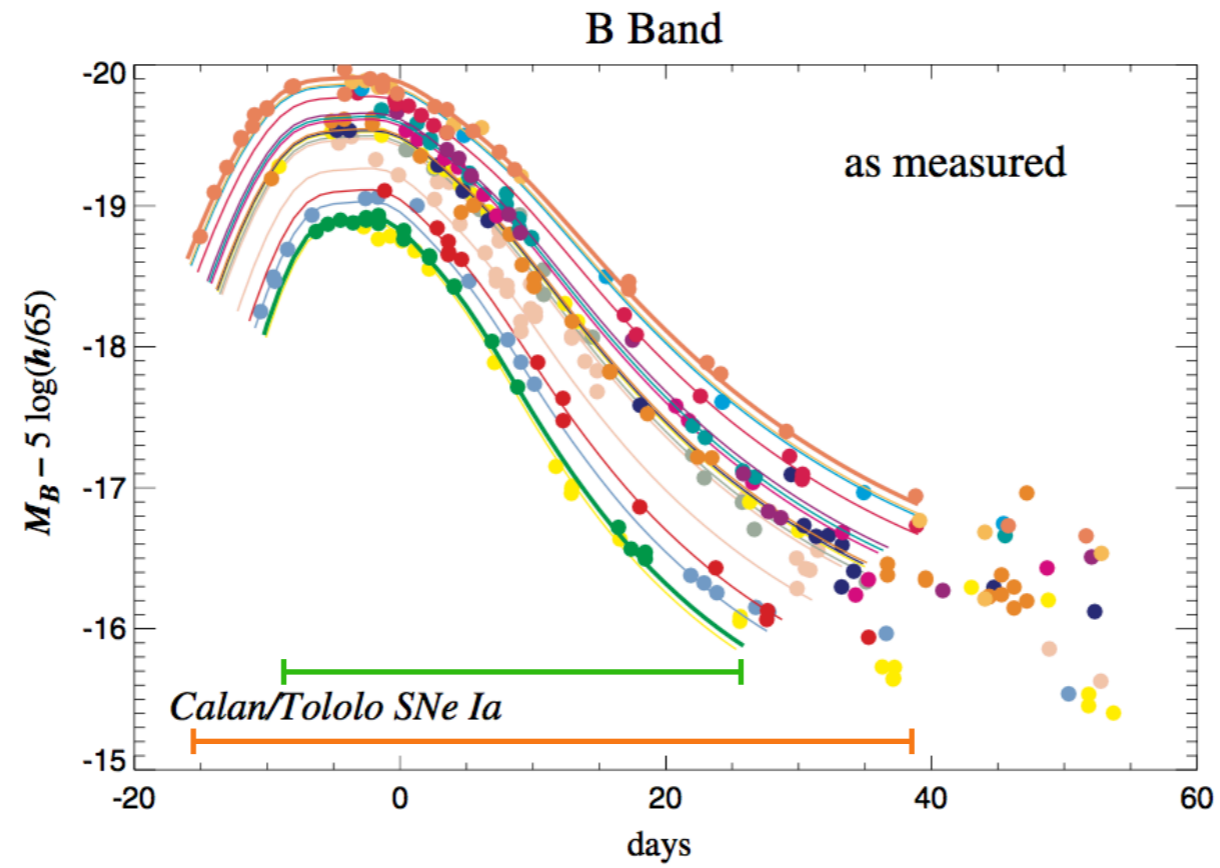


Discover DE

how to discover a SN



SN Ia: standard candle \rightarrow standardised candle

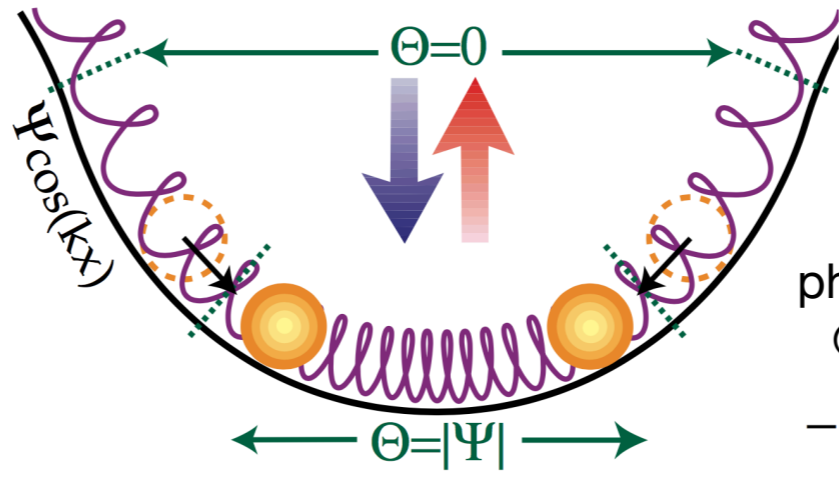
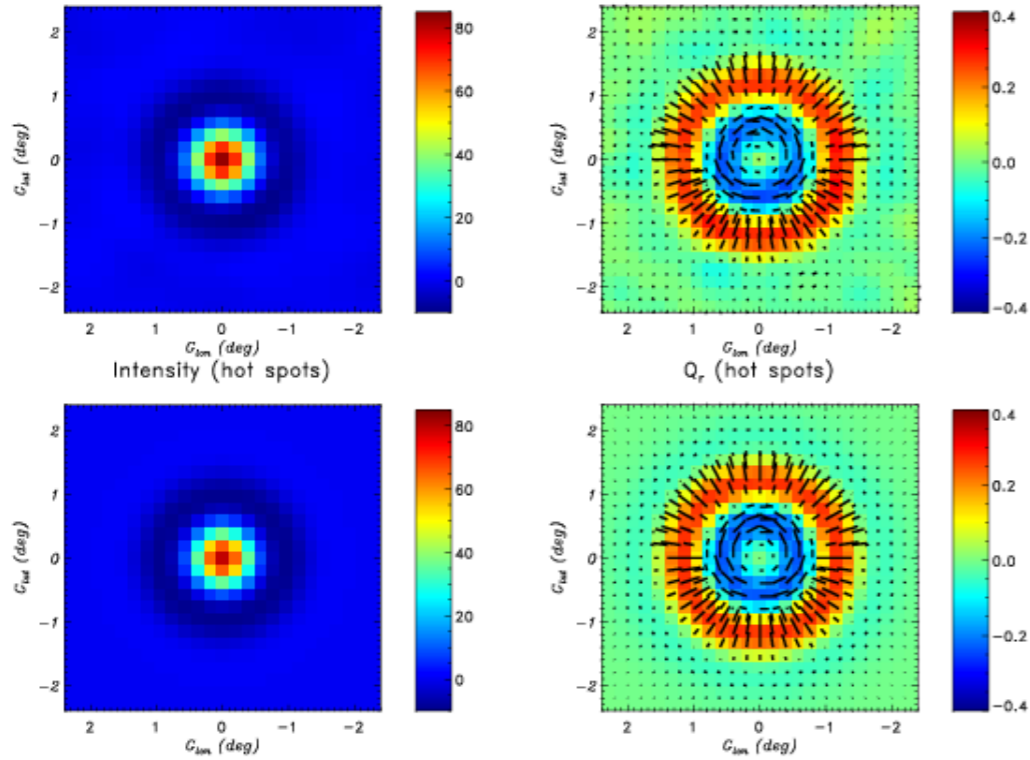
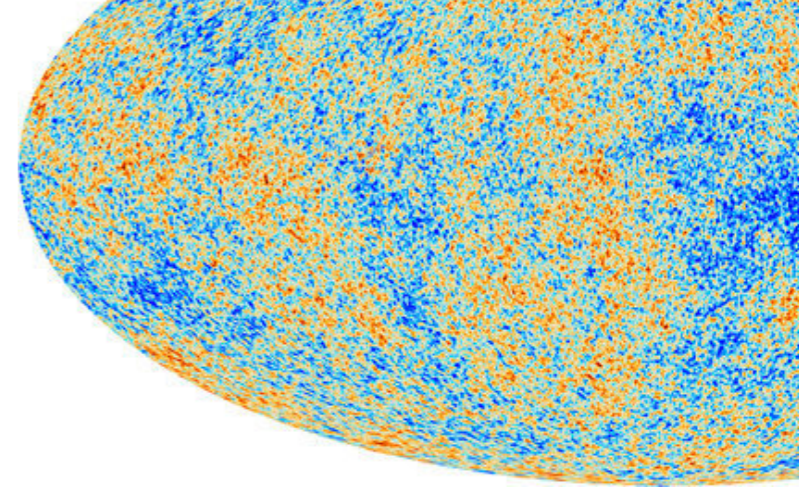


Kim, *et al.* (1997)

time-scale stretch

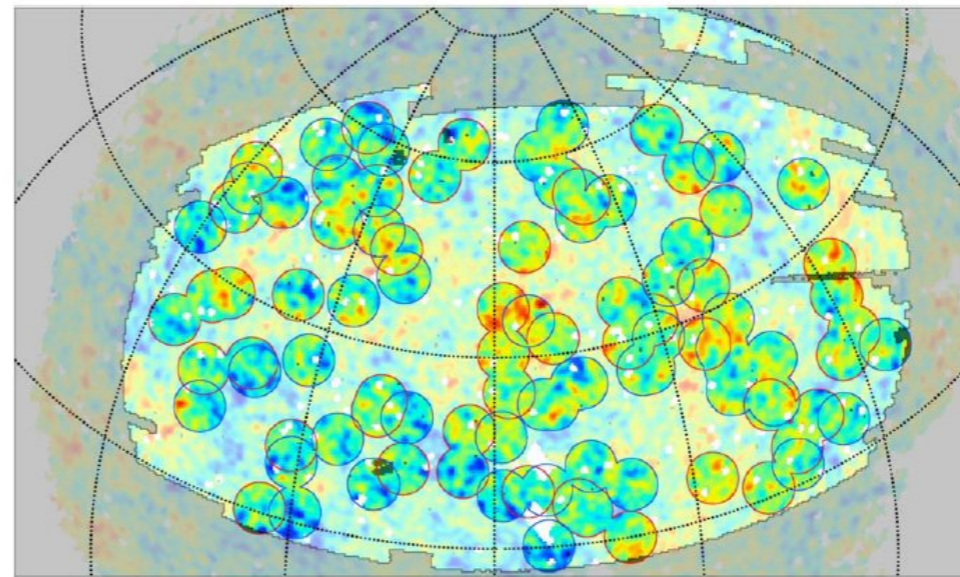
6. standard ruler

→ Baryon Acoustic Oscillation



photon pressure balance gravity
 @ 150Mpc (co-moving scale)
 — physical scale ~ 150 kpc

$$\chi \sim 150 \text{ Mpc} / D_A \sim 150 \text{ kpc}$$



$$D_A = \frac{\chi}{1+z}$$

+

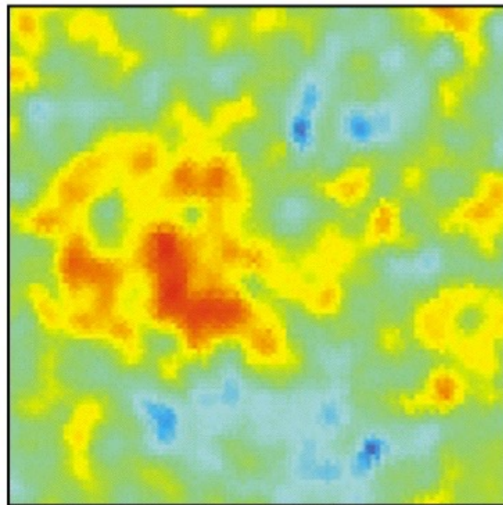
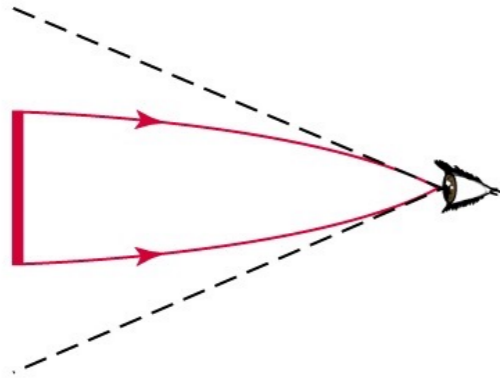
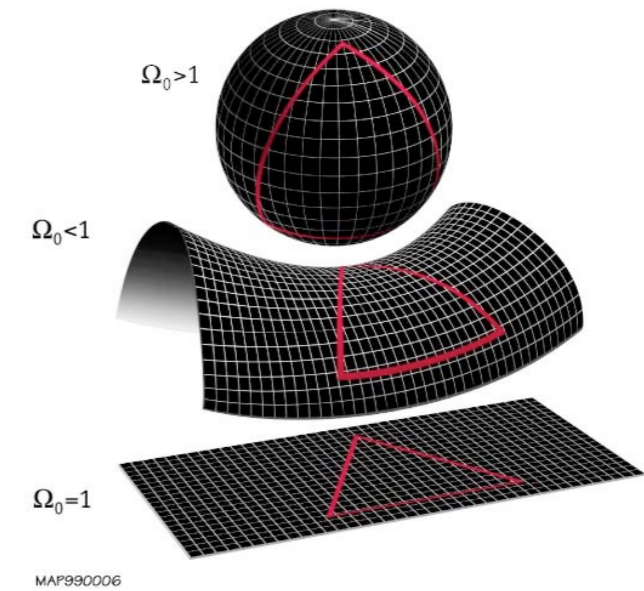
$z \sim 1100$

$$\chi = \int_{z_1}^{z_0} dr = \int_{z_1}^{z_0} \frac{dt}{a} = \int_{z_1}^{z_0} \frac{dz}{H(z)}$$

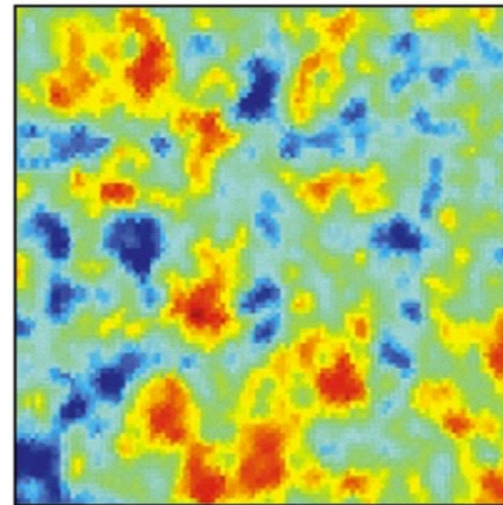
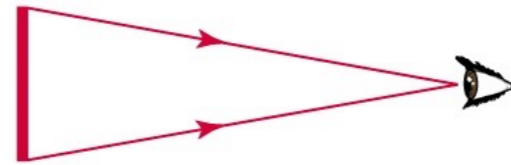
test cosmology

e.g. measure the 3d spatial curvature

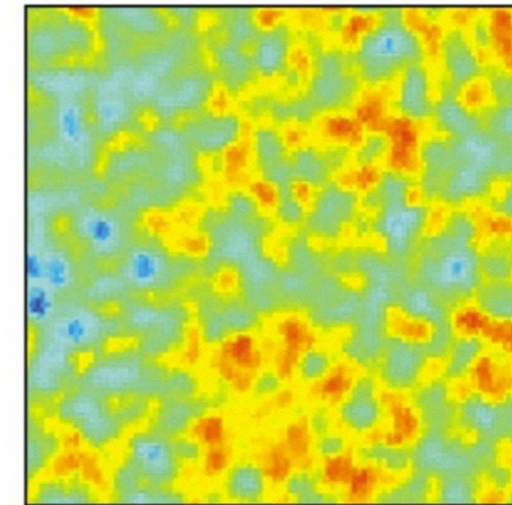
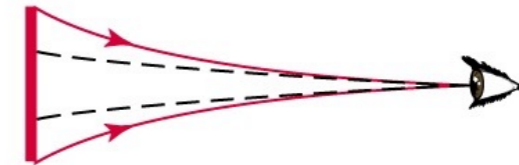
$$ds^2 = -dt^2 + a^2(t) \left[\frac{dr^2}{1 - K(t)r^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right]$$



a If universe is closed, "hot spots" appear larger than actual size

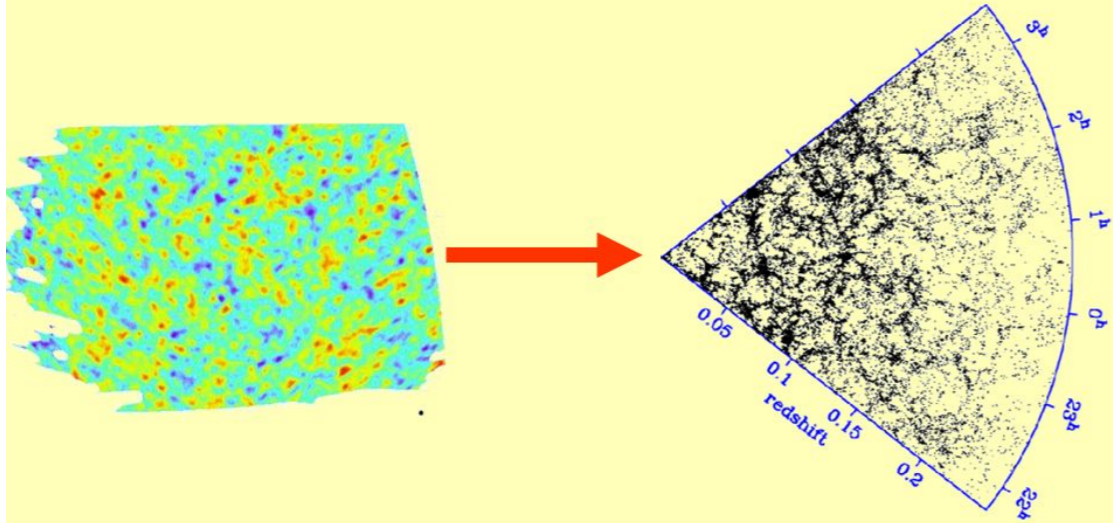


b If universe is flat, "hot spots" appear actual size



c If universe is open, "hot spots" appear smaller than actual size

Measuring large-scale structure in the universe with the 2dF Galaxy Redshift Survey

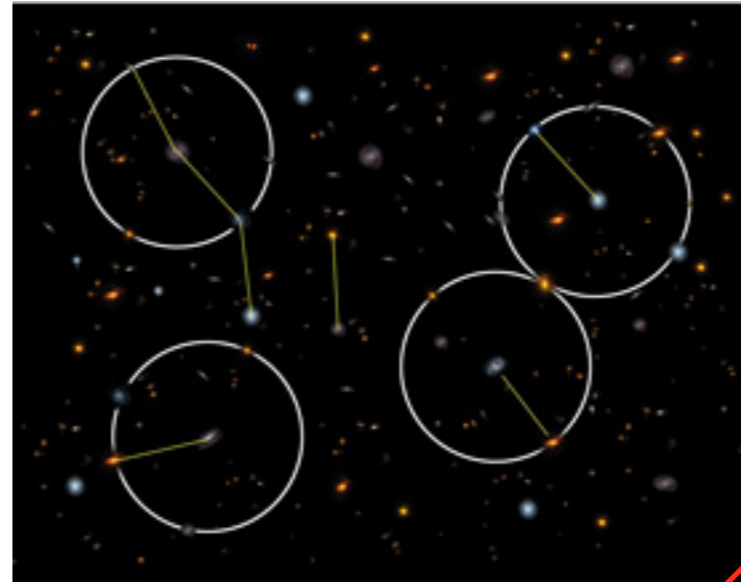


John Peacock

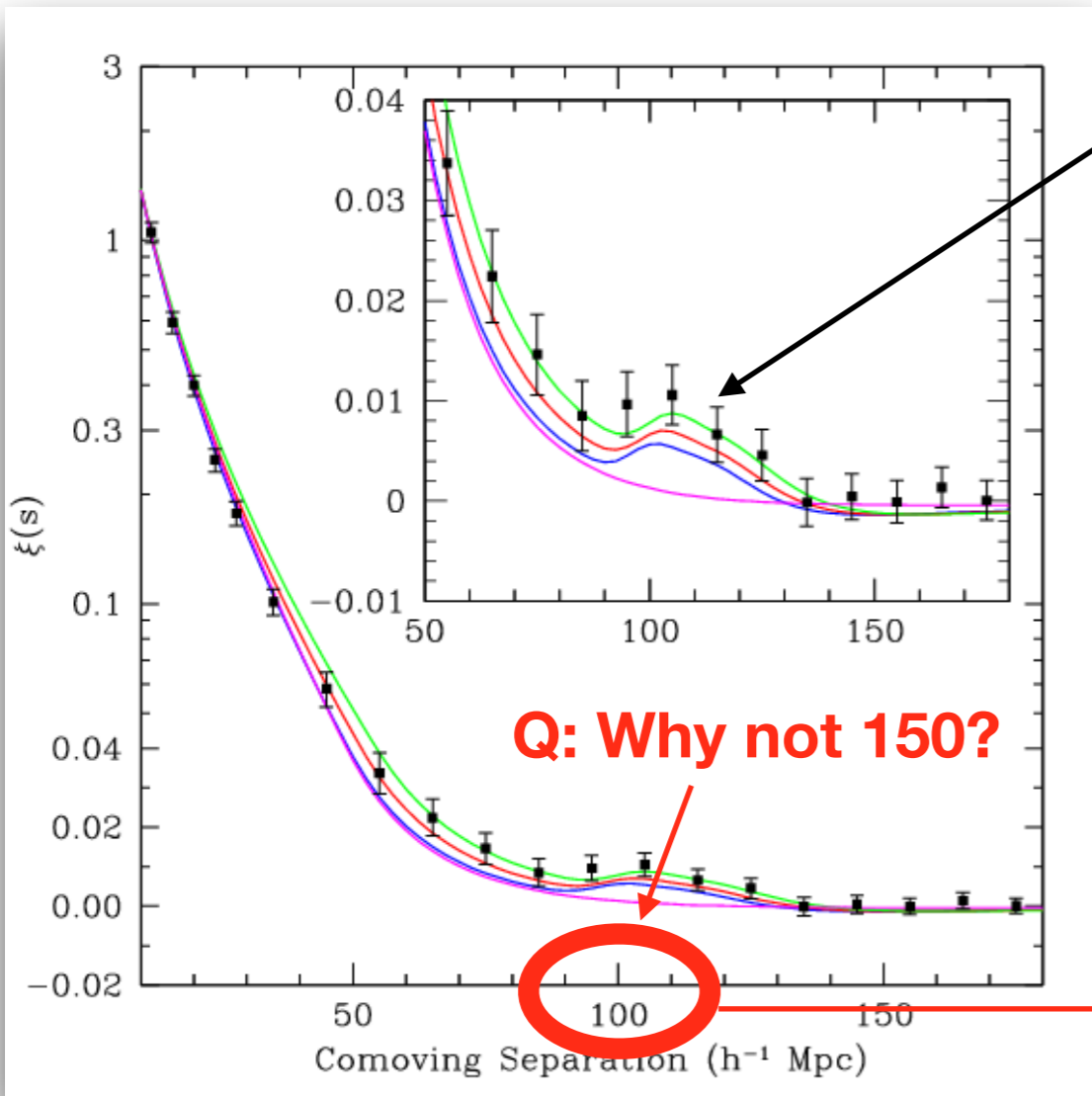
Garching

December 2001

BAO signal can also imprint on the matter distribution, e.g. galaxies



characteristic break

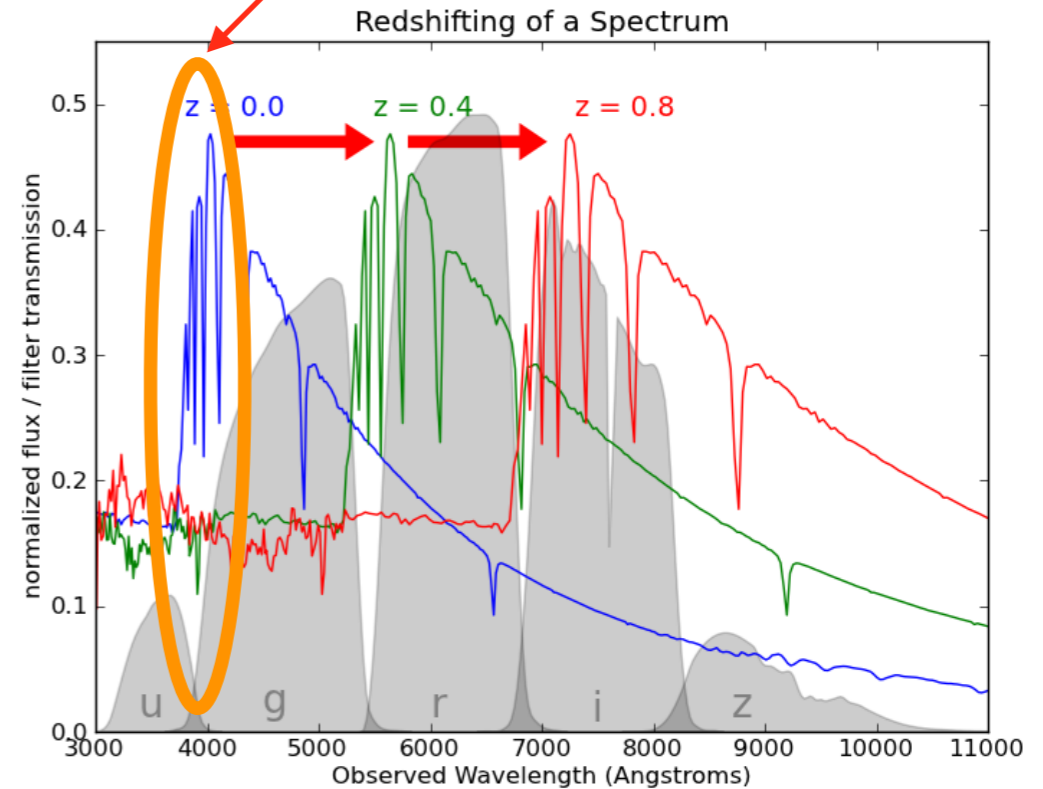


exceed

$$\xi(r) = \frac{DD(r)}{RR(r)} - 1$$

photo-z

(photometric redshift)

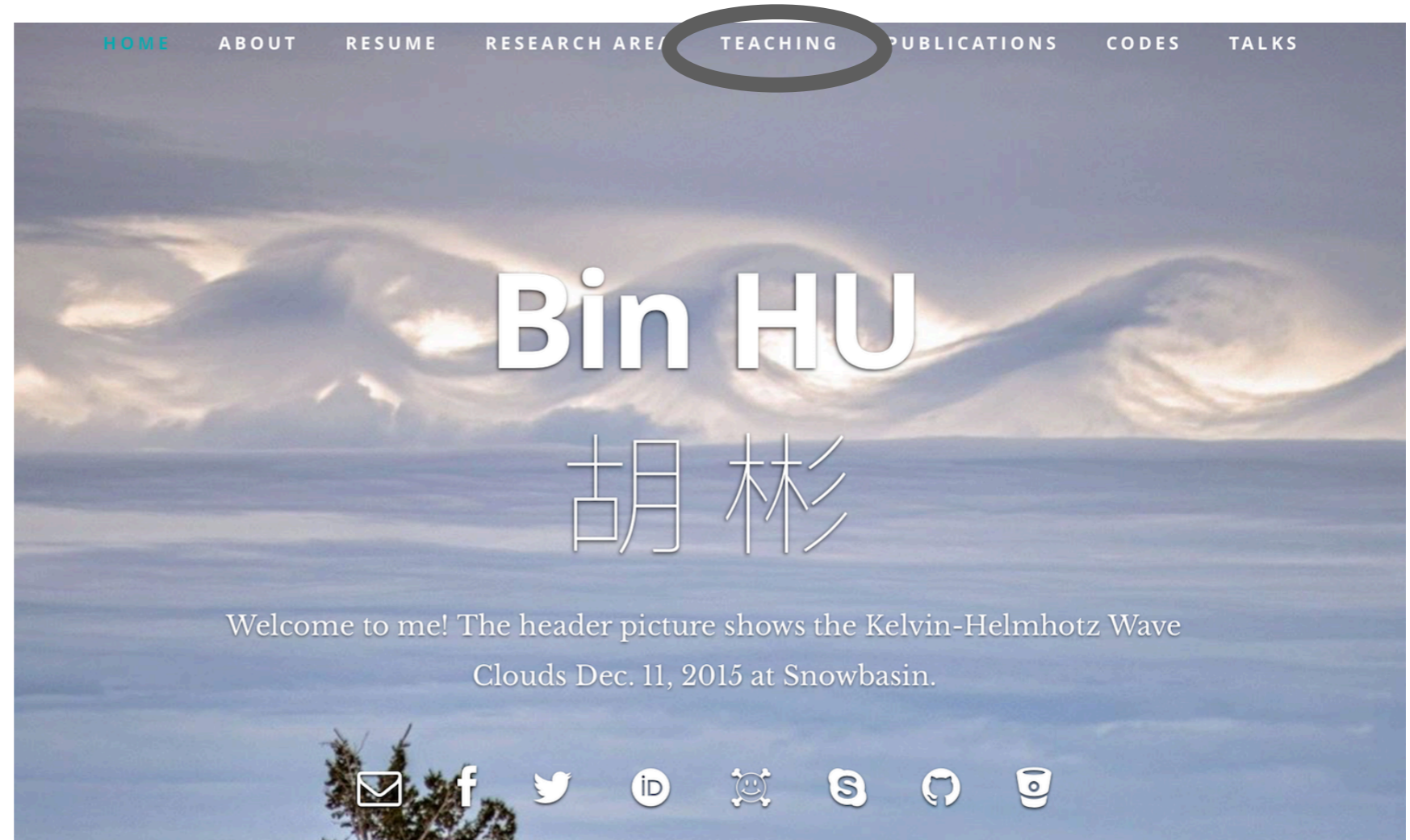


$$\chi = \int_{z_1}^{z_0} dr = \int_{z_1}^{z_0} \frac{dt}{a} = \int_{z_1}^{z_0} \frac{dz}{H(z)}$$

test cosmology

Further reading:

- Baumann Lecture note/Chapter 1
- 宇宙大尺度结构的形成 向守平、冯珑珑/Chapter 1,2,3



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http://astrowww.bnu.edu.cn/sites/hubin/bh_bnu_homepage/#teach