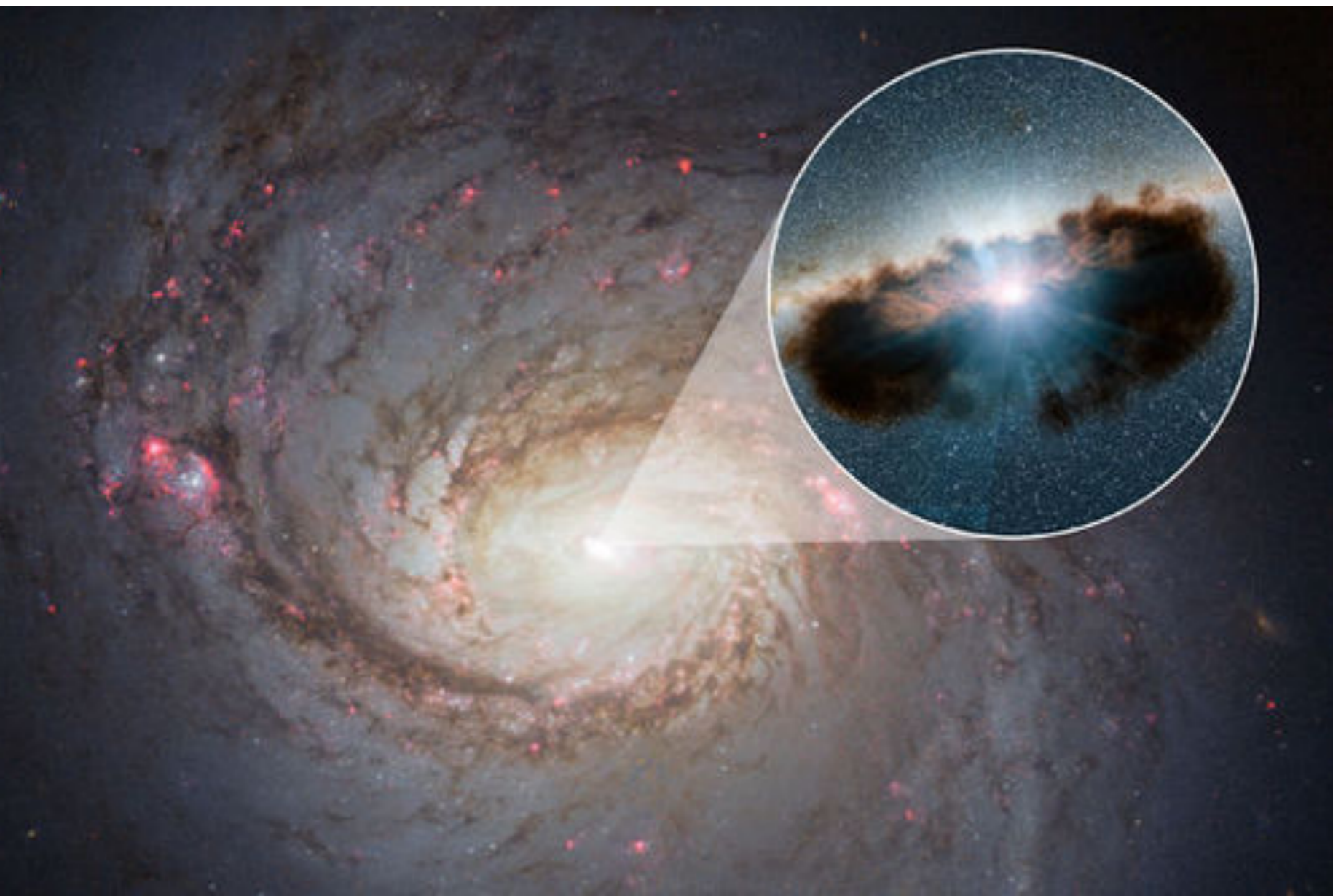


Co-evolution of super massive BHs with galaxies

—stochastic GWB & galaxy clustering



[arXiv:1802.03925]
arguing w. MNRAS

w. Qing Yang@BNU
Xiao-Dong Li@SYSU

Bin HU @ BNU

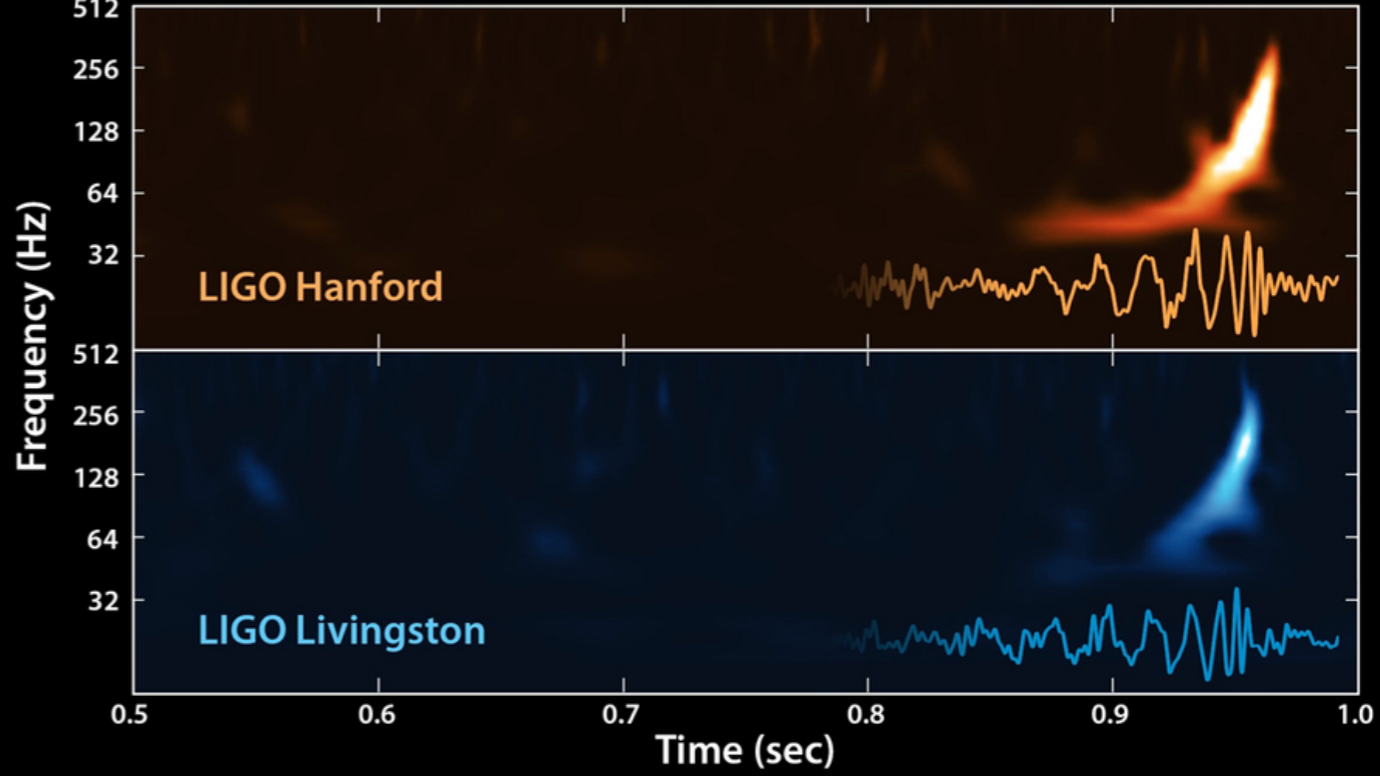
2018/05 Yangzhou



Big Bang



Supermassi



The Gravitational Wave Detectors

age of the universe

years

Wave Period

hours

seconds

milliseconds

10^{-16}

10^{-14}

10^{-12}

10^{-10}

10^{-8}

10^{-6}

10^{-4}

10^{-2}

1

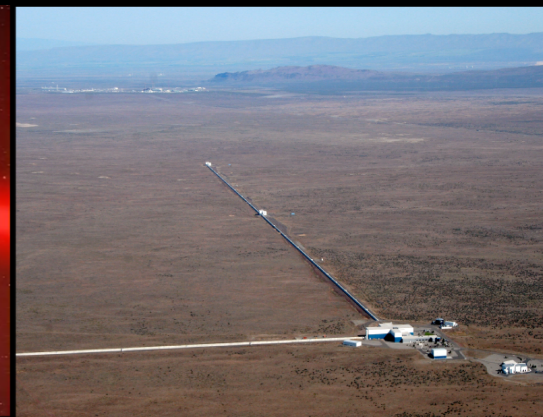
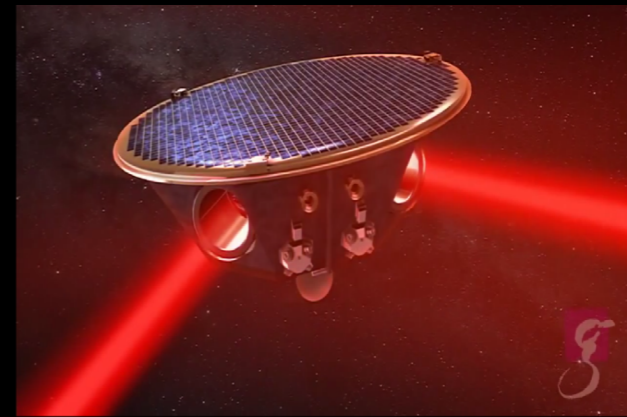
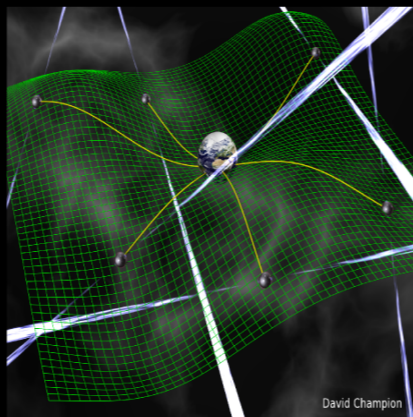
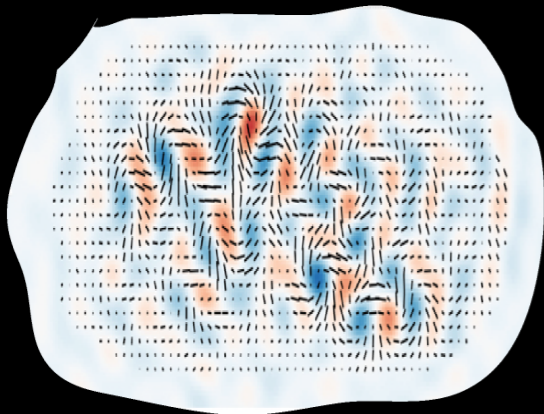
10^2

Wave Frequency

CMB Polarization

Radio Pulsar Timing Arrays

Space-based interferometers **Terrestrial interferometers**



~ kpc (10^{16} km)

$5 \cdot 10^6$ km

4km



李柯伽@北大

如果**2005**年，你问“我们什么时候能探测到随机引力波背景信号？”

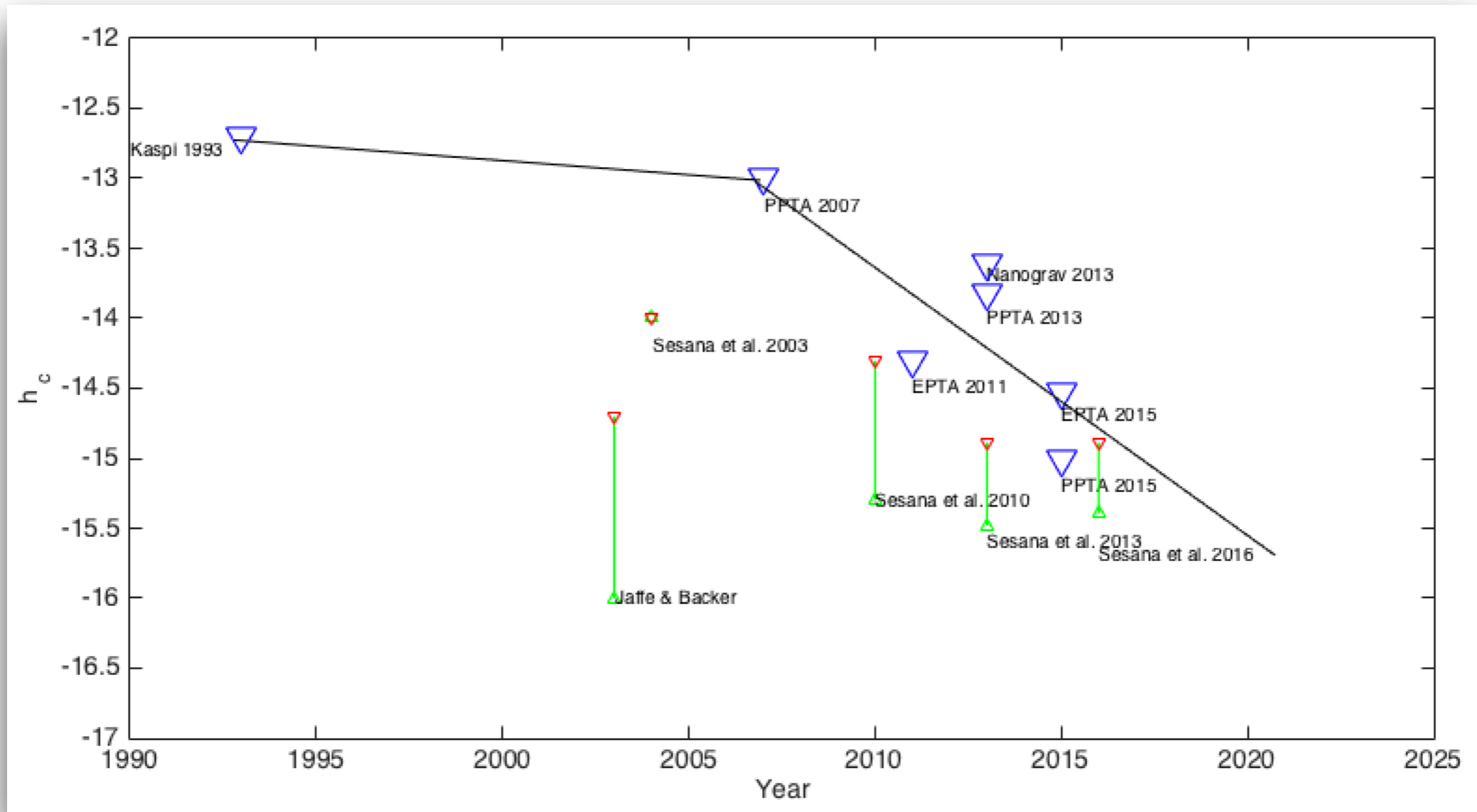
答案：**五年**



李柯伽@北大

如果2010年，你问“我们什么时候能探测到随机引力波背景信号？”

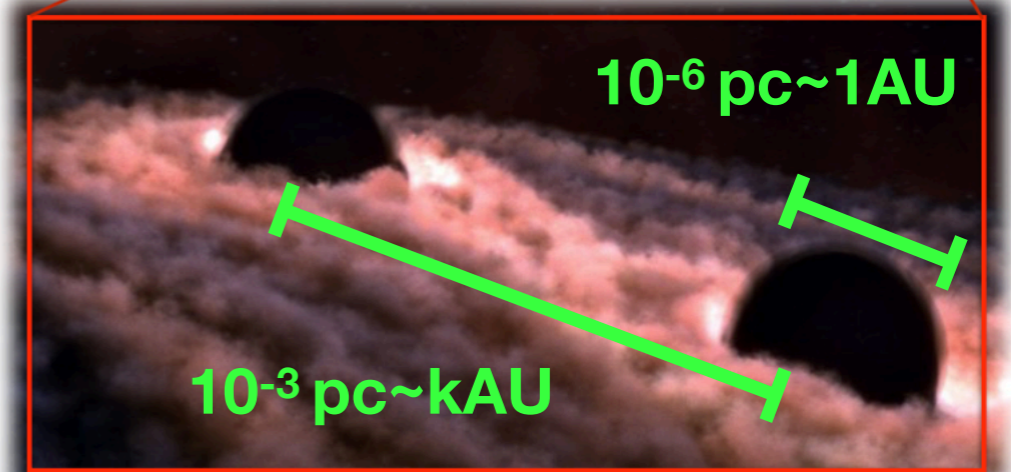
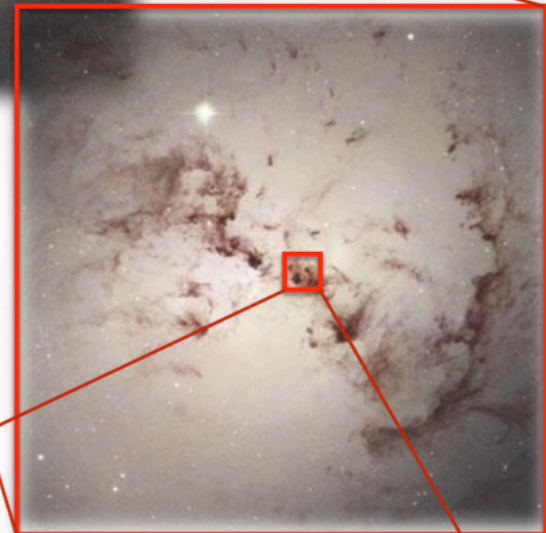
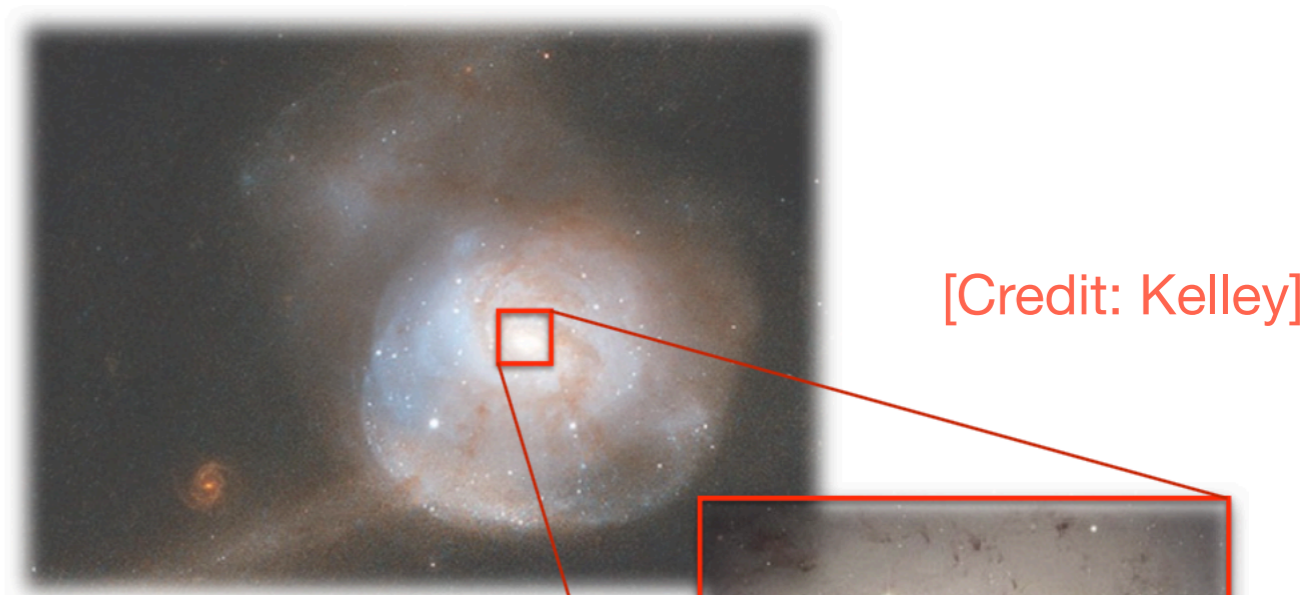
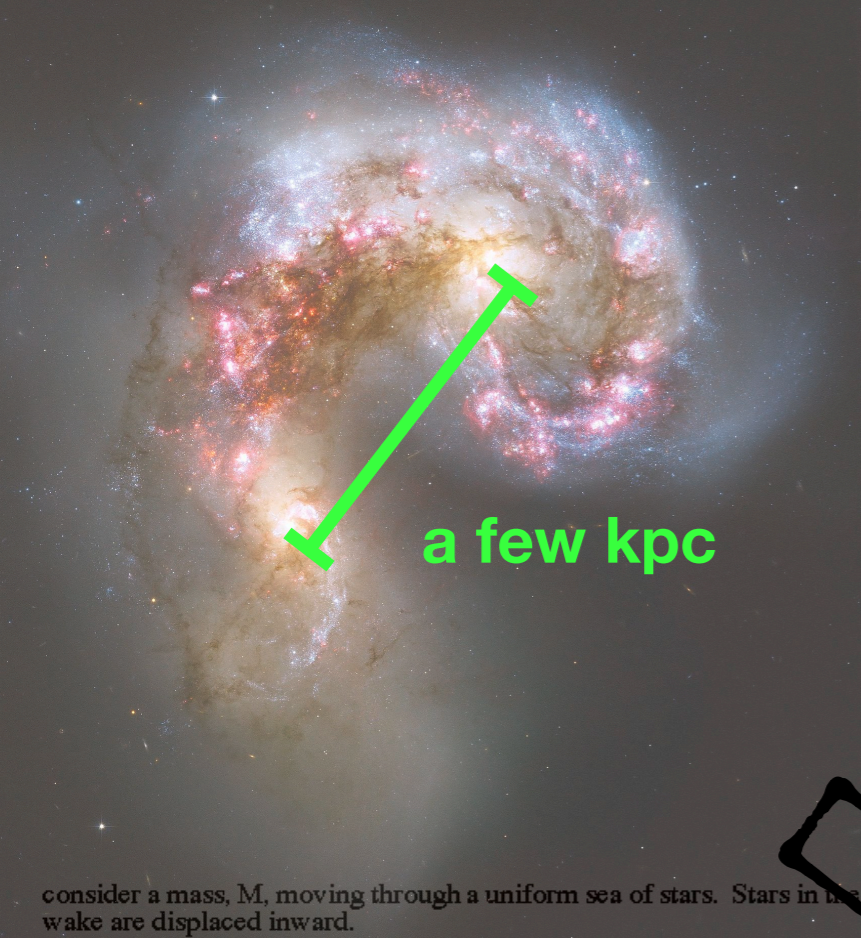
答案：**依然**是五年



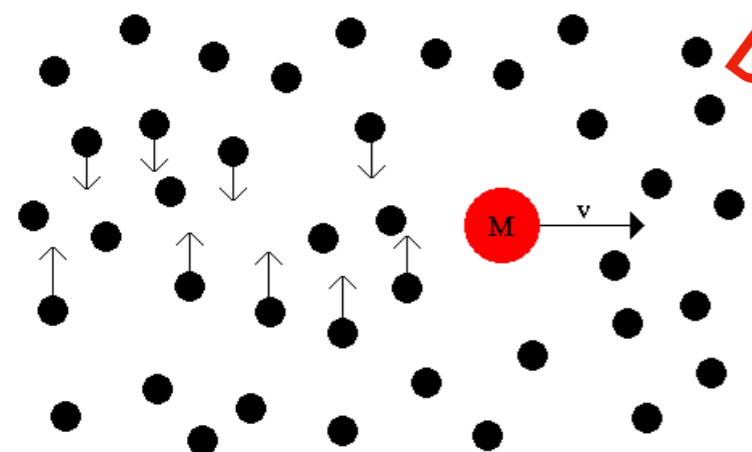
预测未来总是容易的，但是预测过去总是很困难！

Q: 如何**靠谱**地计算随机引力波背景信号的大小?

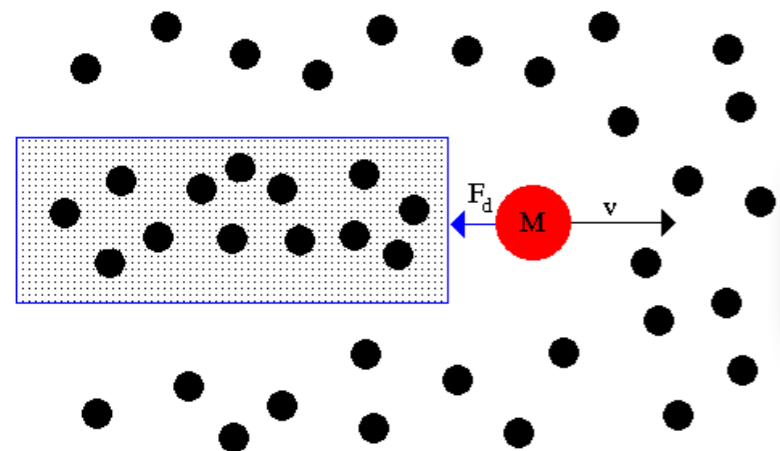




consider a mass, M , moving through a uniform sea of stars. Stars in the wake are displaced inward.



this results in an enhanced region of density behind the mass, with a drag force, F_d known as dynamical friction

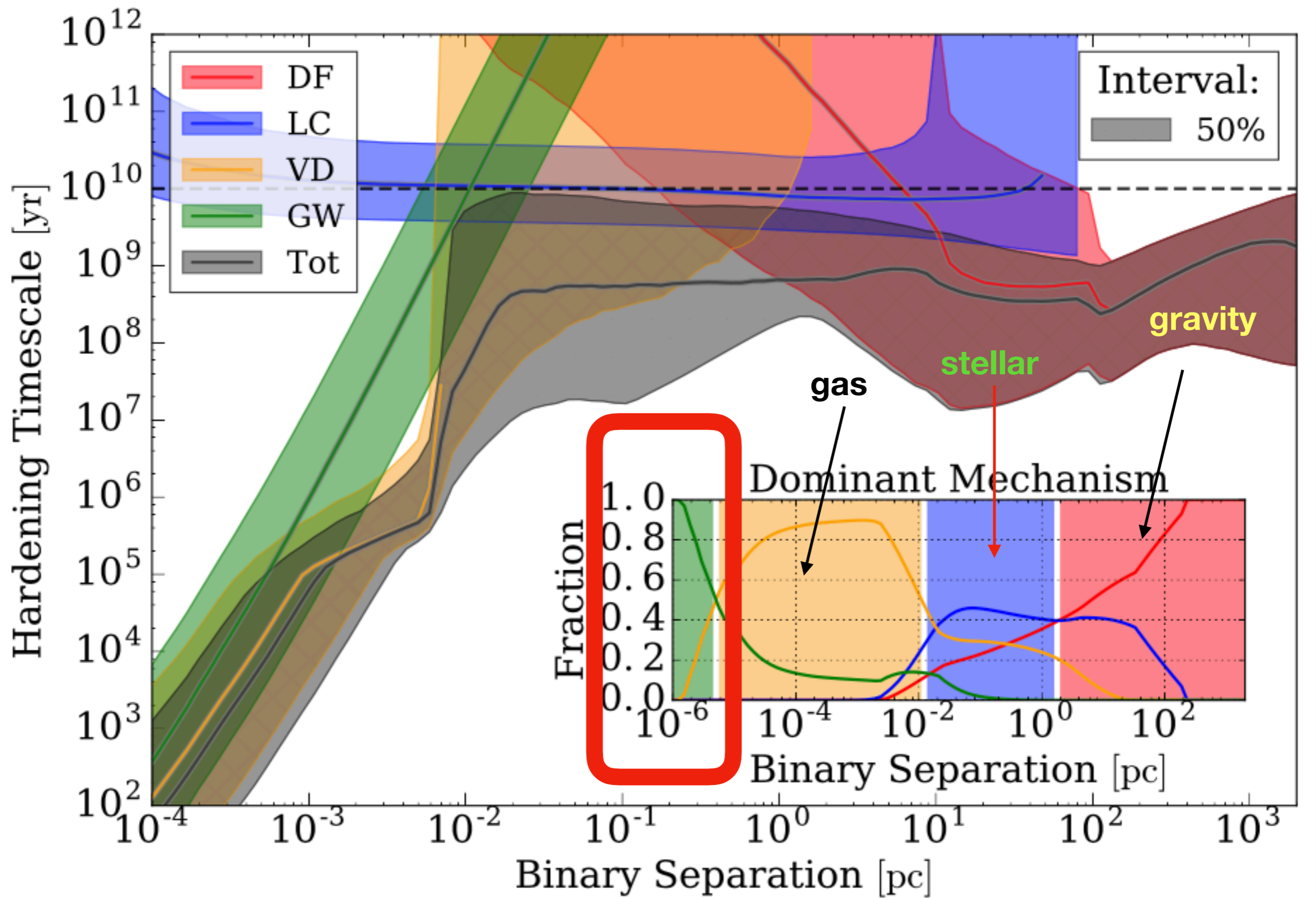


galaxy merger
time ~ 1 Gyr

Dynamical Friction

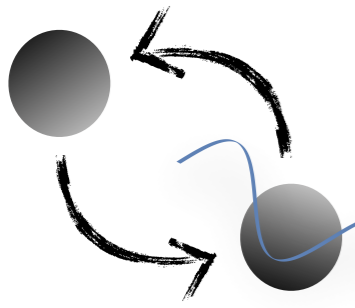
$$\frac{dv}{dt} = -\frac{2\pi G^2 (M + m) \rho}{v^2} \ln \Lambda$$

merger time with GW ~ a few Myr



[Kelley, Blecha, Hernquist, 2017]

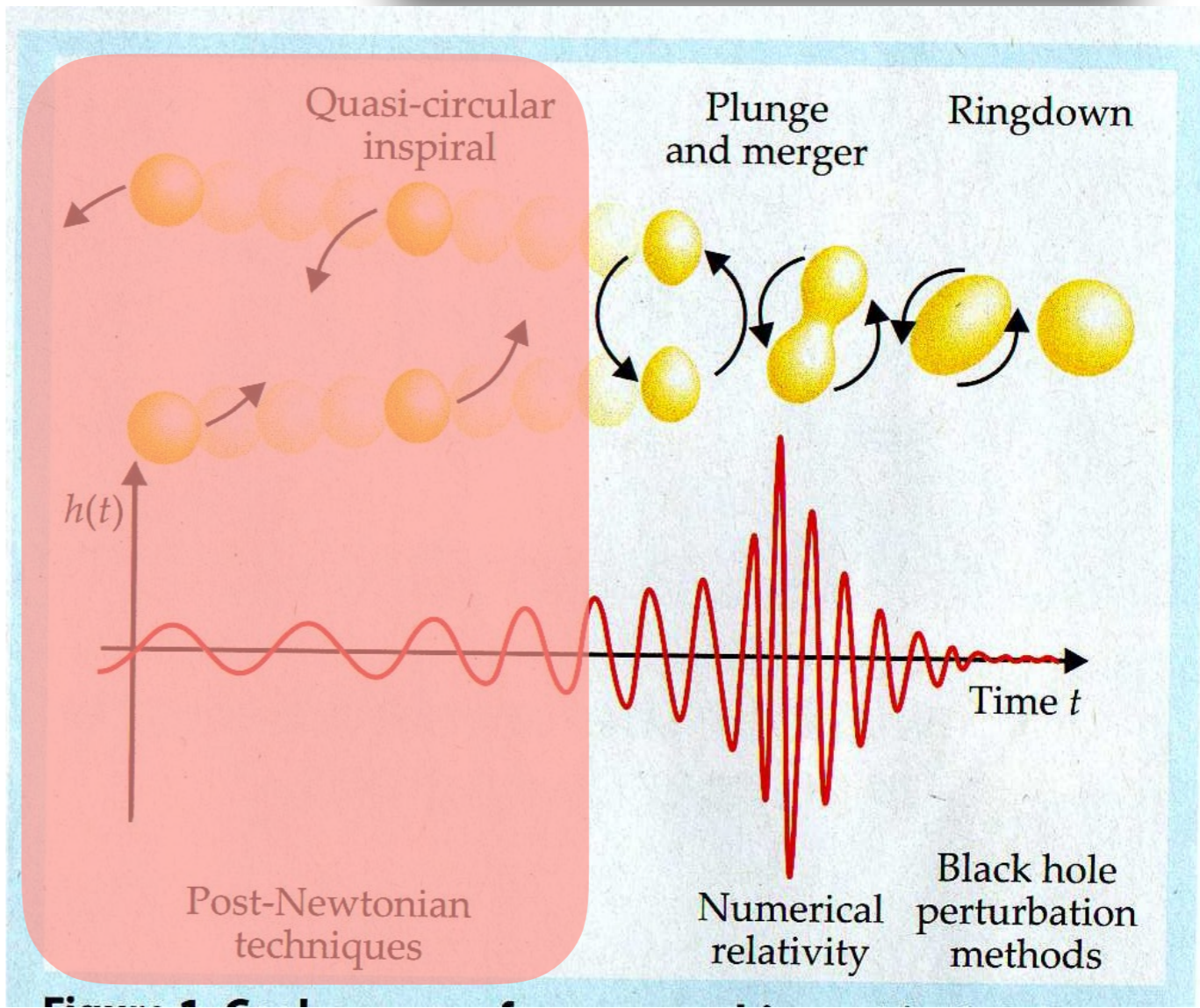
Single binary ~ **circular orbit**, **Quadrupole formula is enough!**



$$f_{GW} = 2f_K \sim [5yr]^{-1}$$

$$\bar{h}_{ij}(t, r) = \frac{2G}{c^4 r} \ddot{I}_{ij}(t - r/c),$$

We can **NOT** observe the inspiral phase, except it is very very nearby!

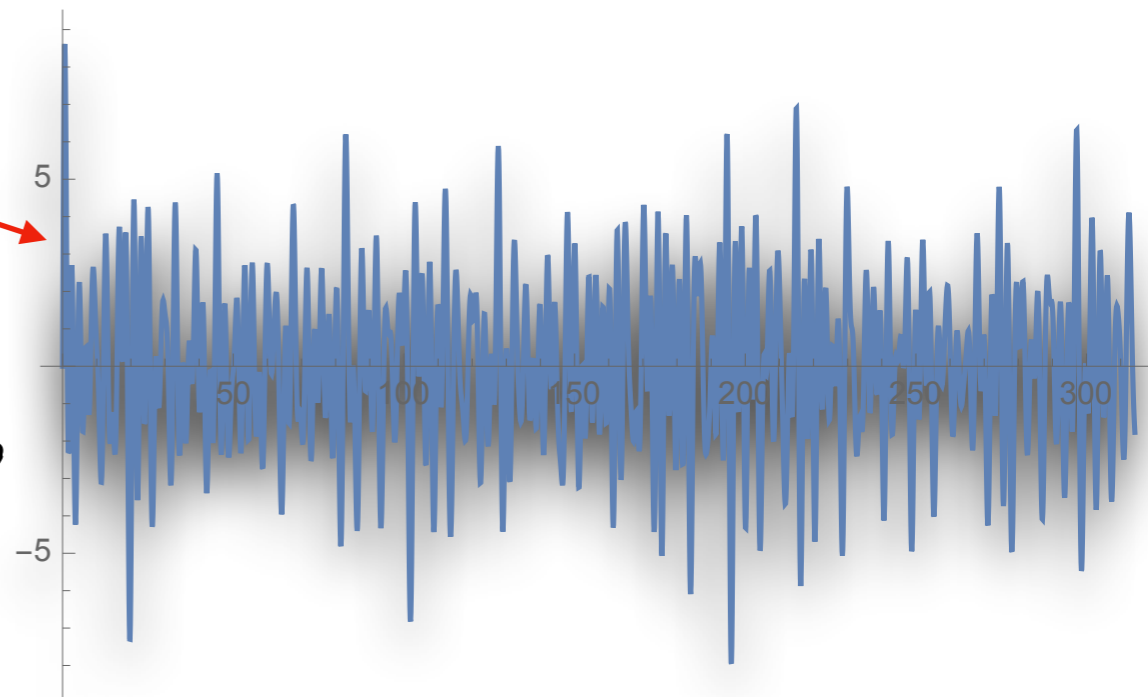
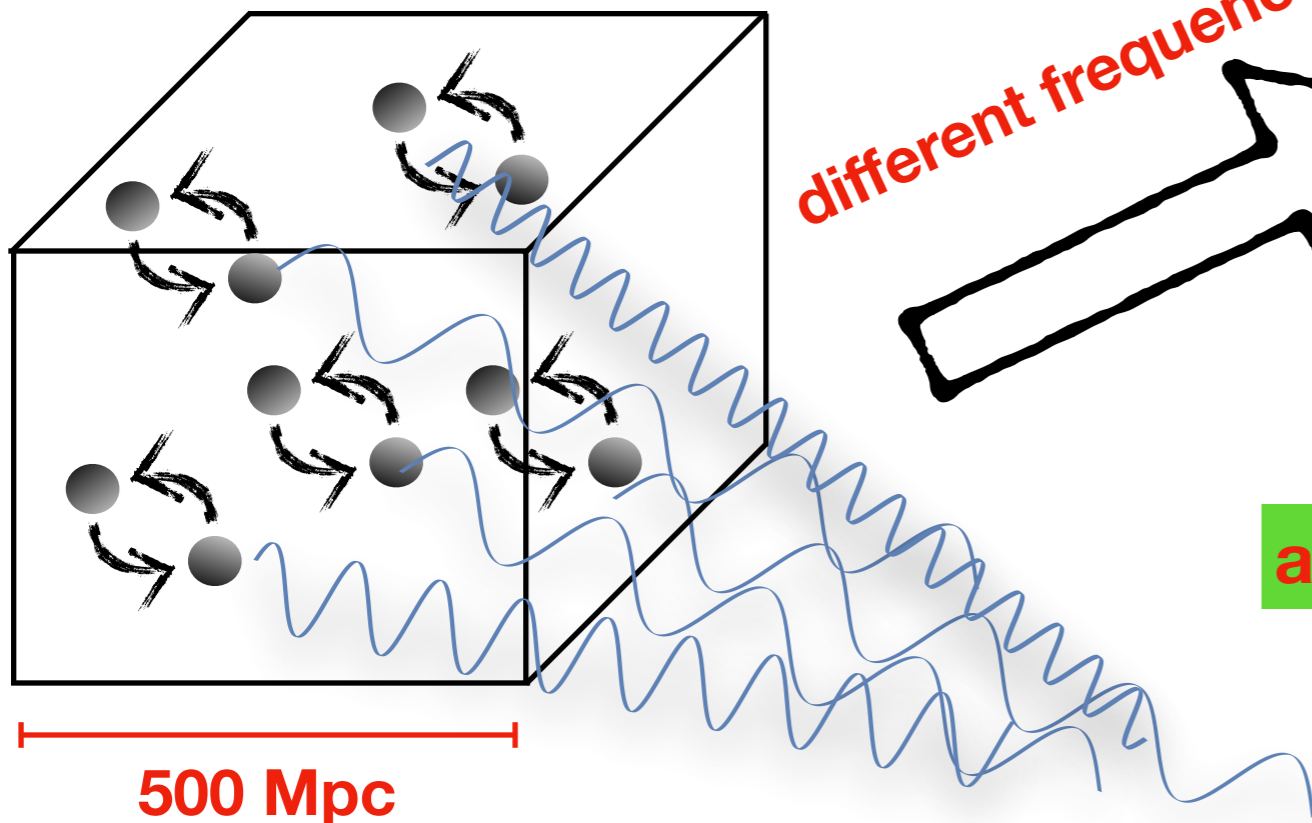


[Credit: 蔡少芬 & wangyi]

Multi-binaries \longrightarrow **GWB**

GW signal we want! (noise-like)

different frequencies



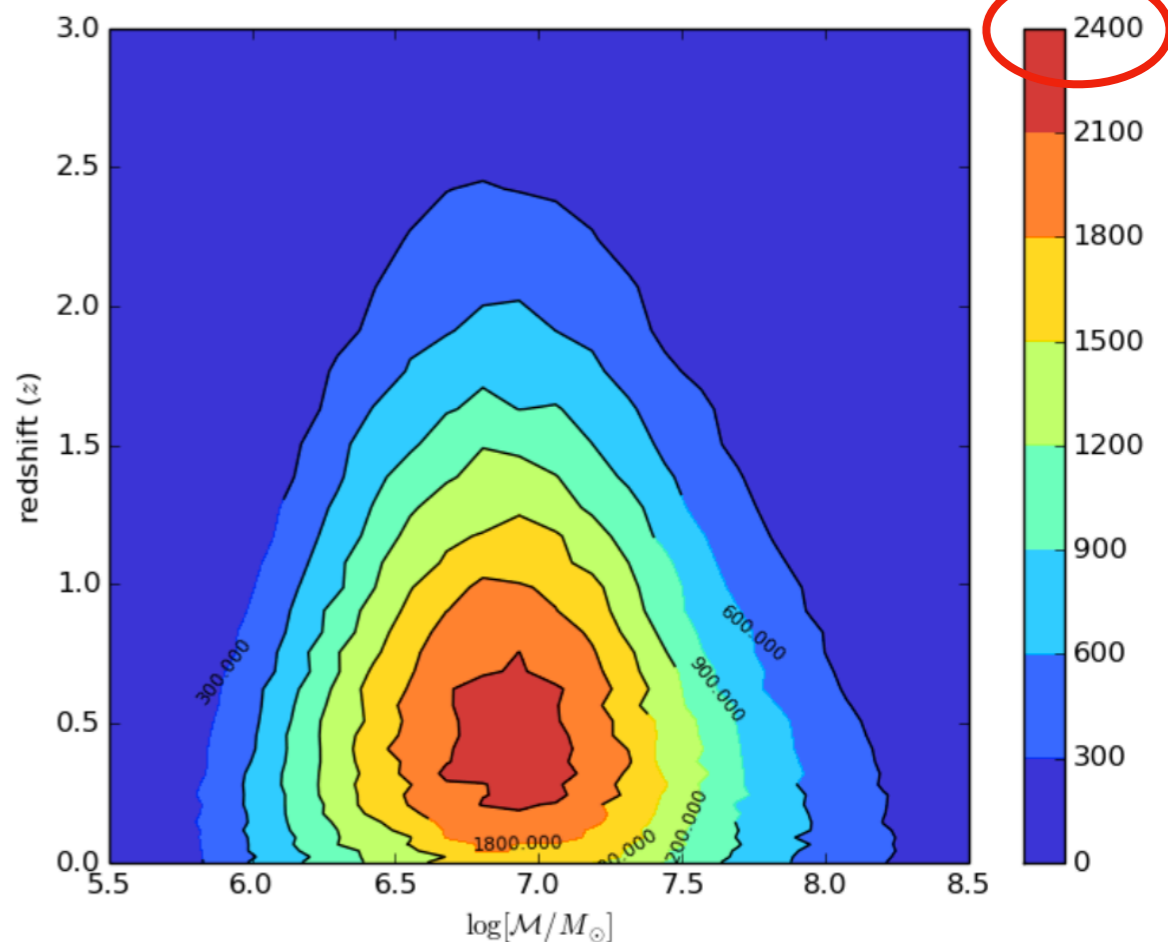
Stochastic in time sequence!

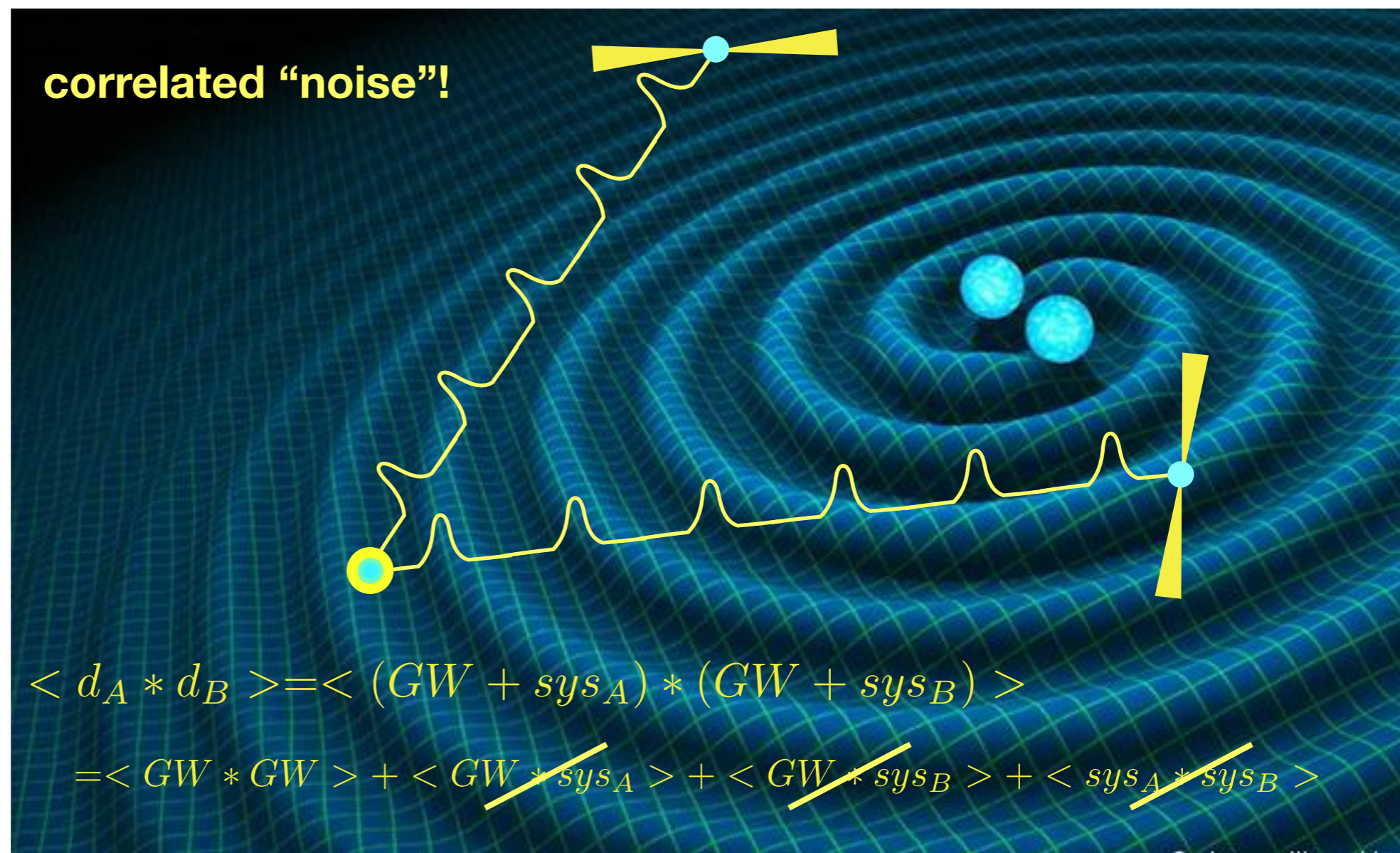
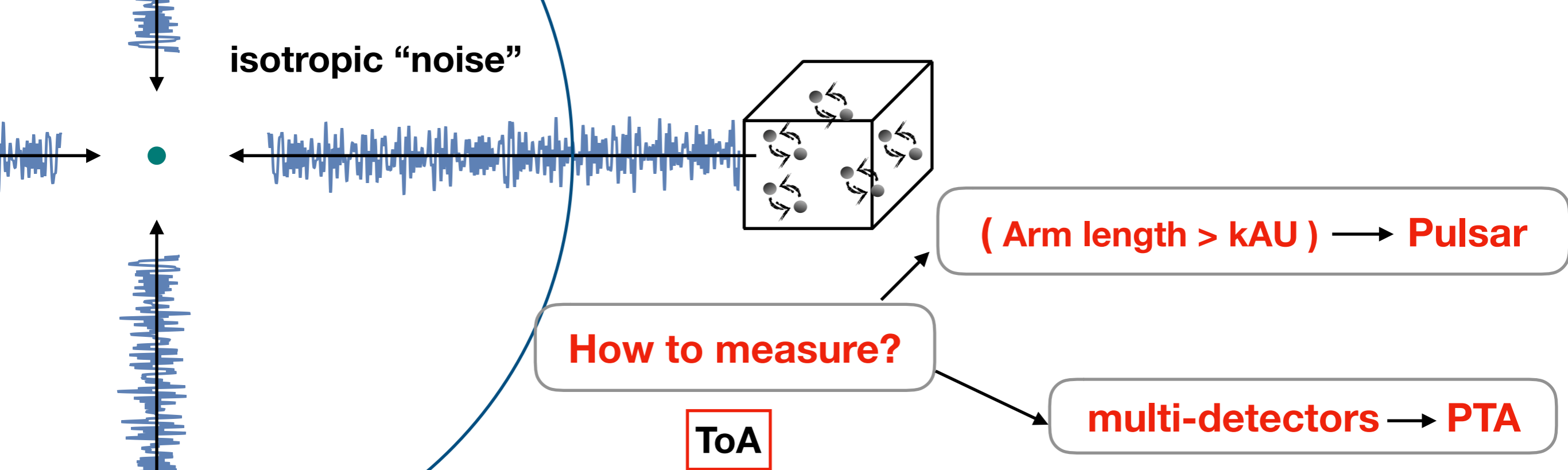
a few 10^3 BBH w. Chirp mass $10^7 M_{\text{solar}}$

$$h_c^2(f) \propto \int \frac{1}{1+z} \frac{dn}{dz} \frac{d\varepsilon_{\text{GW}}}{d \ln f_r} \Big|_{f_r=f(1+z)} dz$$

[Phinney 2001]

Merger rate!



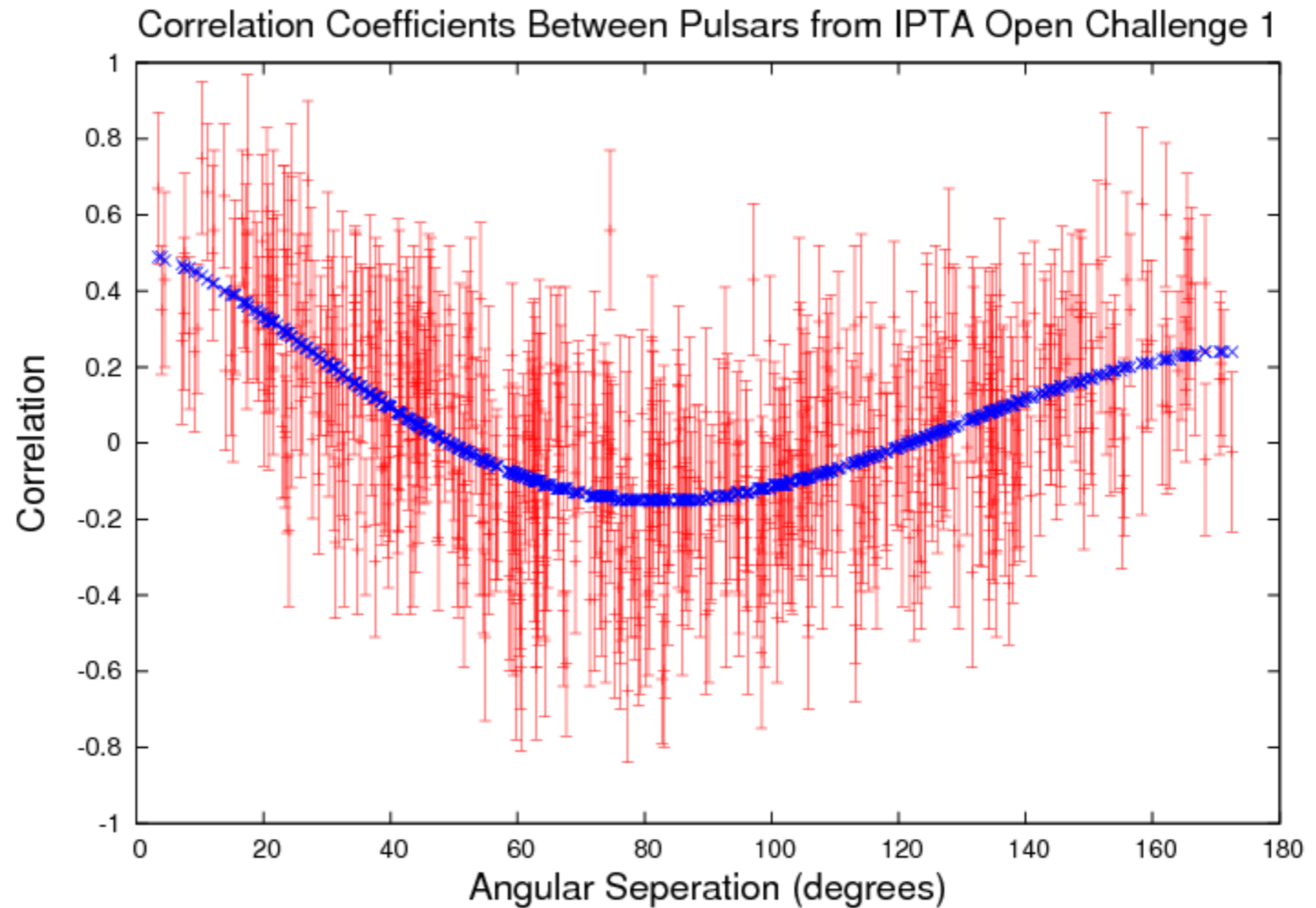


$\langle \text{GW}^* \text{GW} \rangle = \alpha_{ij} \equiv \frac{1}{4\pi} \int \alpha_i \alpha_j d\Omega = \frac{1 - \cos \gamma_{ij}}{2} \ln \left(\frac{1 - \cos \gamma_{ij}}{2} \right) - \frac{1}{6} \frac{1 - \cos \gamma_{ij}}{2} + \frac{1}{3},$ (5)

average over GW from all direction

where γ_{ij} is the angle between the two pulsars.

[Hellings & Downs 1983]



$$h_c^2(f) \propto \int \frac{1}{1+z} \frac{dn}{dz} \frac{d\varepsilon_{\text{GW}}}{d \ln f_r} \Big|_{f_r=f(1+z)} dz$$

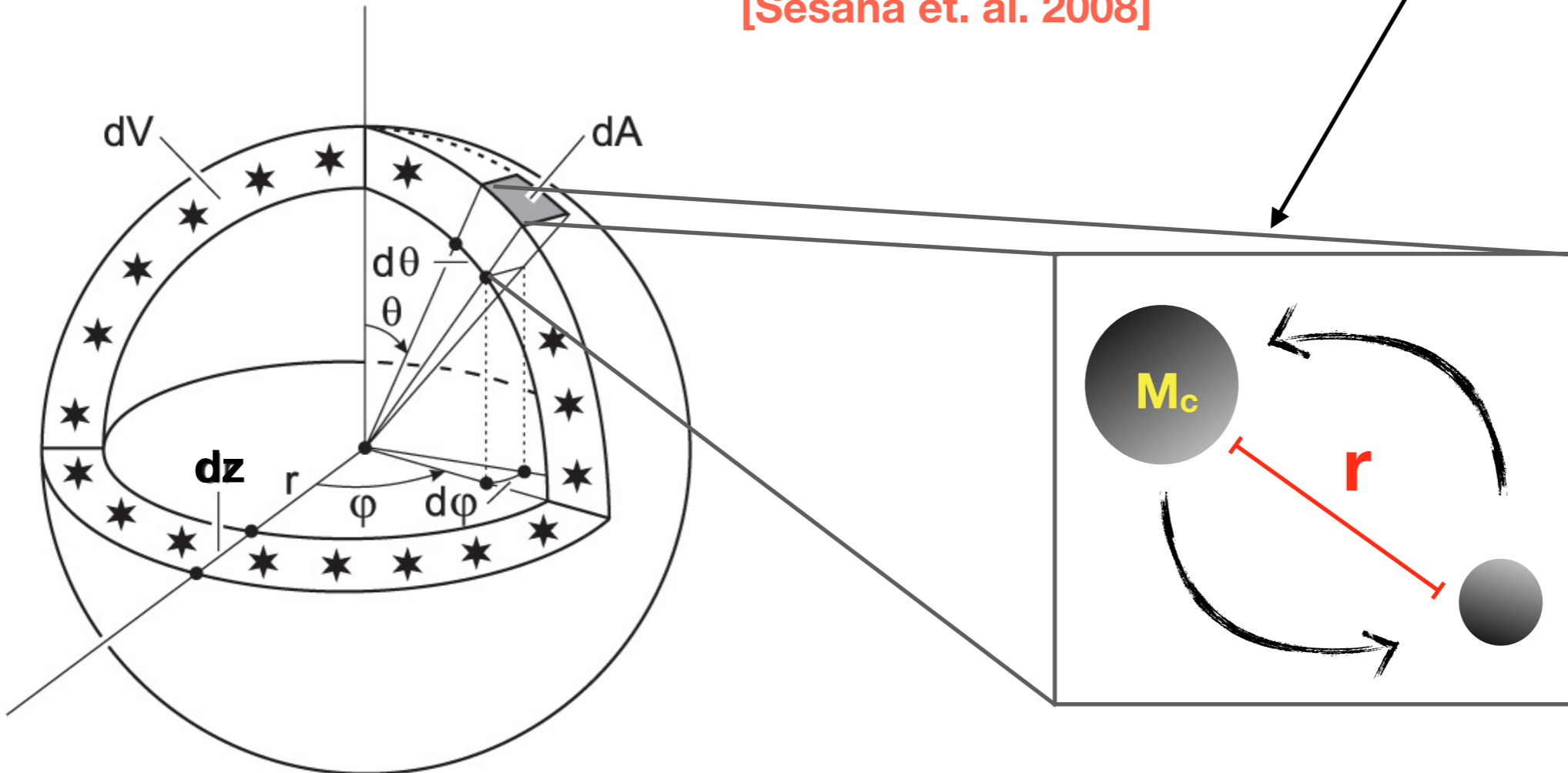
[Phinney 2001]

more accurate

$$h_c^2(f) = \int_0^\infty dz \int_0^\infty d\mathcal{M} \frac{d^3 N}{dz d\mathcal{M} d \ln f_r} h^2(f_r),$$

BBH number per comoving volume, in such configuration

[Sesana et. al. 2008]



[Peters 1964]

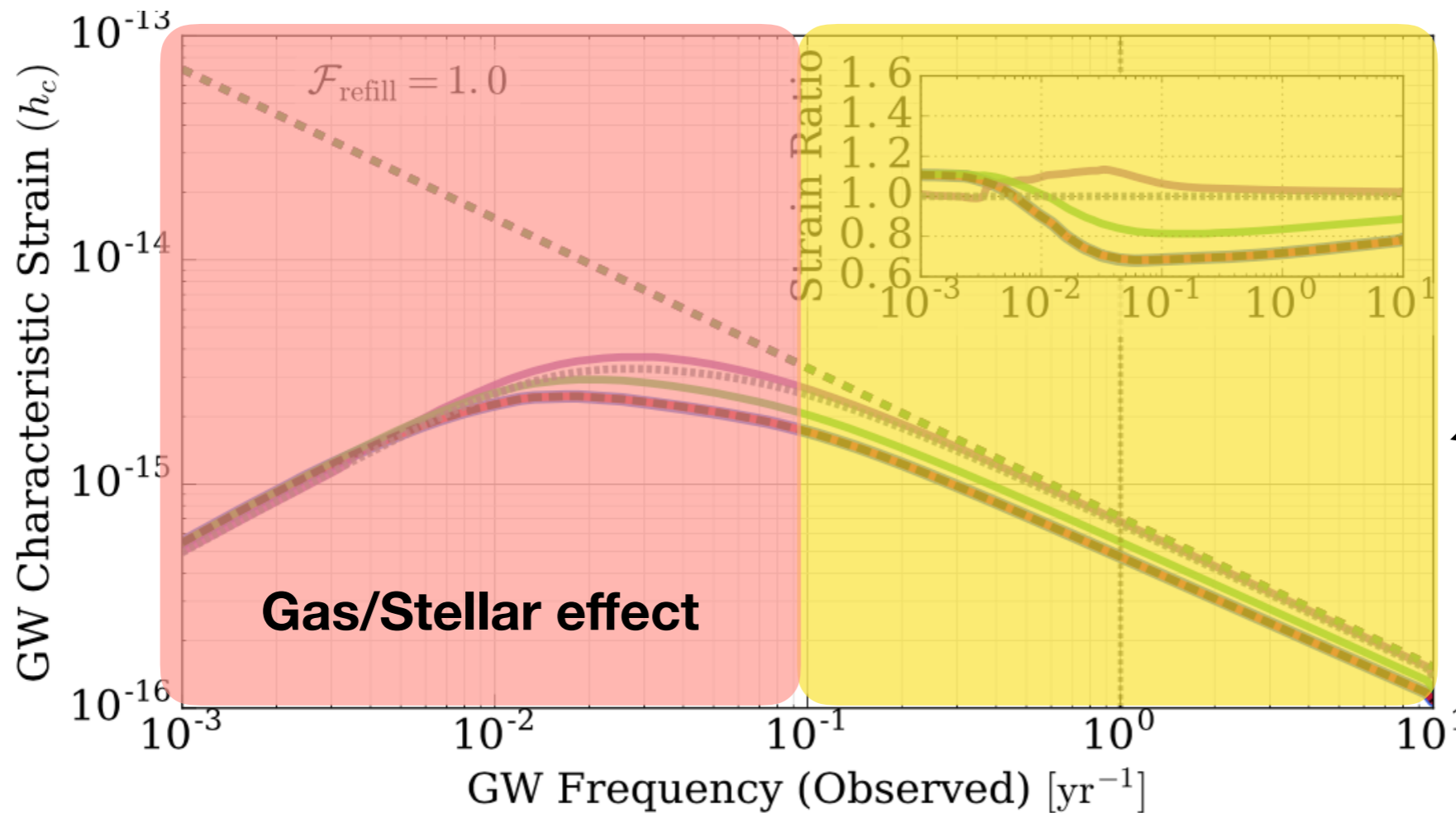
$$dt/d\ln f = \frac{5}{64\pi^{8/3}} \mathcal{M}^{-5/3} f_r^{-8/3}$$

time spend in per logarithmic frequency

$$h_c = A(f/f_0)^{-2/3}$$

Major eq.

$$h_c^2(f) = \frac{4f^{-4/3}}{3\pi^{1/3}} \int \int dz d\mathcal{M} \frac{d^2 n}{dz d\mathcal{M}} \frac{1}{(1+z)^{1/3}} \mathcal{M}^{5/3}$$



energy loss
only due to
GW radiation
during inspiral

[Credit: Kelley]

$$\frac{d^2 n}{dz dM}$$

phenomenology

[Sesana 2012]

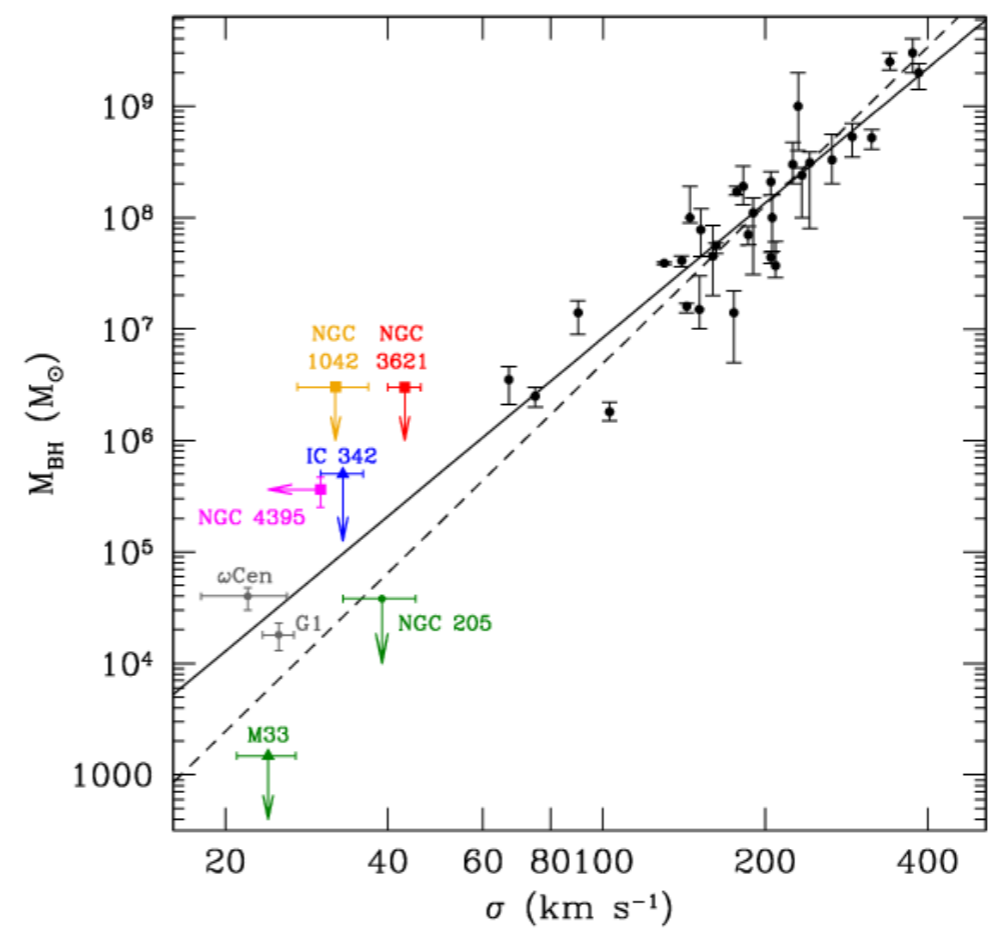
Preliminary
Semi-analytic
Modelling

SMBH merger modelling
[Kelley et. al. 2017]

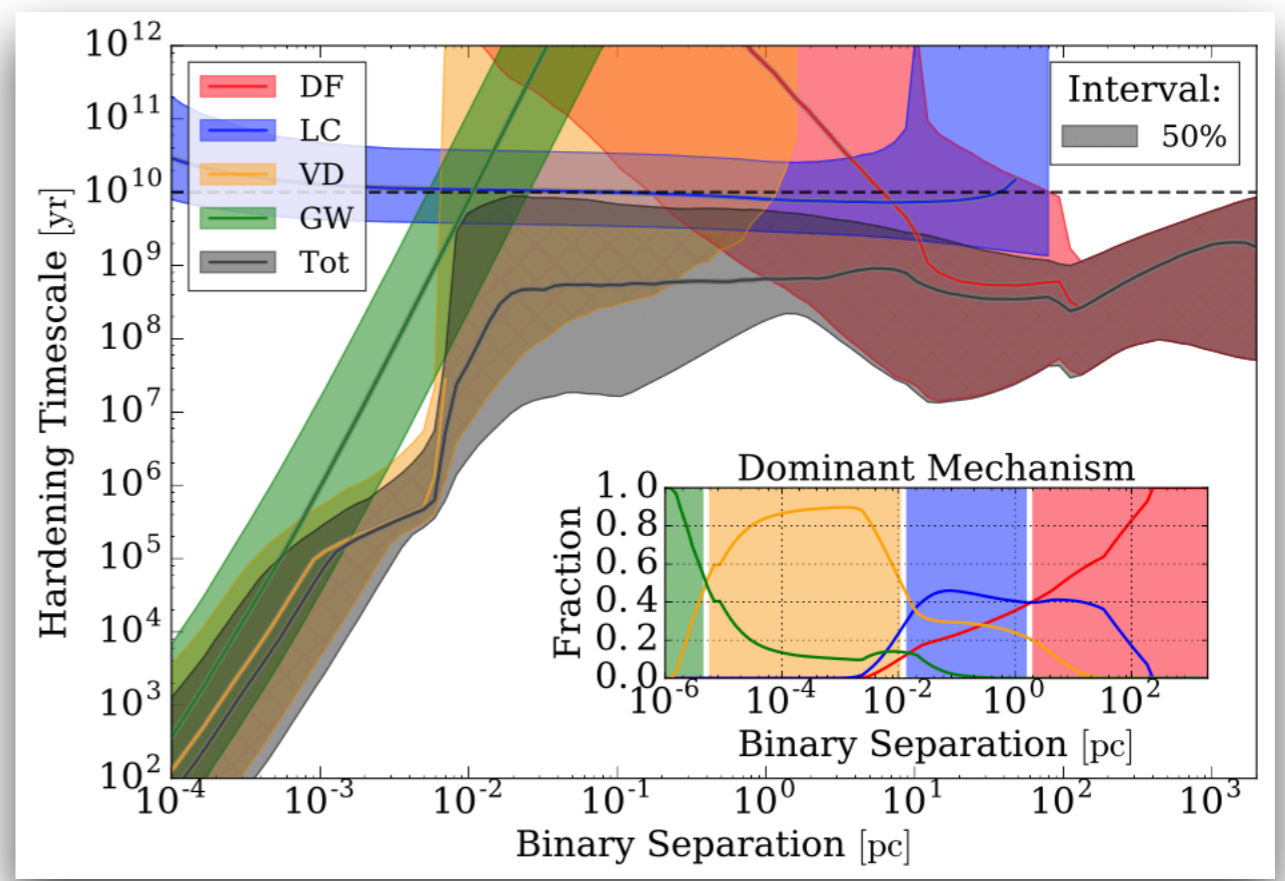
$$\frac{d^2 n_g}{dz dM}$$

[Sesana 2008]

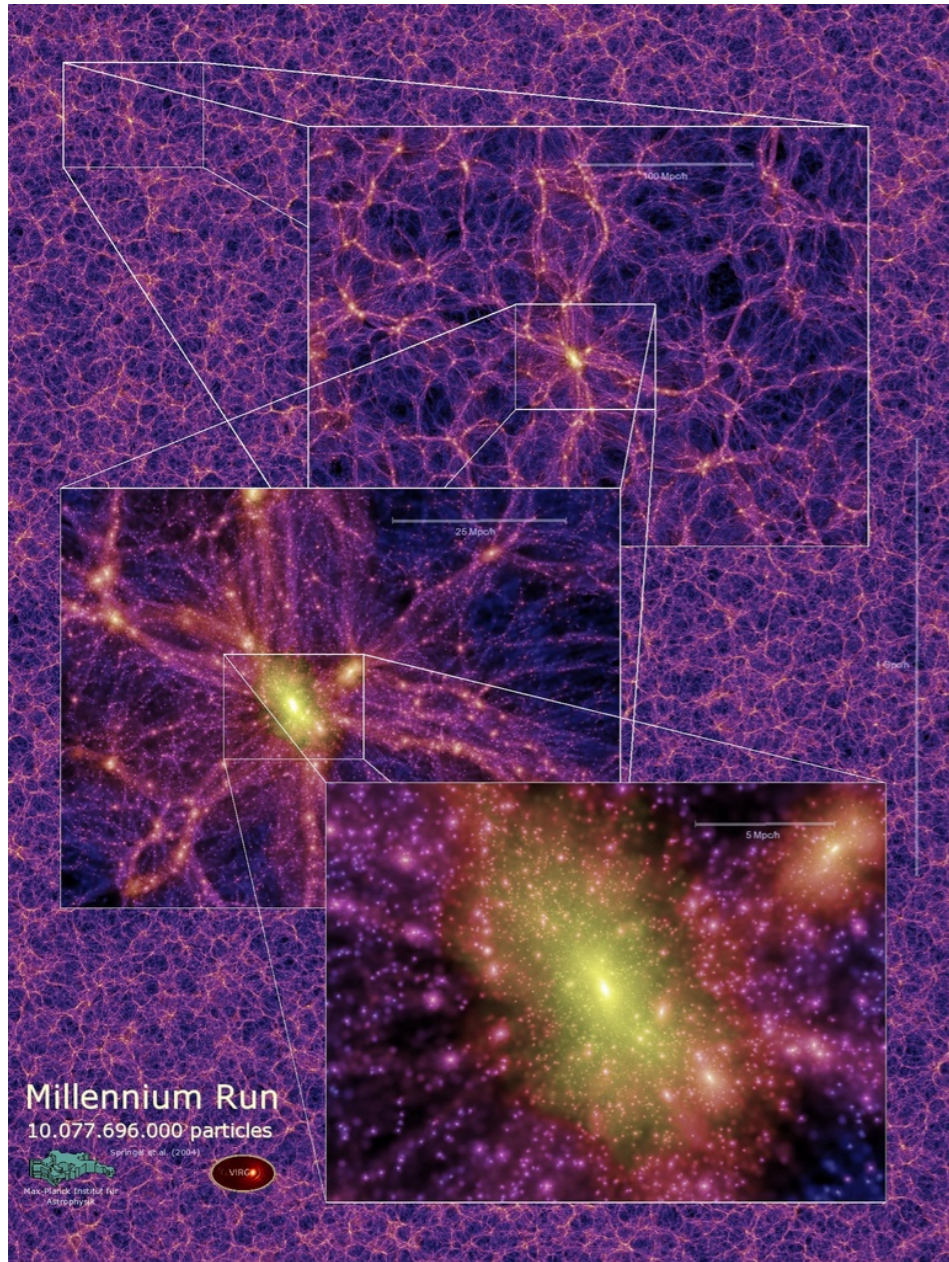
e.g. galaxy mass function
is calculated via EPS formalism
& with only hundreds of DM halos



$M - \sigma$
relation



Our method: Semi-Analytic Model (SAM) of galaxy formation



$V \sim 500^3 \text{ Mpc}^3$

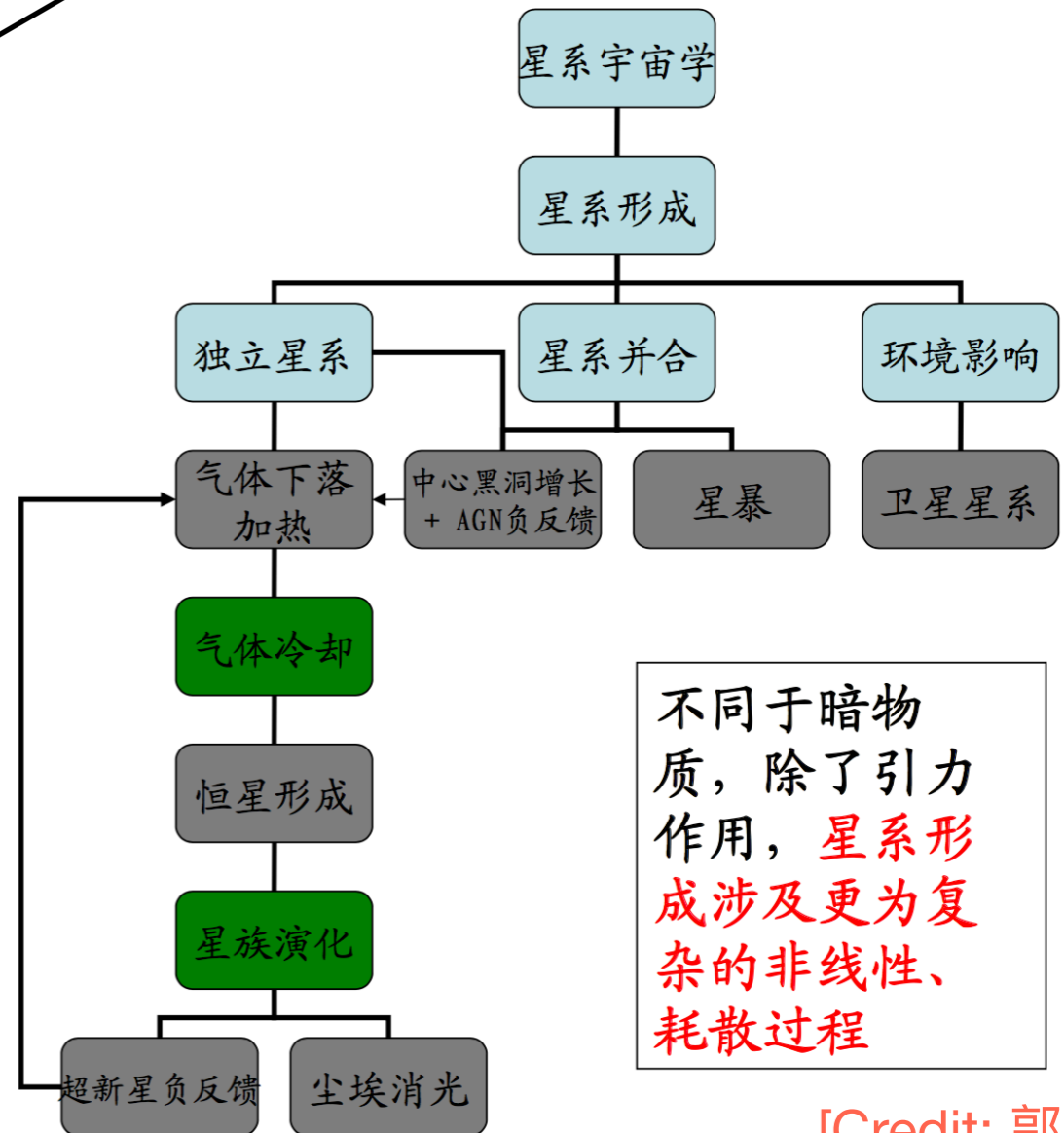
8668809 SMBHs,
51538704 galaxies
in total

code: L-galaxies

1. Run N-body simulation \longrightarrow DM halo merge tree
2. Add SN, AGN, hot/cold gas, stellar, galaxies, BHs

directly read
BH mass function

$$\frac{d^2 n}{dz dM}$$



[Credit: 郭琦]

BH Self-regulated growth & feedback

Quasar mode: (gas-rich merger)

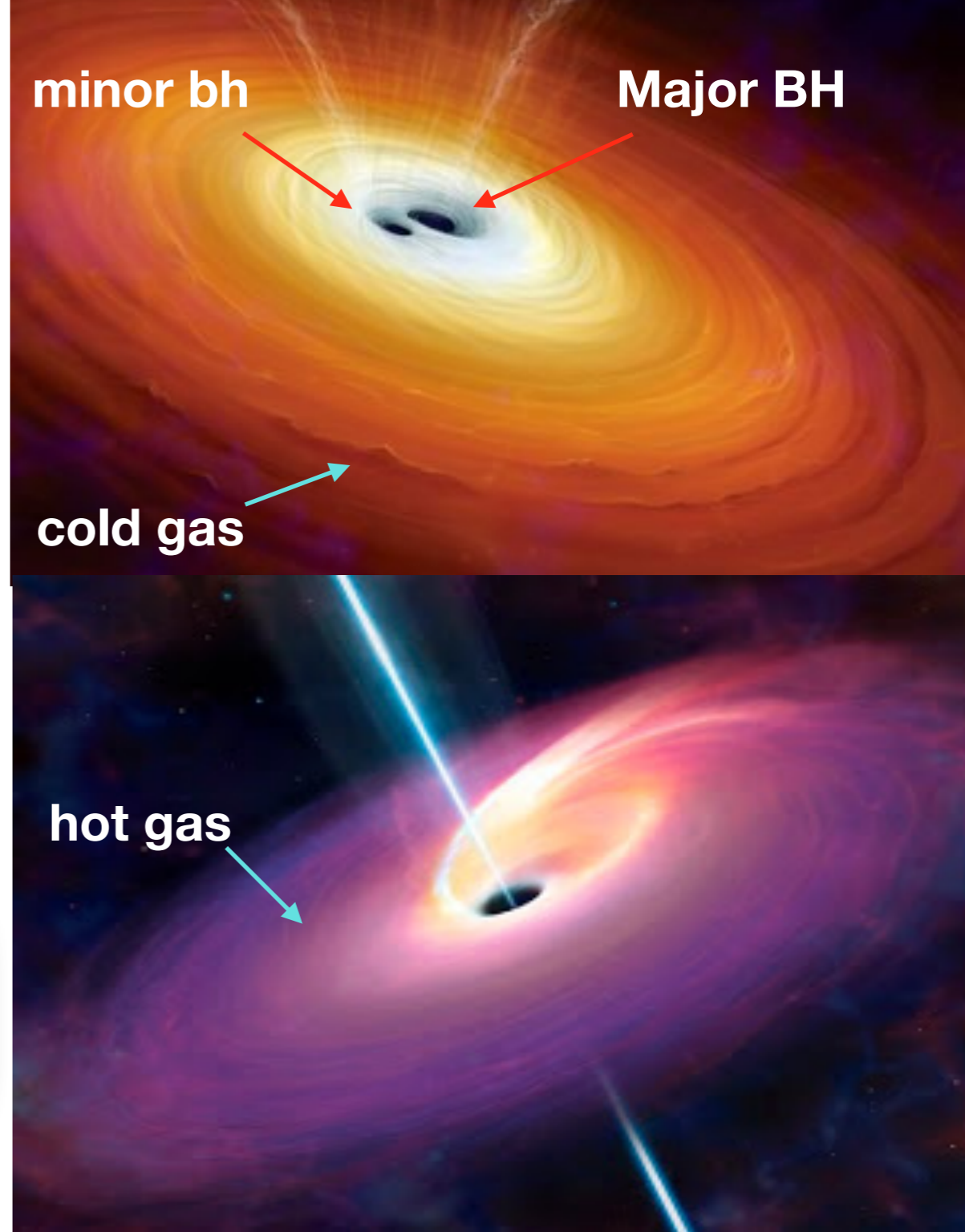
$$M_{\text{bh},f} = M_{\text{bh},\text{maj}} + M_{\text{bh},\text{min}} + \Delta M_{\text{bh},Q},$$
$$\Delta M_{\text{bh},Q} = \frac{f_{\text{bh}}(M_{\text{min}}/M_{\text{maj}})M_{\text{cold}}}{1 + 280 \text{ km s}^{-1}/V_{\text{vir}}},$$

Radio mode: (hot gas accretion)

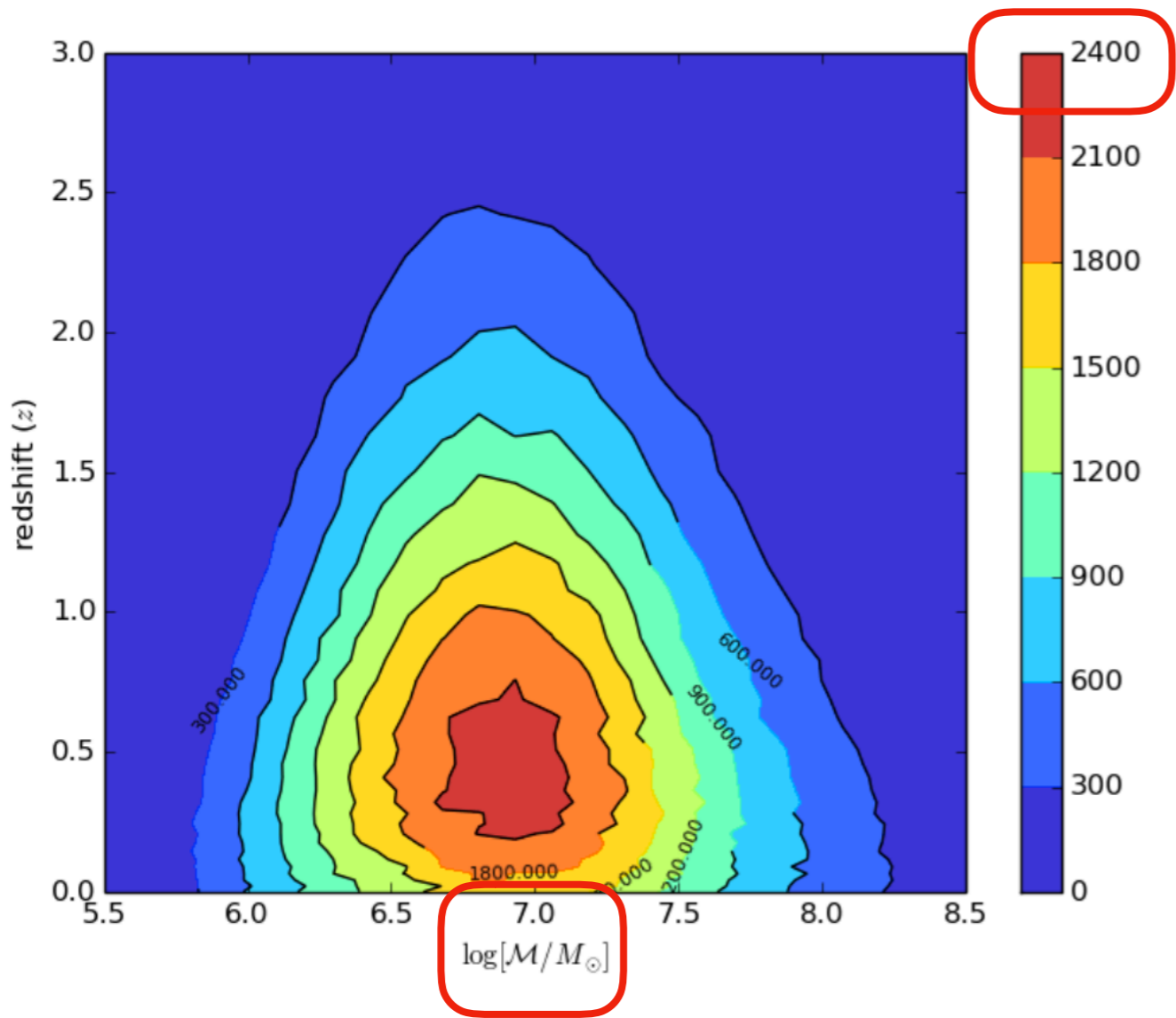
$$\dot{M}_{\text{bh}} = \kappa \left(\frac{f_{\text{hot}}}{0.1} \right) \left(\frac{V_{\text{vir}}}{200 \text{ km s}^{-1}} \right)^3 \left(\frac{M_{\text{bh}}}{10^8 h^{-1} M_{\odot}} \right) M_{\odot} \text{ yr}^{-1}$$

$$\dot{E}_{\text{radio}} = 0.1 \dot{M}_{\text{bh}} c^2$$

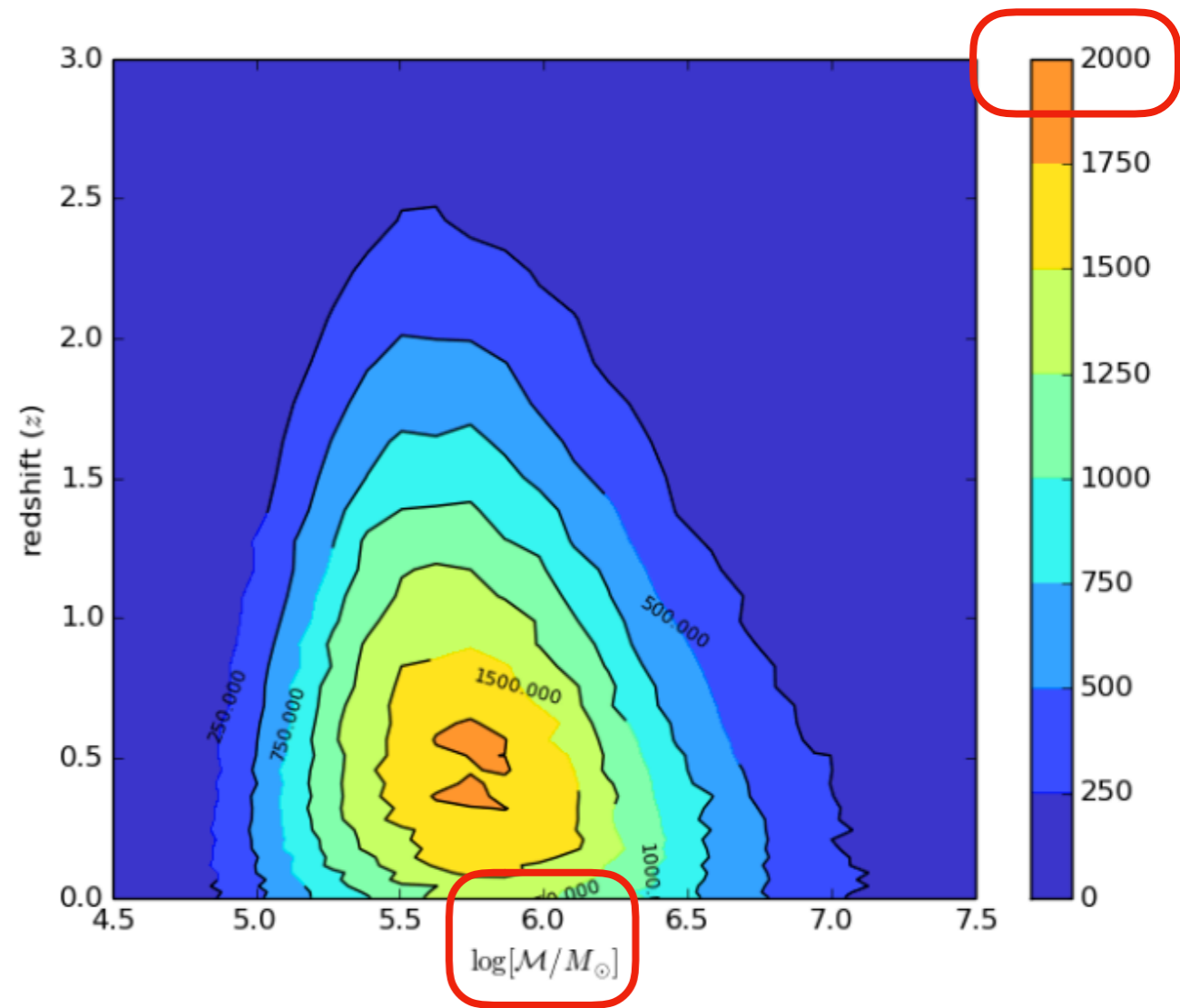
10% energy deposit into relativistic jet



$$\frac{d^2n}{dzdM}$$

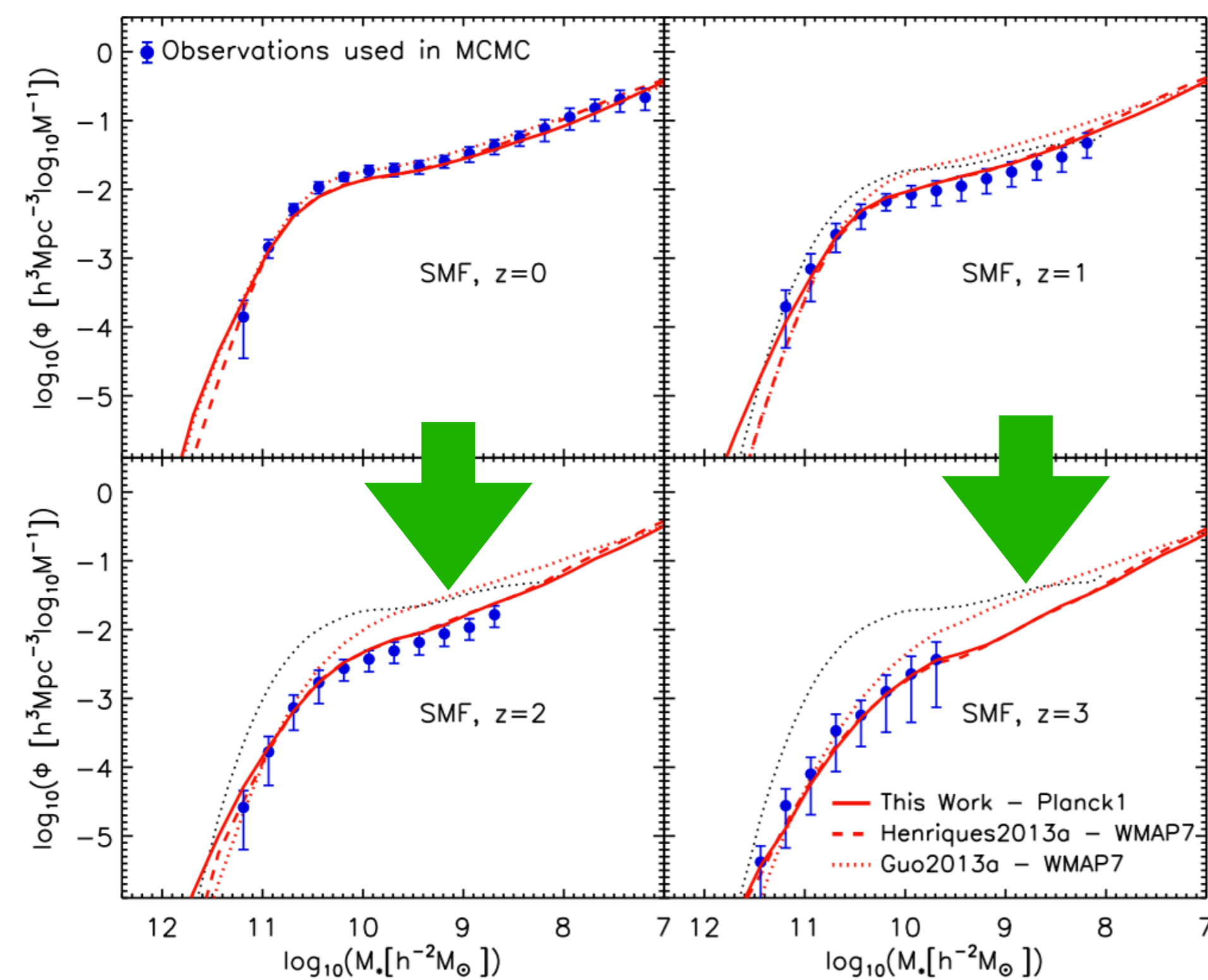


Guo 2013
based WMAP7 cosmology

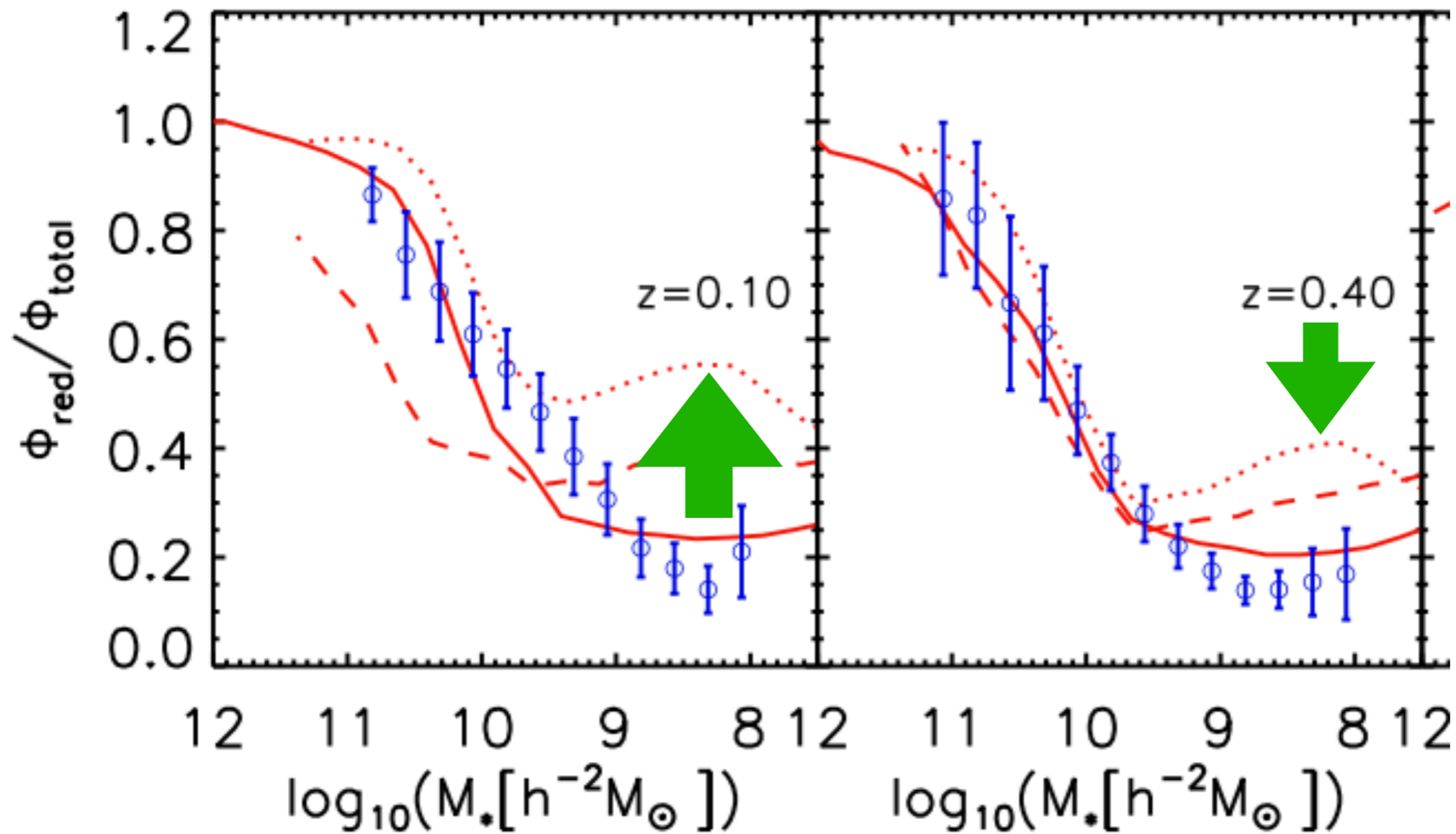


Henriques 2015
based Planck cosmology

[Henriques et. al. 2015]



overly early formation of low-mass galaxies in Guo2013



overly large fraction of them that are passive at late times in Guo 2013

星系宇宙学

星系形成

独立星系

星系并合

环境影响

气体下落
加热

中心黑洞增长
+ AGN 负反馈

星暴

卫星星系

气体冷却

恒星形成

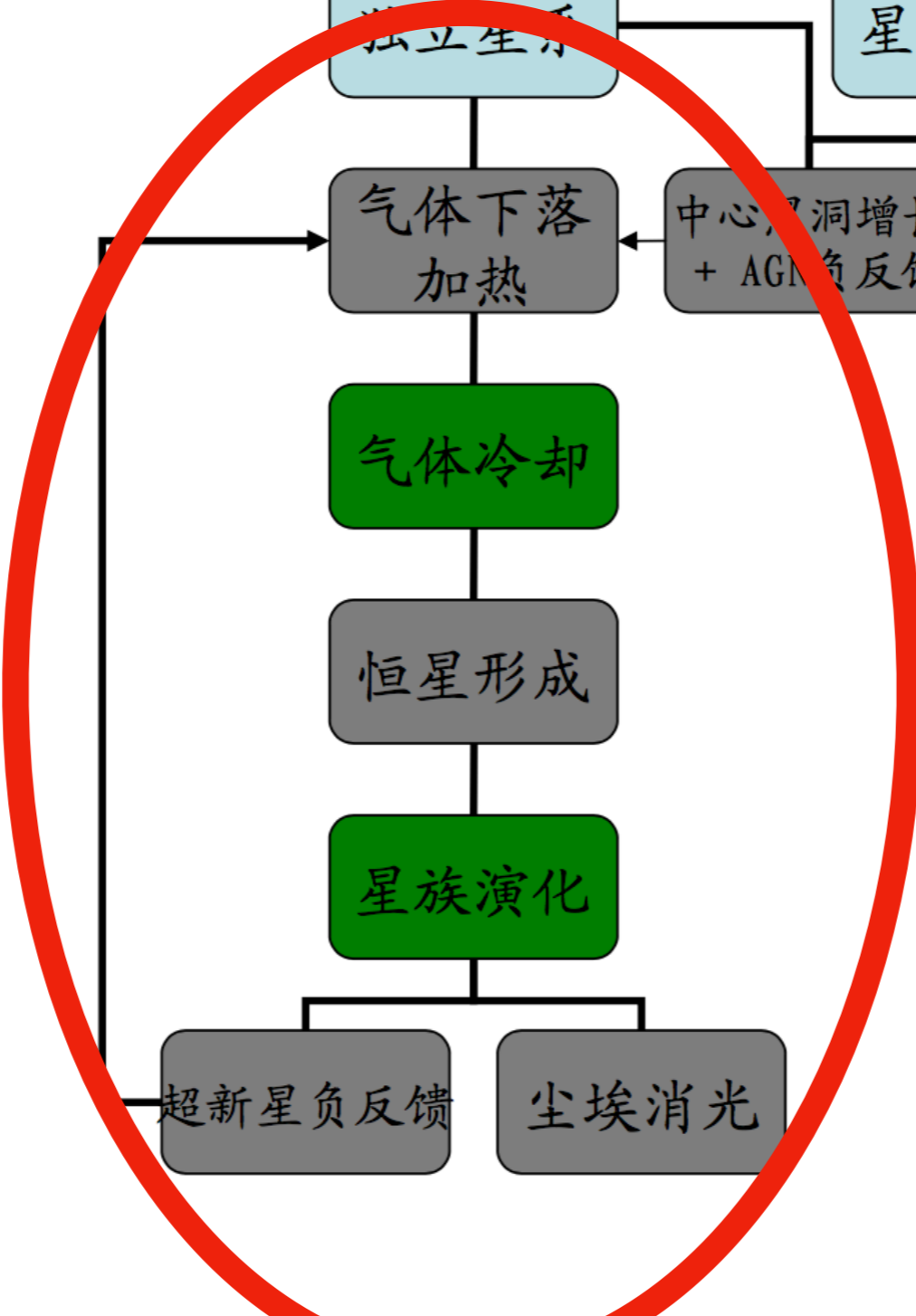
星族演化

超新星负反馈

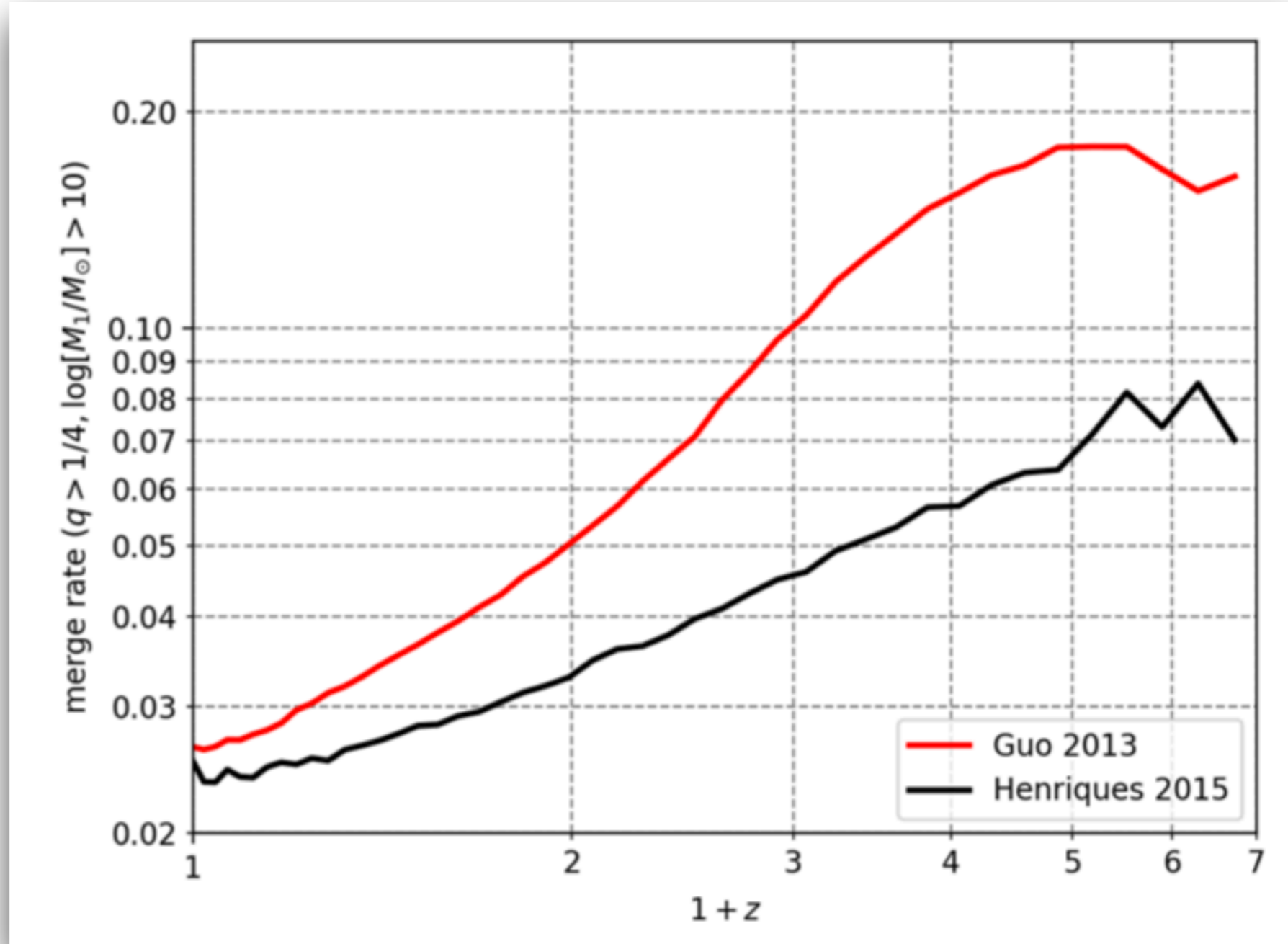
尘埃消光

解决方案：
拉长
该过程
的时标

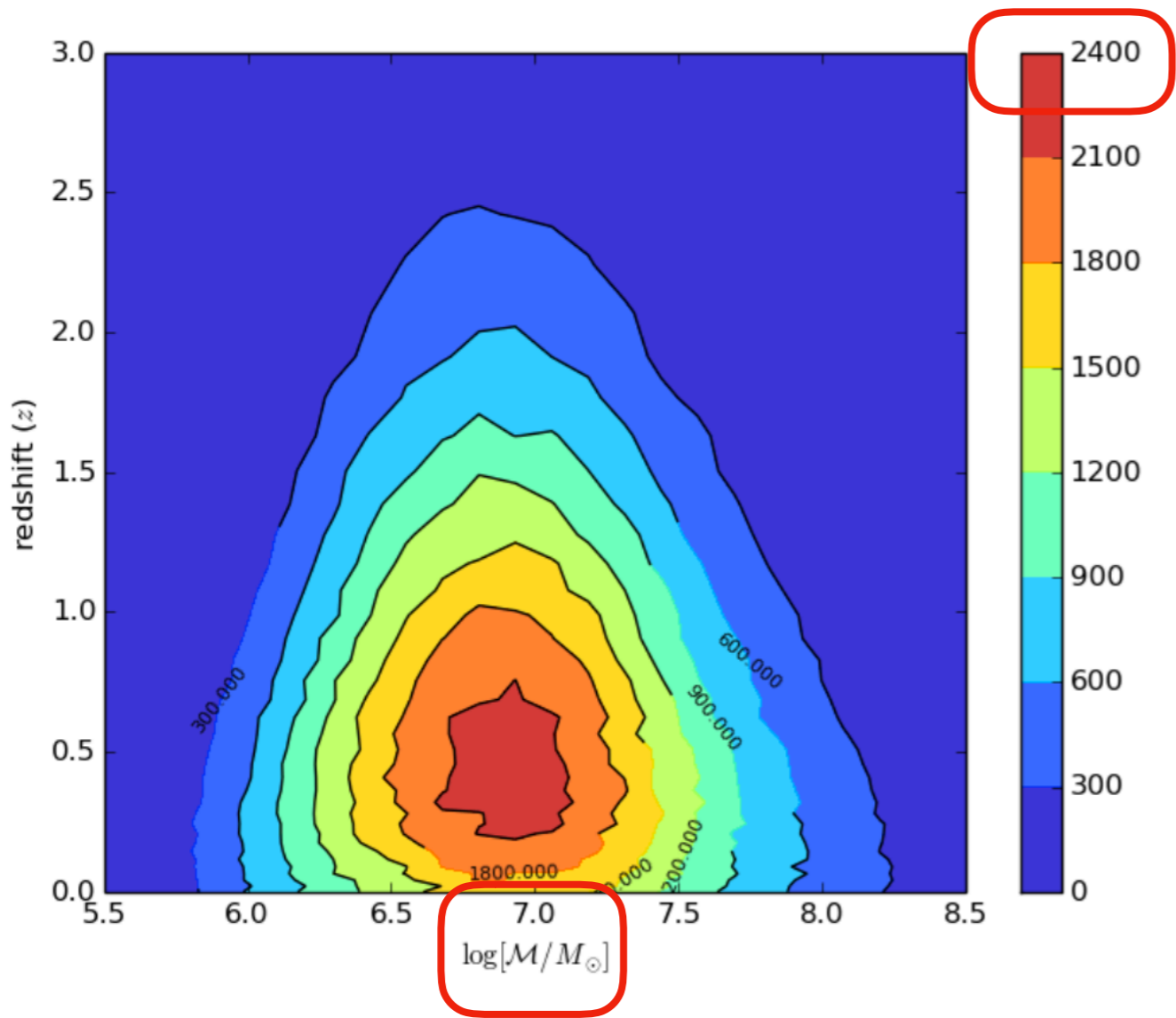
不同于暗物质，除了引力作用，星系形成涉及更为复杂的非线性、耗散过程



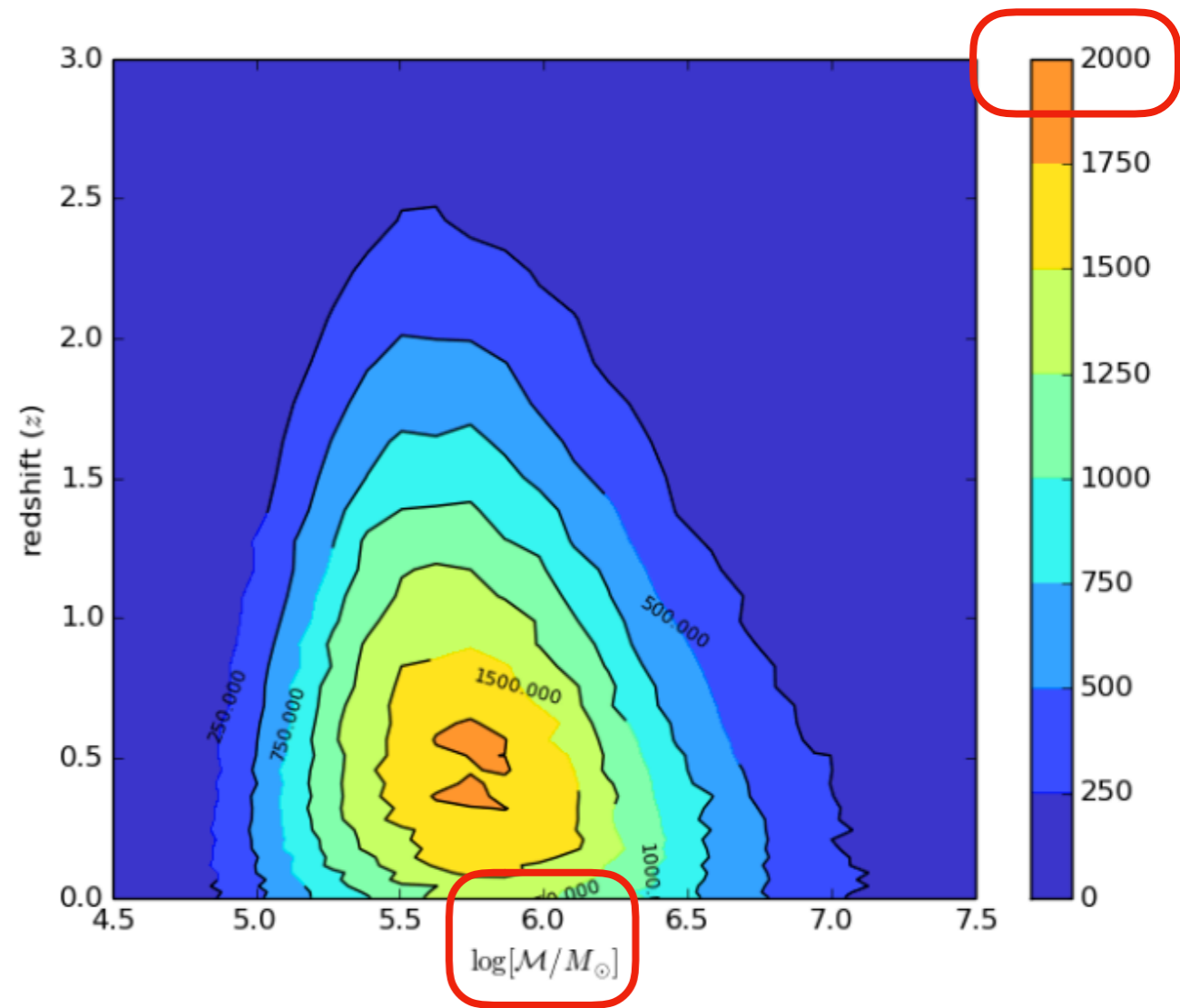
$$\frac{dn_g}{dz}$$



$$\frac{d^2n}{dzdM}$$

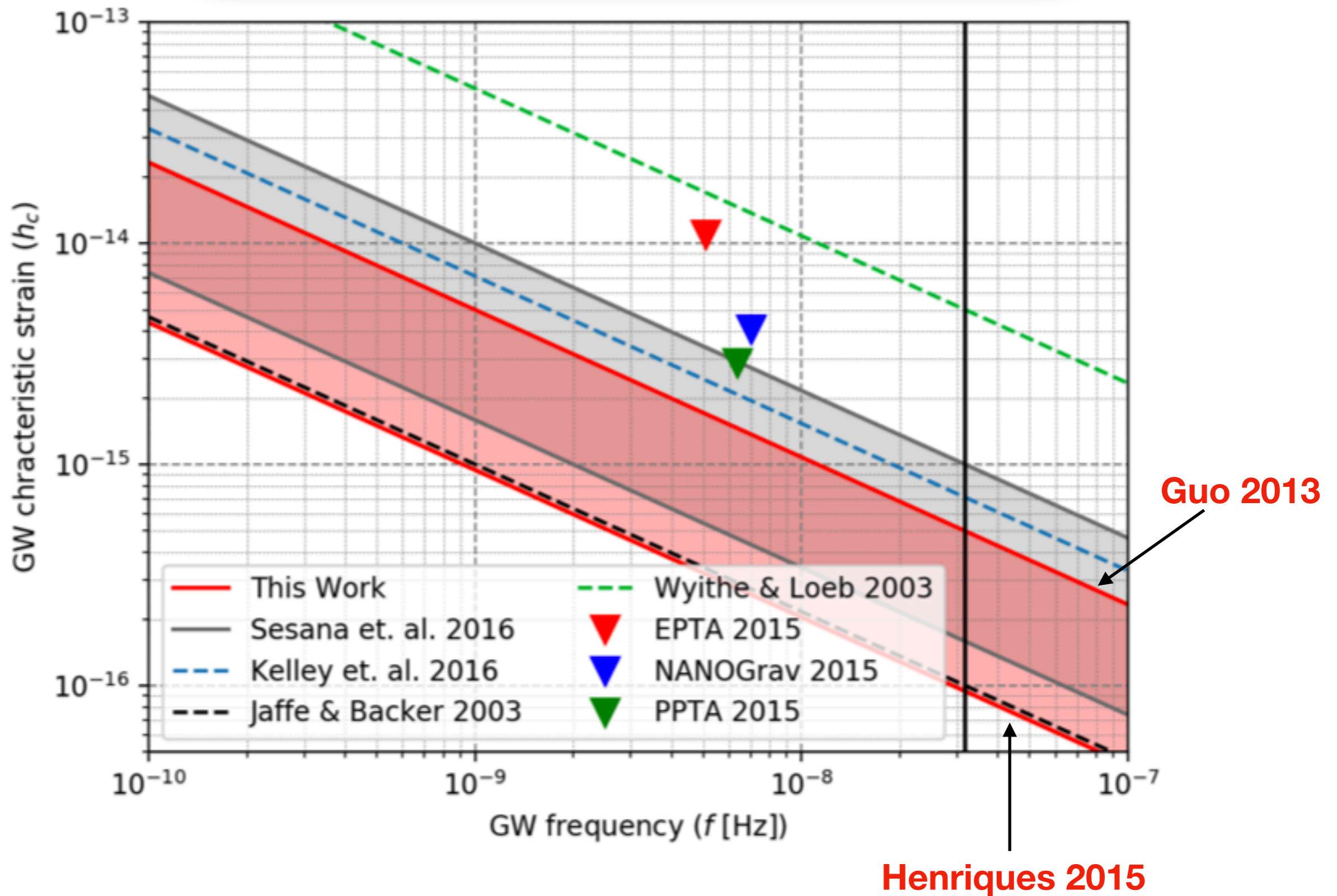


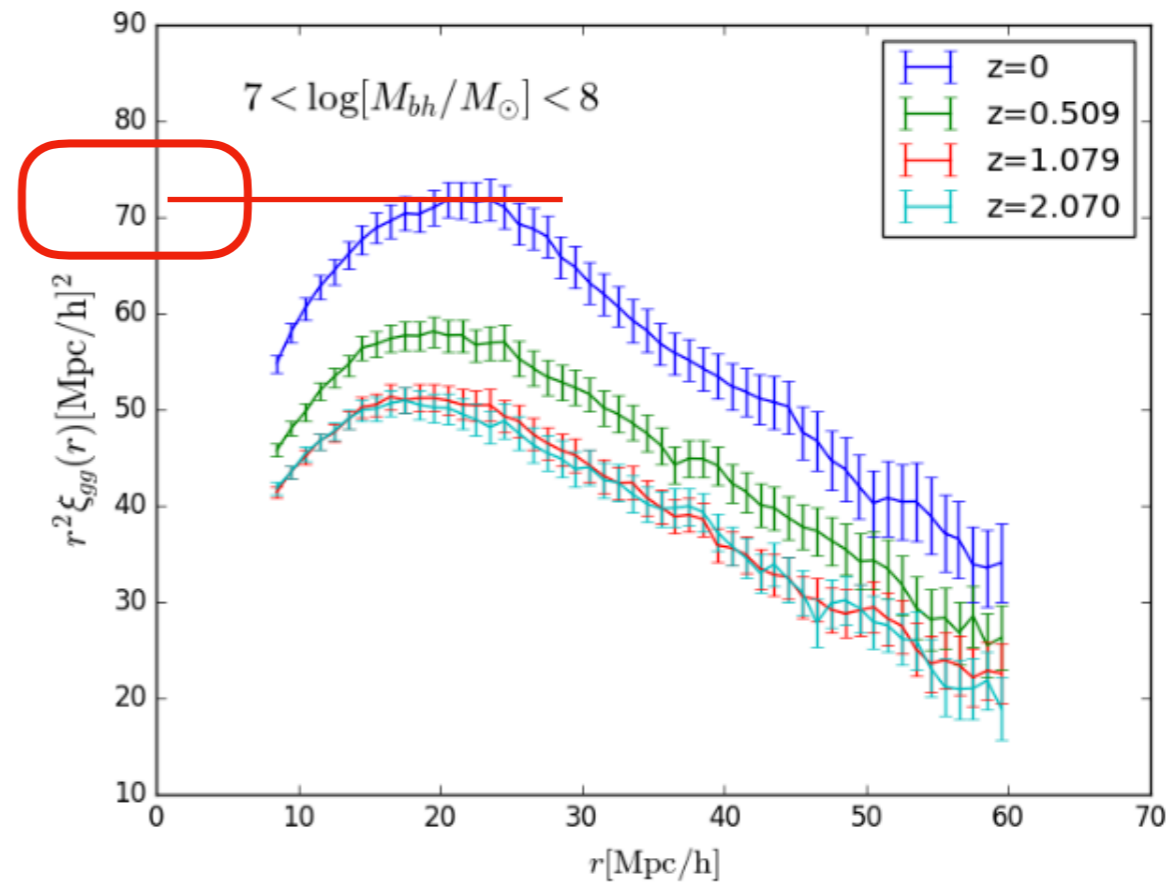
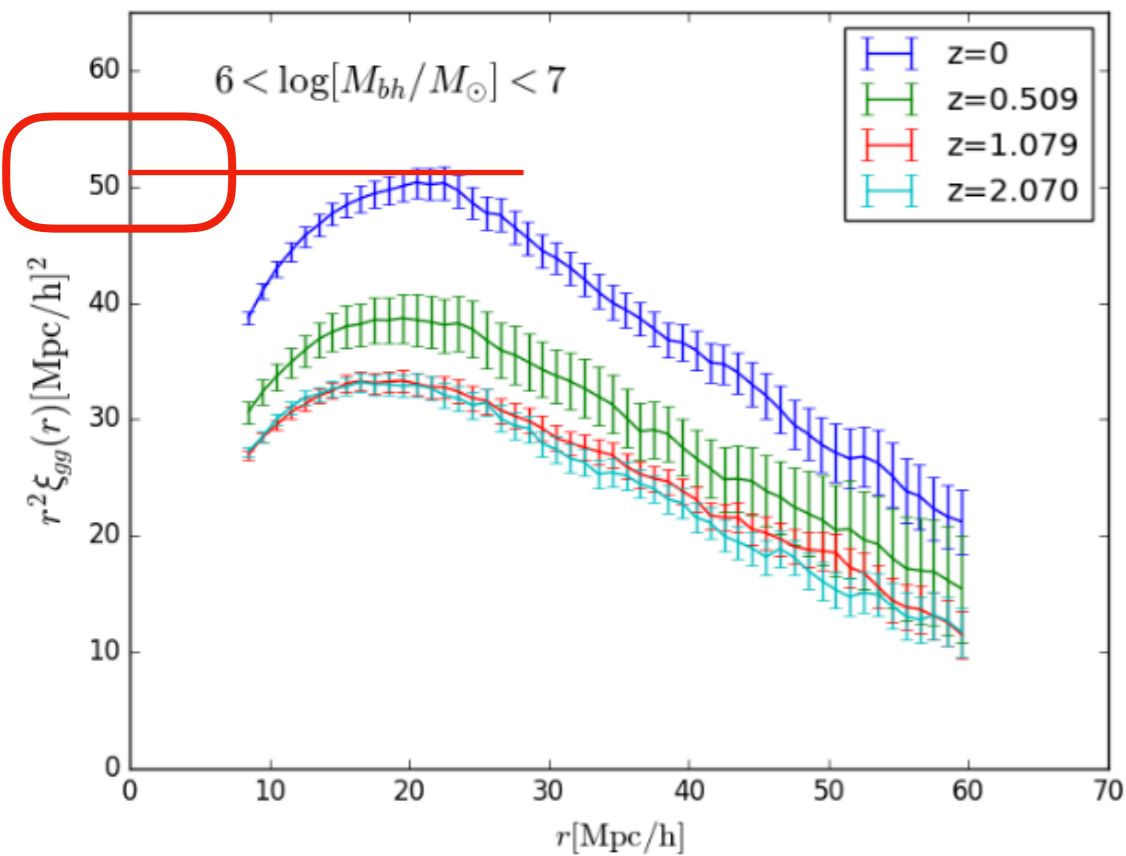
Guo 2013
based WMAP7 cosmology



Henriques 2015
based Planck cosmology

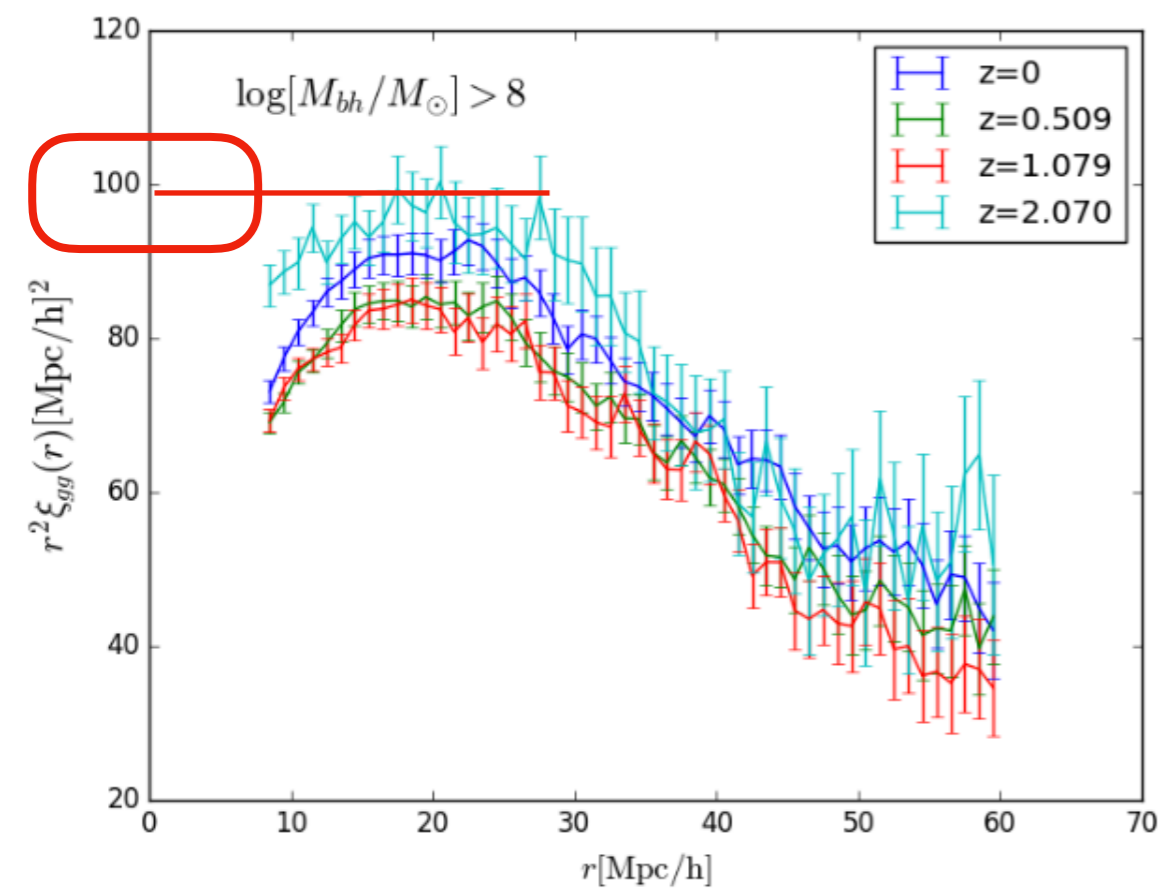
$$A_{yr^{-1}} = 5.00 \times 10^{-16} \text{ and } A_{yr^{-1}} = 9.42 \times 10^{-17}$$



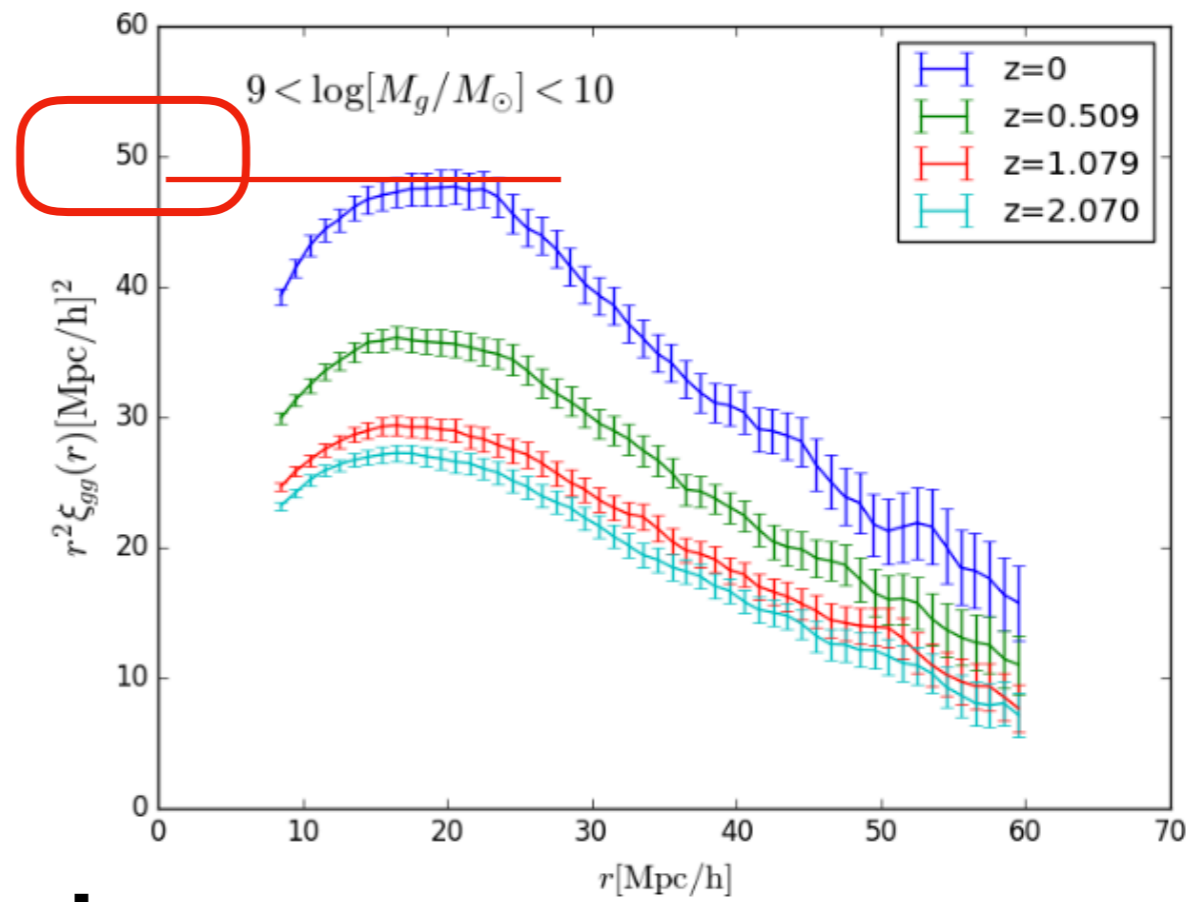
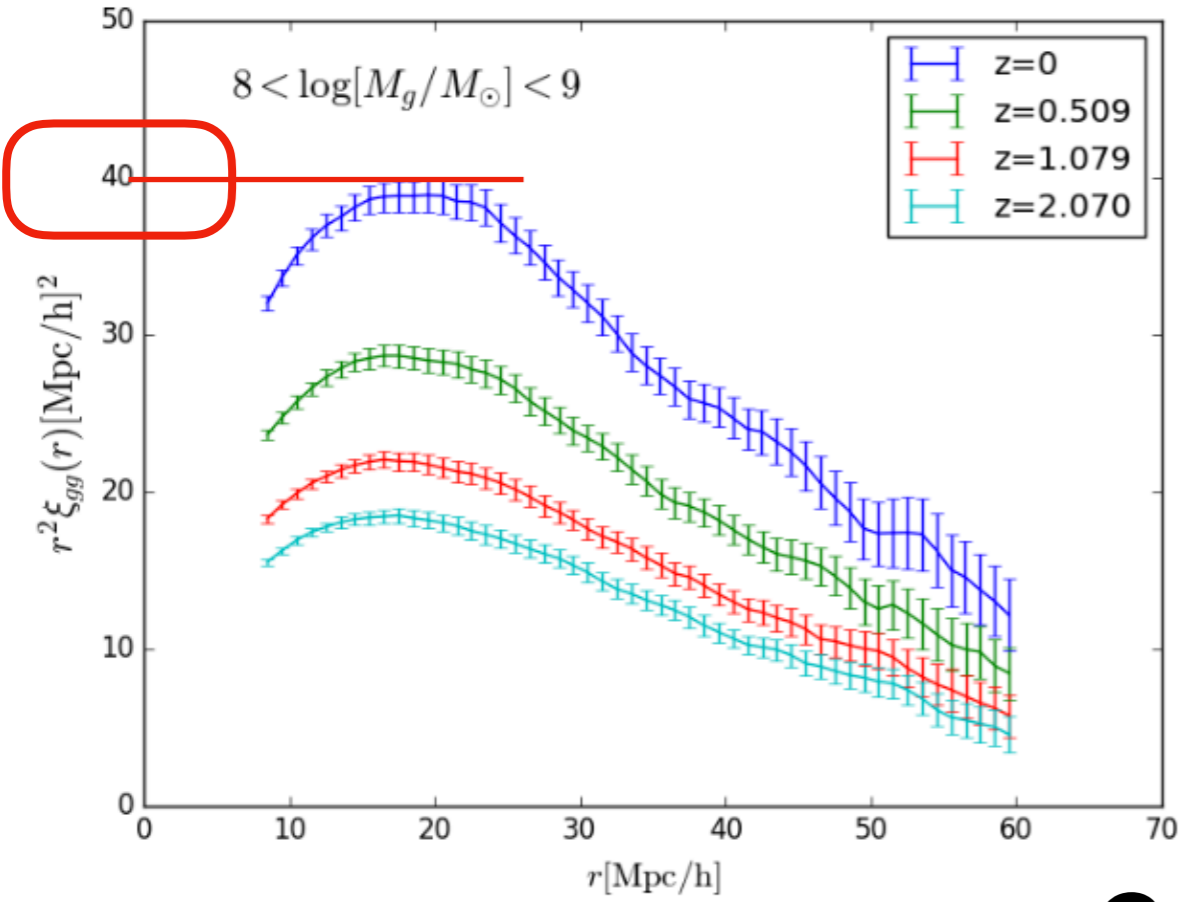


BHs

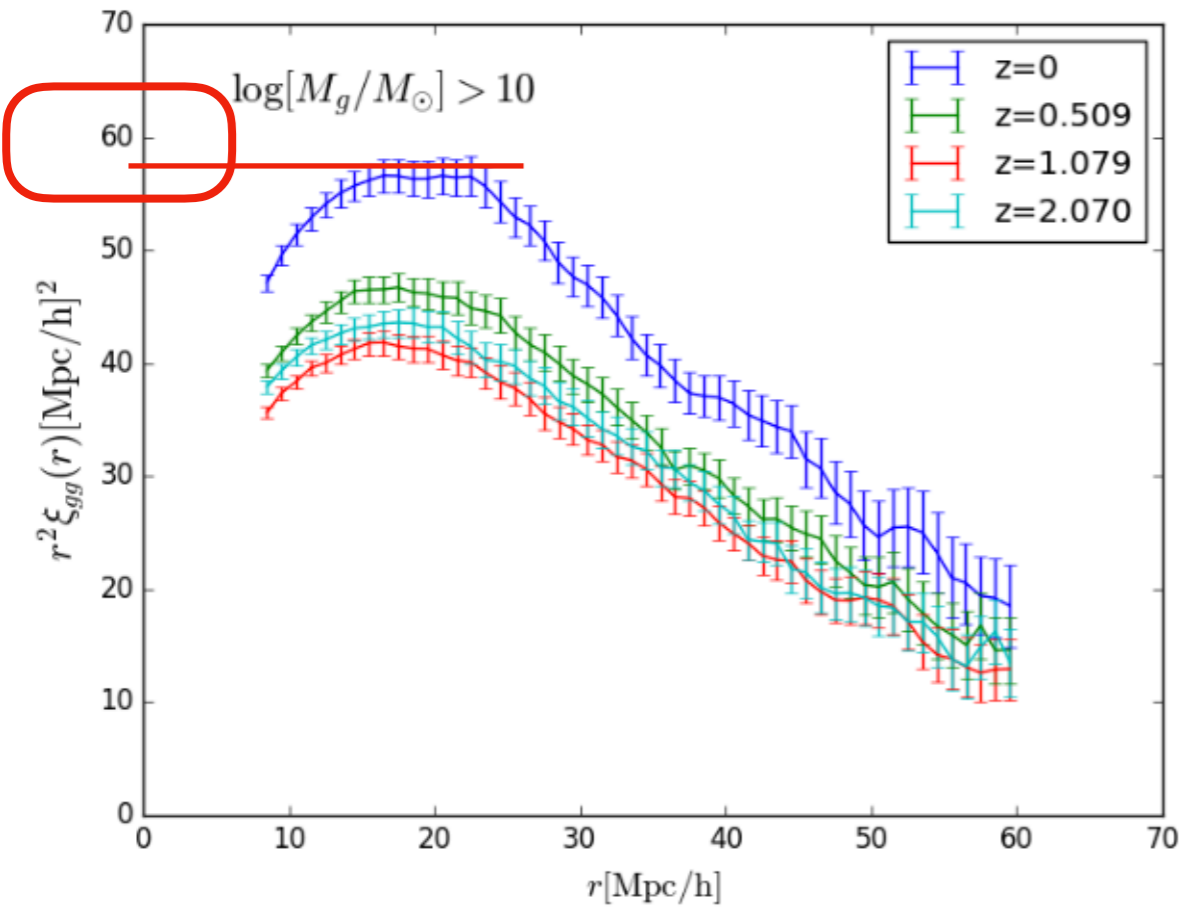
$$\xi(r) = \frac{DD(r) - RR(r)}{RR(r)}$$



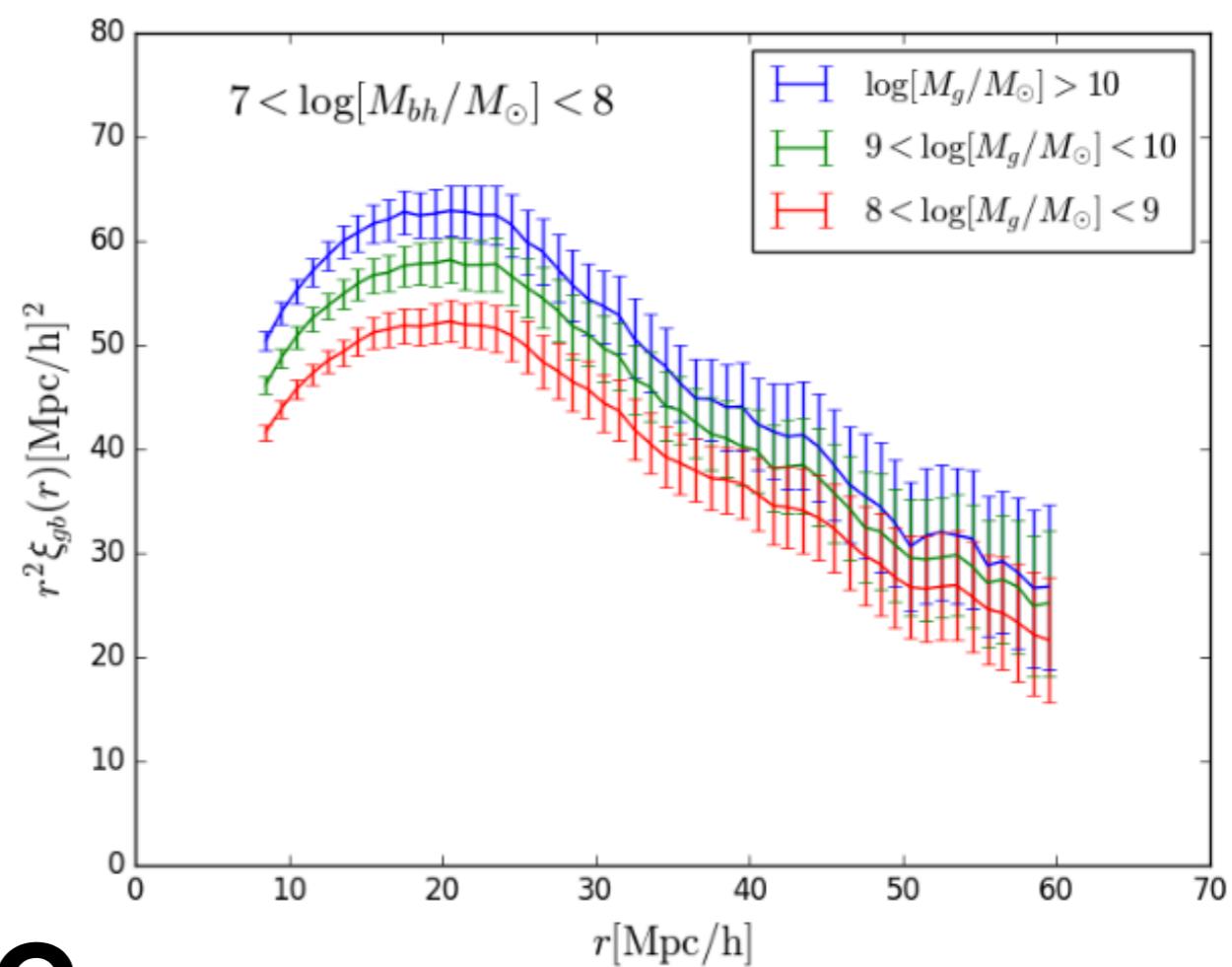
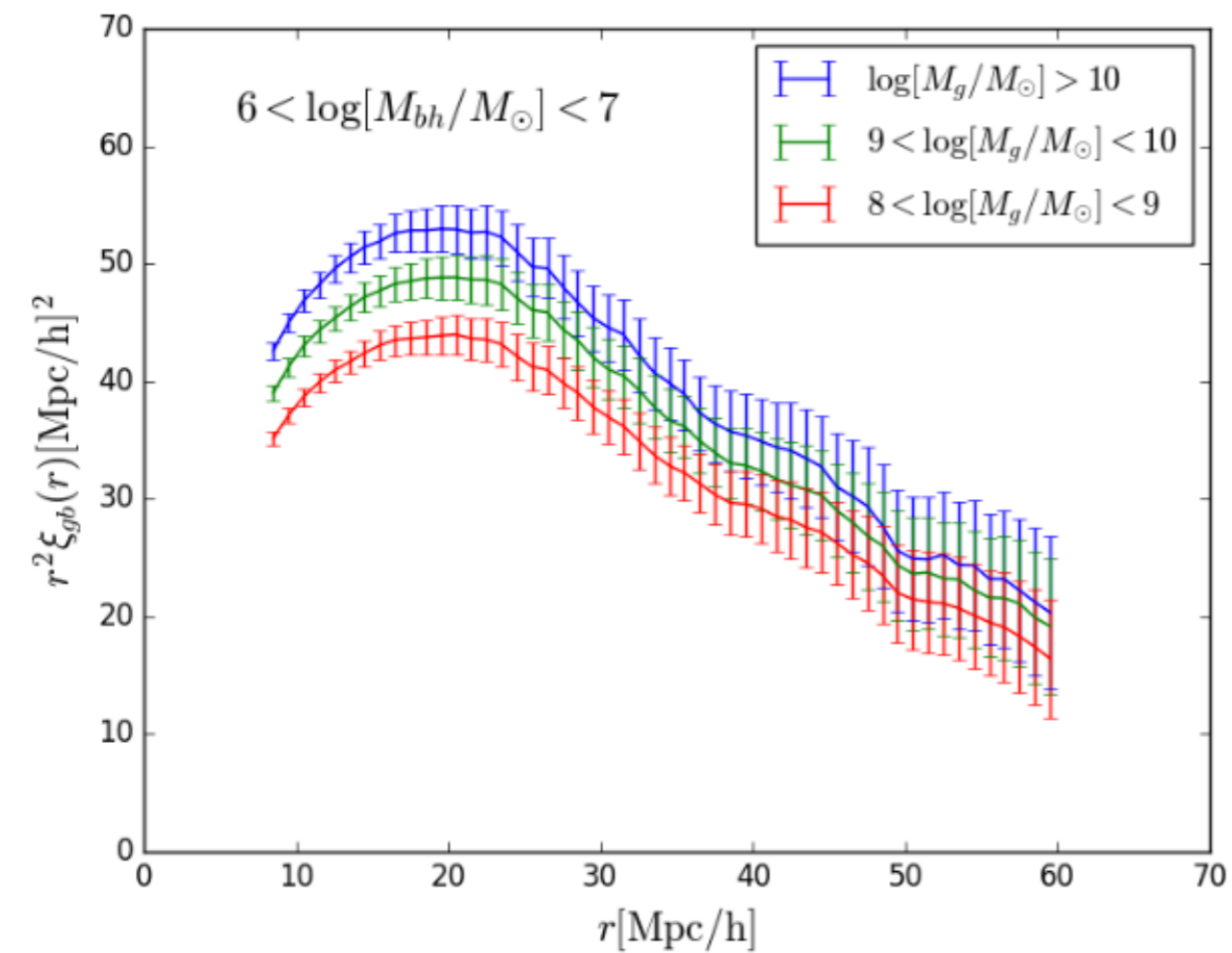
1. Clustering is enhanced in the lower redshift
2. Clustering is enhanced with mass increasing



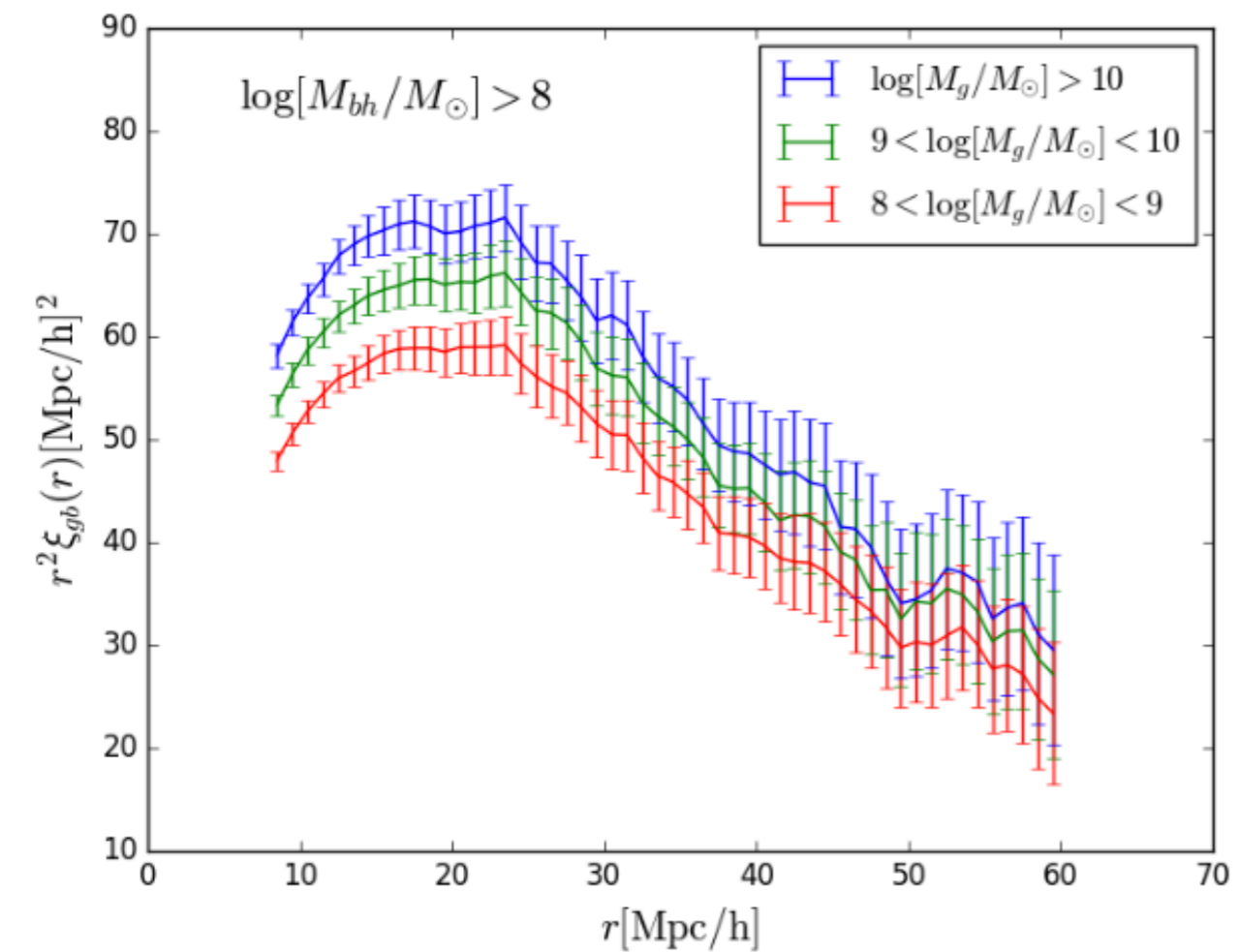
Galaxies



Same as BHs



XC



M_{bh} ↗

M_g ↗

XC ↗

Summary

1. We compare the different GW prediction from different SAM model, namely Guo 2013 & Henriques 2015.

$$A_{yr^{-1}} = 5.00 \times 10^{-16} \text{ and } A_{yr^{-1}} = 9.42 \times 10^{-17}$$

2. Clusterings of SMBHs share great similarity as galaxies:

2.1 increase with mass

2.2 enhanced at low redshift

Thanks!