

# Computation of CMB-LSS XC spectra in EFT Cosmologies-II



**Bin HU**  
*Lorentz Institute, Leiden University*

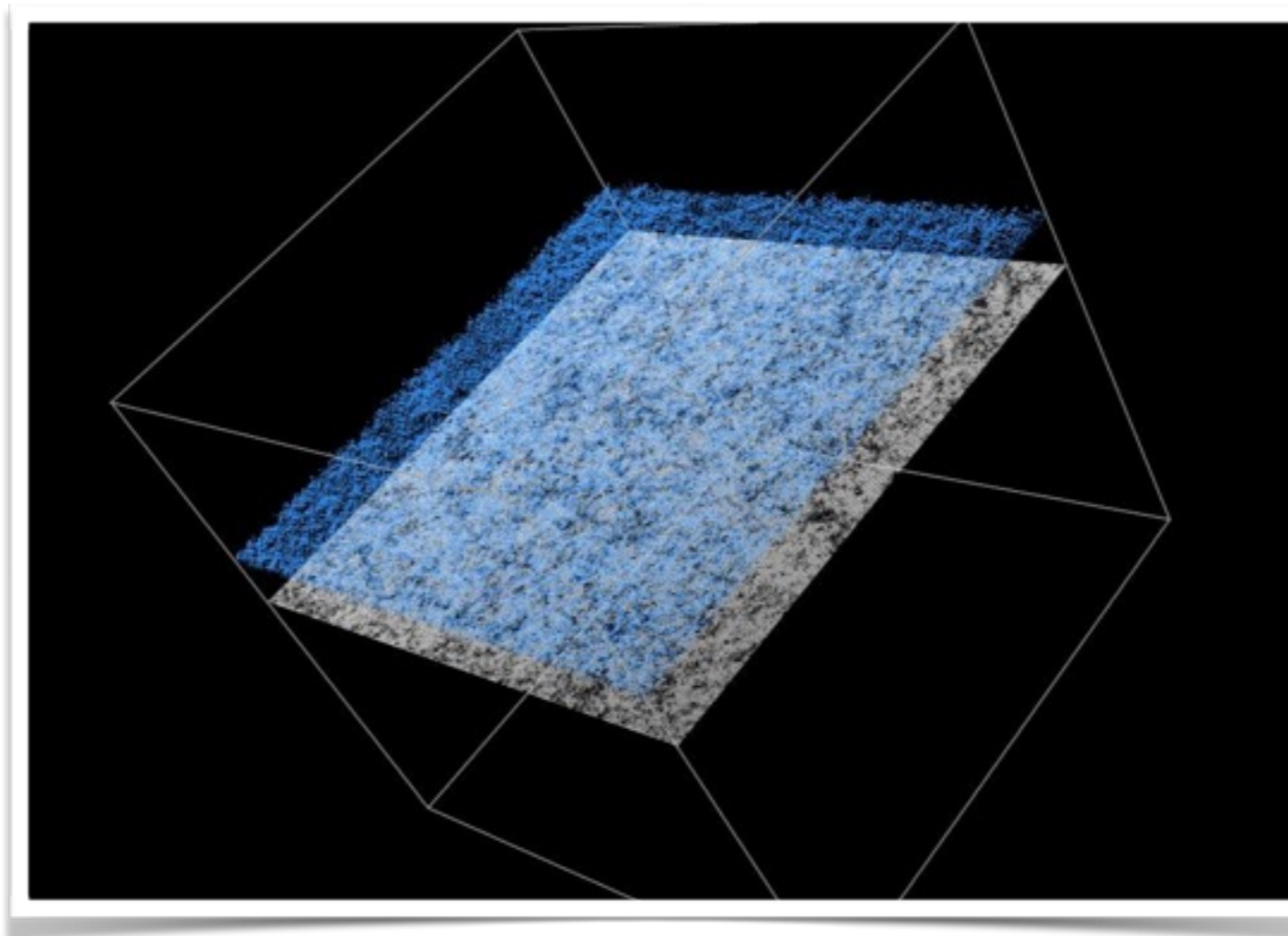
**Marco RAVERI**  
*SISSA, Trieste*

*Euclid Consortium Meeting 2015, Lausanne, June-2015*

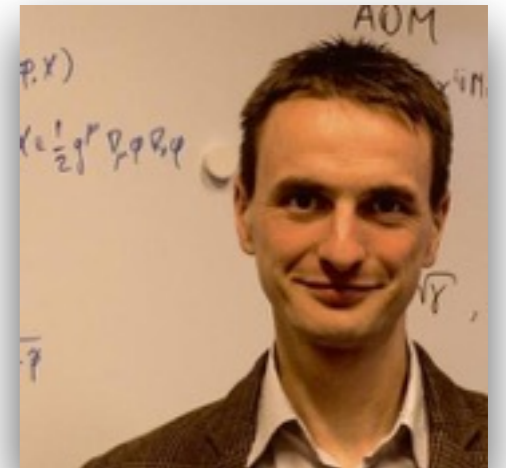




The interface between **EFTCAMB** with  
IC of N-body simulation of DE/MG—  
**FalconIC**

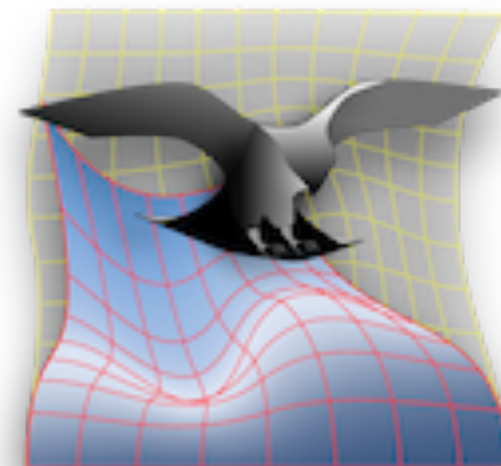


C++ code



[<http://falconic.org>]

[Wessel Valkenburg, BH, arXiv:1505.05865]



# FalconIC

(developed by Wessel Valkenburg)

- Integrated with **CAMB/CLASS/EFTCAMB**
- Work for GR and DE/MG model
- Generates IC at arbitrary scales, of arbitrary size
- Compile with MPI and OpenMP

### Parameter set 0

nGrid	256
convolveWindowFunction	<input type="checkbox"/>
boxSize	200 Mpc
randomSeed	FalconIC rocks.
zStart	49
reality	Full GR
coordinateGauge	CAMB
linearPowerSpectrum	CLASS
H0	✓ EFTCAMB
T <sub>CMB</sub>	TabulatedPLANCK2015
	2.7255 K

EFTflag	2
EFTwDE	1
PureEFTmodel $\Omega$	1
PureEFTmodelA1	0
PureEFTmodelA2	0
PureEFTmodelA3	0
PureEFTmodelA4	0

# Zeldovich Approximation

In the sub-Horizon regime, GR gives

$$\ddot{\delta}_m + 2H\dot{\delta}_m - 4\pi G\rho_m\delta_m = 0$$

The growth rate of CDM only depends on time!

The displacement field  $\vec{x} = \vec{y} - \mathcal{D}(\tau) \frac{1}{\nabla_y^2} \vec{\nabla}_y \Delta_c(\tau_i, \vec{y})$

In GR: CDM particles trajectory is **straight** line!

$$k^2\psi = -4\pi G \mu(a, k) a^2 \rho \Delta ,$$
$$\frac{\phi}{\psi} = \gamma(a, k) .$$

DE/MG:

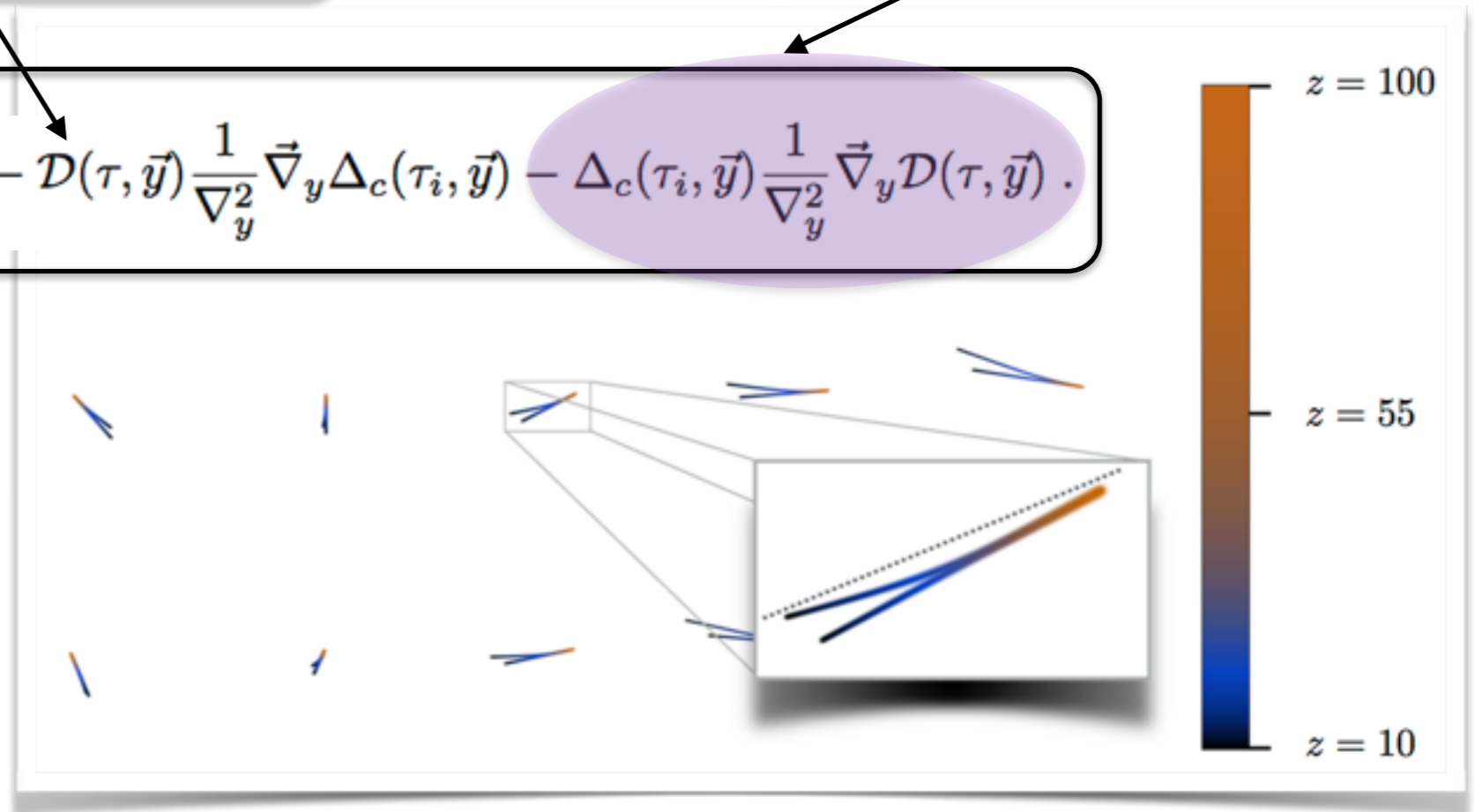
$$\ddot{\delta}_m + 2H\dot{\delta}_m - 4\pi G_{\text{eff}}(t, k)\rho_m\delta_m = 0$$

# Beyond Zeldovich Approximation

DE/MG: at linear regime  
growth rate of CDM  
**depends** on the scales!

Deflection by the  
gravitational potential

$$\vec{x} = \vec{y} - \mathcal{D}(\tau, \vec{y}) \frac{1}{\nabla_y^2} \vec{\nabla}_y \Delta_c(\tau_i, \vec{y}) - \Delta_c(\tau_i, \vec{y}) \frac{1}{\nabla_y^2} \vec{\nabla}_y \mathcal{D}(\tau, \vec{y}) .$$



Even at linear regime,  
trajectory of CDM particles are **curved**!

# Modified Einstein Eq.

$$m_0^2(1 + \Omega)G_{\mu\nu}[g_{\mu\nu}] = T_{\mu\nu}^{(m)}[\rho_m, \theta_m, \dots] + T_{\mu\nu}^{(\pi)}[\pi, \dot{\pi}, \dots],$$

Define a conserved Fluid EMT ( $\nabla^\nu T_{\mu\nu}^{(Q)} = 0$ )

$$m_0^2 G_{\mu\nu}[g_{\mu\nu}] = T_{\mu\nu}^{(m)}[\rho_m, \theta_m, \dots] + T_{\mu\nu}^{(Q)}[\rho_\pi, \theta_\pi, \rho_m, \dots],$$

$$T_{\mu\nu}^{(Q)}[\rho_\pi, \theta_\pi, \rho_m, \dots] \equiv \frac{1}{1 + \Omega} \left\{ -\Omega T_{\mu\nu}^{(m)}[\rho_m, \theta_m, \dots] + T_{\mu\nu}^{(\pi)}[\pi, \dot{\pi}, \dots] \right\}.$$

At **linear** order the fluid variables

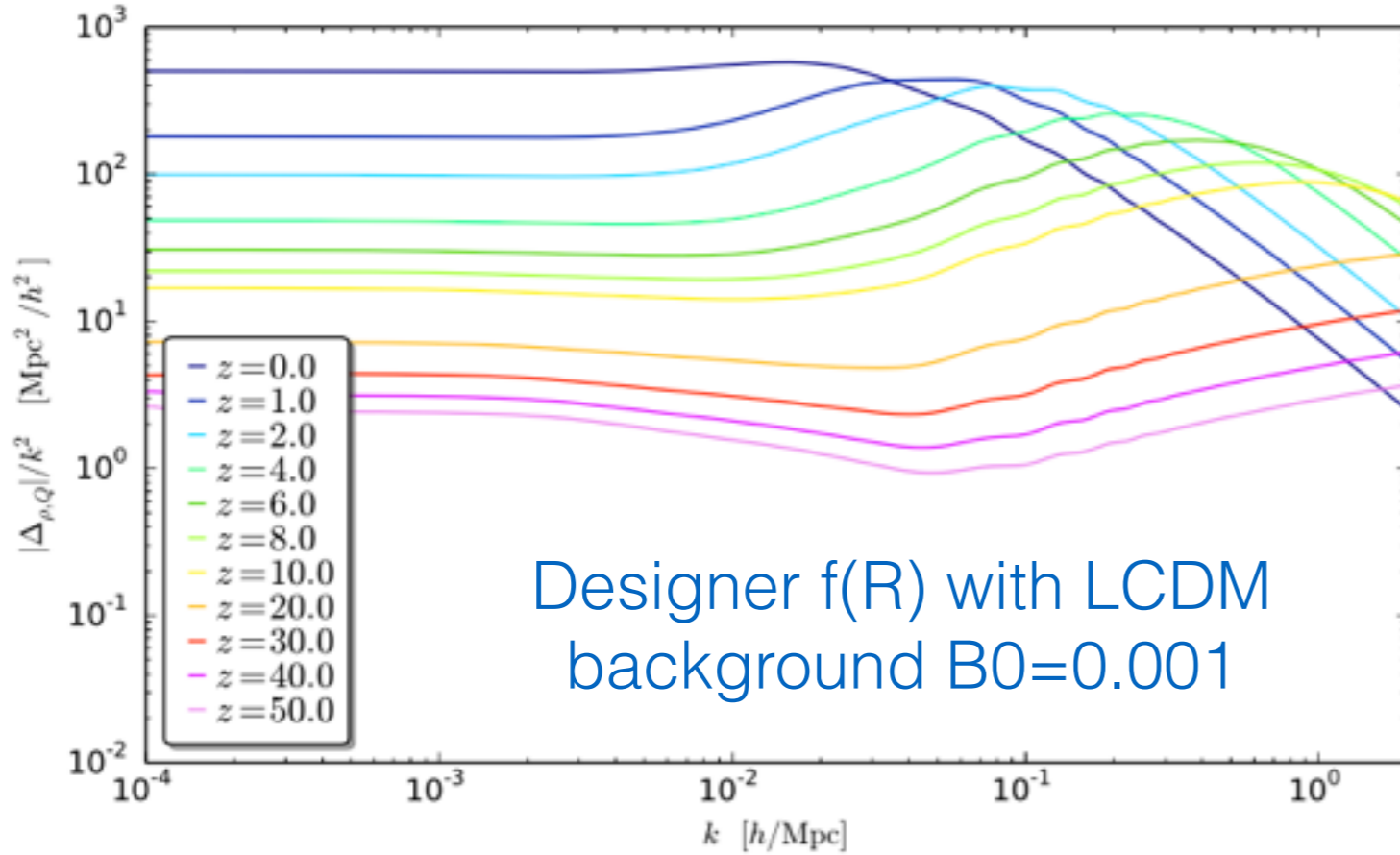
$$\delta\rho_Q^{(\text{syn})} = \frac{1}{(1 + \Omega)} \left\{ -\Omega\delta\rho_m^{(\text{syn})} + \dot{\rho}_Q\pi + 2c(\dot{\pi}^{(\text{syn})} + \mathcal{H}\pi^{(\text{syn})}) - \frac{2m_0^2}{a^2} \left[ \frac{\dot{\Omega}}{4}\dot{h} + \frac{\dot{\Omega}}{2} \left( 3(3\mathcal{H}^2 - \dot{\mathcal{H}})\pi^{(\text{syn})} + 3\mathcal{H}\dot{\pi}^{(\text{syn})} + k^2\pi^{(\text{syn})} \right) \right] \right\} \quad \left| \quad (\rho_{\text{DE}} + P_{\text{DE}})\theta_Q^{(\text{syn})} = \frac{1}{1 + \Omega} \left[ -\Omega(\rho_m + P_m)\theta_m^{(\text{syn})} + (\rho_Q + P_Q)\theta_m^{(\text{syn})} + \frac{2m_0^2}{a^2} k^2 \dot{\Omega} (\dot{\pi}^{(\text{syn})} + \mathcal{H}\pi^{(\text{syn})}) \right], \right.$$

$$\delta P_Q^{(\text{syn})} = \frac{1}{1 + \Omega} \left\{ -\Omega\delta P_m^{(\text{syn})} + P_Q\dot{\pi}^{(\text{syn})} + (\rho_Q + P_Q)(\dot{\pi}^{(\text{syn})} + \mathcal{H}\pi^{(\text{syn})}) + \frac{m_0^2}{a^2} \left[ \frac{1}{3}\dot{\Omega}\dot{h} + \dot{\Omega}\ddot{\pi}^{(\text{syn})} + (\ddot{\Omega} + 3\mathcal{H}\dot{\Omega})\dot{\pi}^{(\text{syn})} + \left( \mathcal{H}\ddot{\Omega} + 5\mathcal{H}^2\dot{\Omega} + \dot{\mathcal{H}}\dot{\Omega} + \frac{2}{3}k^2\dot{\Omega} \right) \pi^{(\text{syn})} \right] \right\}$$

$$(\rho_{\text{DE}} + P_{\text{DE}})\sigma_Q^{(\text{syn})} = \frac{1}{1 + \Omega} \left[ -\Omega(\rho_m + P_m)\sigma_m^{(\text{syn})} + \frac{m_0^2}{3a^2} \dot{\Omega} \left( \dot{h} + 6\dot{\eta} + 2k^2\pi^{(\text{syn})} \right) \right]$$



# Transfer function of Q-fluid



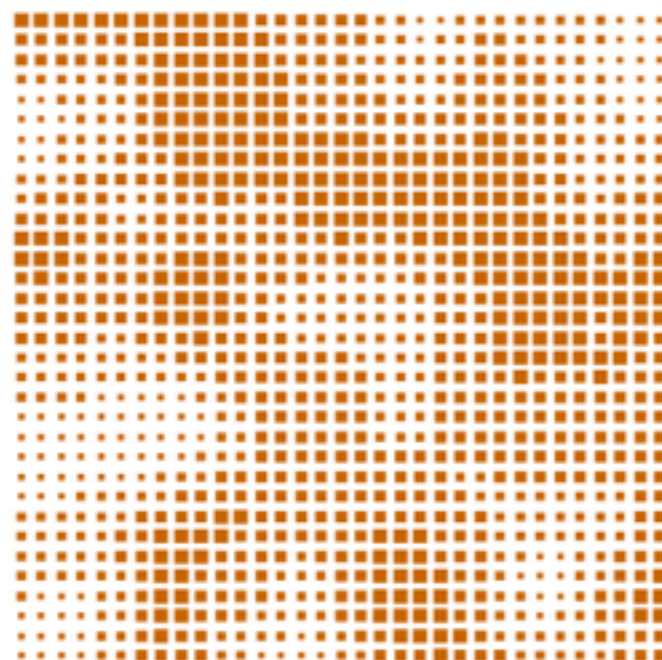
Take Newtonian limit

$$-k^2\psi = \frac{16\pi G}{3}\delta\rho_{\text{cdm}} - \frac{\delta R}{6}.$$

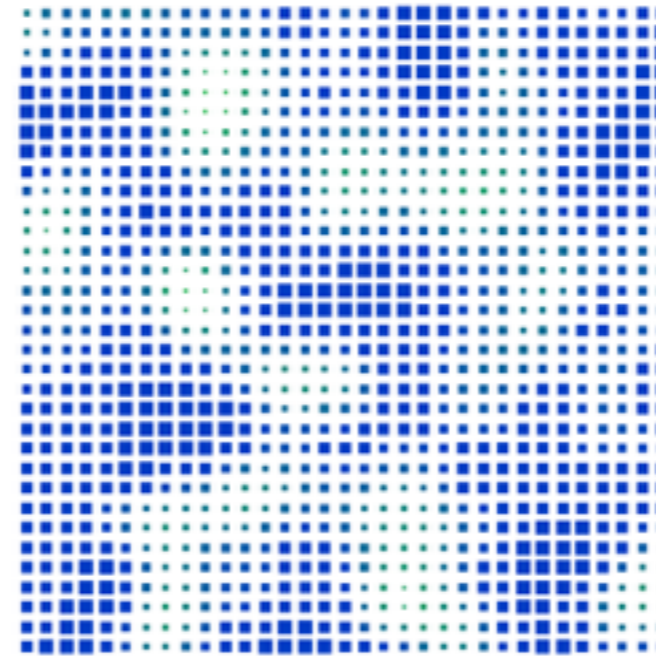
$$-k^2\psi = \frac{16\pi G}{3}(\delta\rho_{\text{cdm}} + \delta\rho_Q)$$

$$\delta\rho_Q \propto -\delta R$$

In the CDM over dense regime, Q-fluid is under dense!

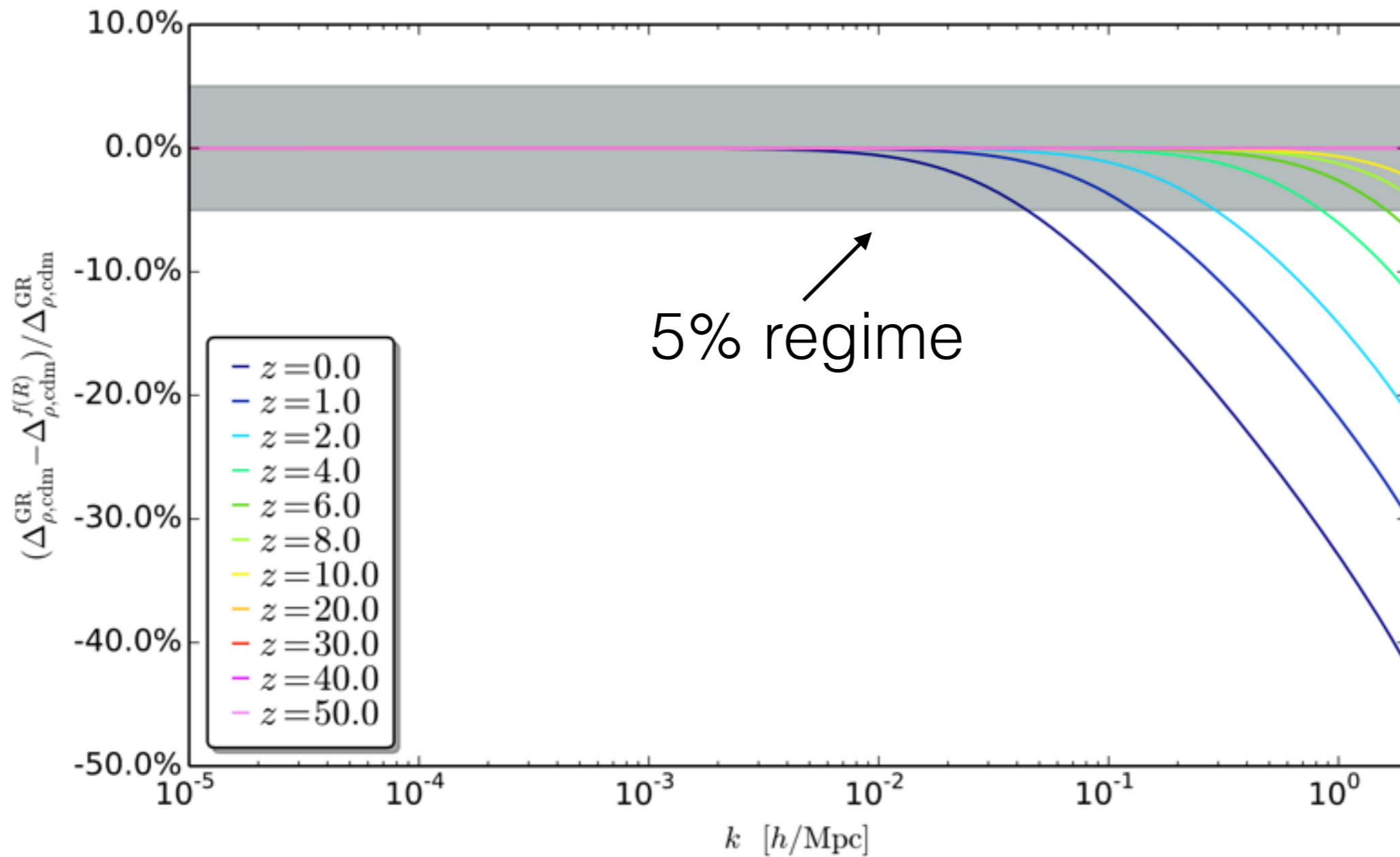


Dark Matter (Eulerian)



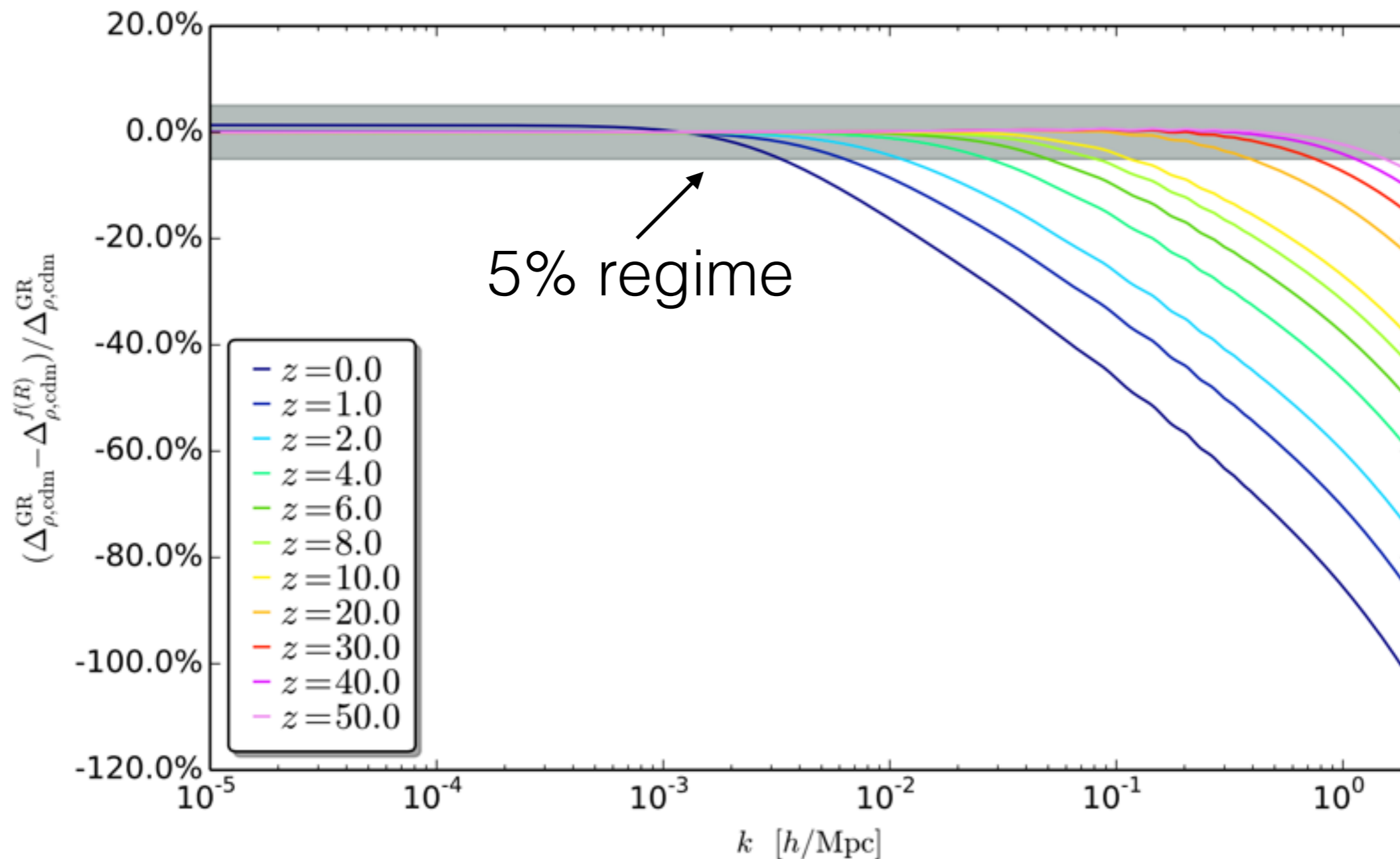
Q-fluid (Eulerian)

# Transfer function of CDM



Designer  $f(R)$  with LCDM background  
 $B_0 = 0.001$

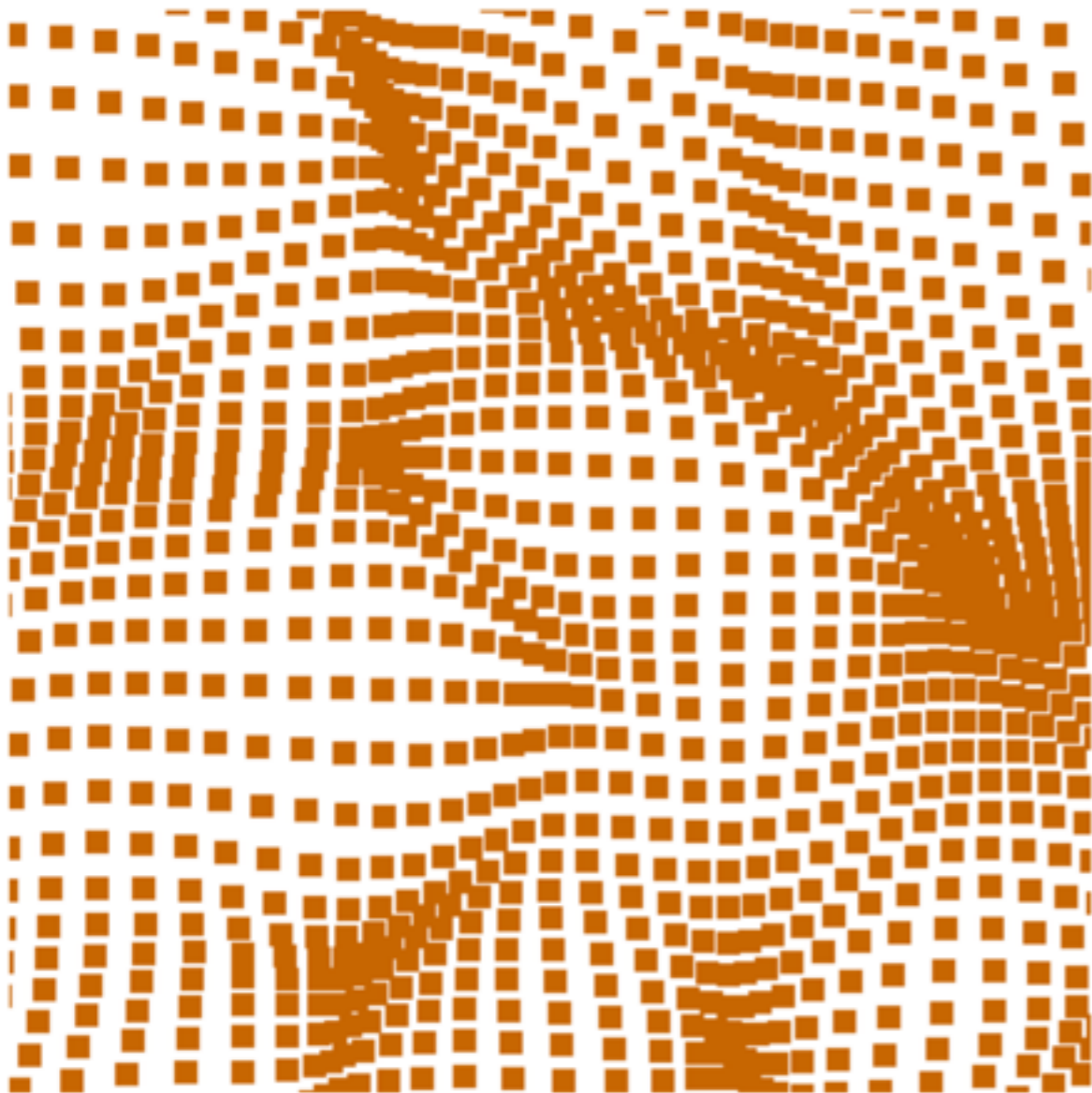




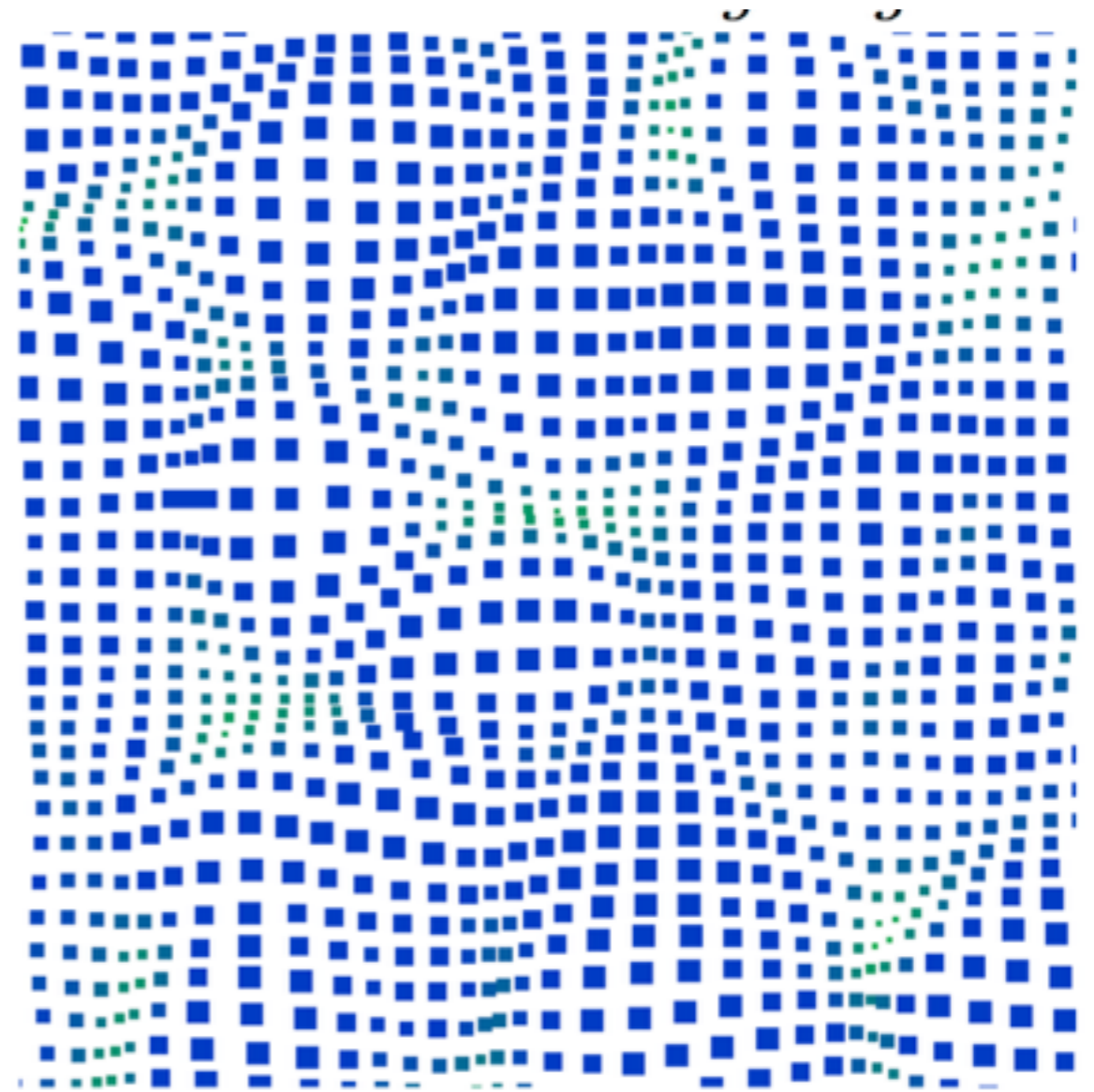
Designer  $f(R)$  with  $w$ CDM background  
 $B_0=0.01$  and  $w=-0.95$

(box size  $\rightarrow$  particle mass)

(box color  $\rightarrow$  pressure)



Dark Matter



Q-particles (Langrangian)

- CDM particle mass is conserved, Pressureless
- Q-paricle mass is non-conserved, Pressure

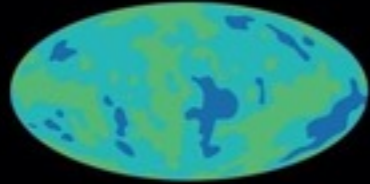


# Conclusion

- The EFT of Cosmic Acceleration provides a generic and powerful framework to efficiently study DE/MG
- The EFT framework has been implemented in the Einstein/Boltzmann code CAMB, EFTCAMB (HiCLASS, E. Bellini, M. Zumalacarregui et. al.)
- EFTCAMB is publicly available, does not rely on QSA
- XC is a valuable complementary probe for DE/MG test
- IC for N-body can be important for some DE/MG

A graphic consisting of a series of curved lines in shades of blue and green, radiating from the top left and curving towards the right, resembling a stylized wing or a field of lines.

EFTCAMB



KEEP  
CALM  
AND  
TEST  
GRAVITY

the EFTCAMB team

Thank you!

